

The GEOMETRY of the OCTONIONS



Tevian Dray

- I: Octonions
- II: Rotations
- III: Lorentz Transformations
- IV: Dirac Equation
- V: Exceptional Groups

DIVISION ALGEBRAS

Real Numbers:

$$\mathbb{R}$$

Complex Numbers:

$$\mathbb{C} = \mathbb{R} + \mathbb{R}i$$

$$z = x + yi$$

$$i^2 = -1$$

DIVISION ALGEBRAS

Real Numbers:

$$\mathbb{R}$$

Quaternions:

$$\mathbb{H} = \mathbb{C} + \mathbb{C}j$$
$$q = (a + bi) + (c + di)j$$

Complex Numbers:

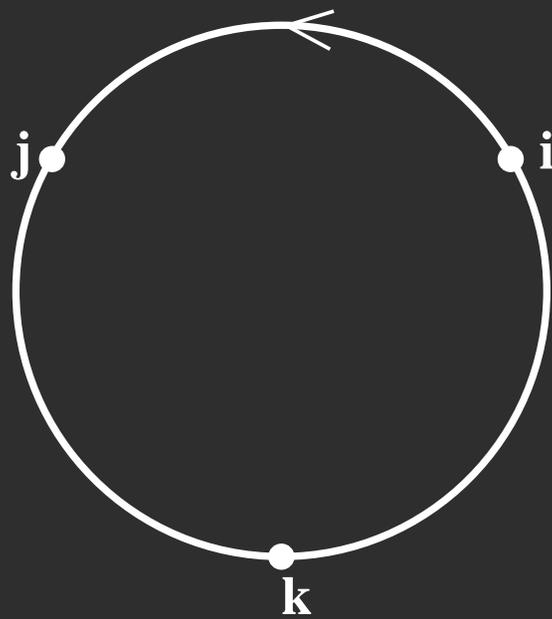
$$\mathbb{C} = \mathbb{R} + \mathbb{R}i$$
$$z = x + yi$$

$$i^2 = j^2 = -1$$

● i

j●

●i



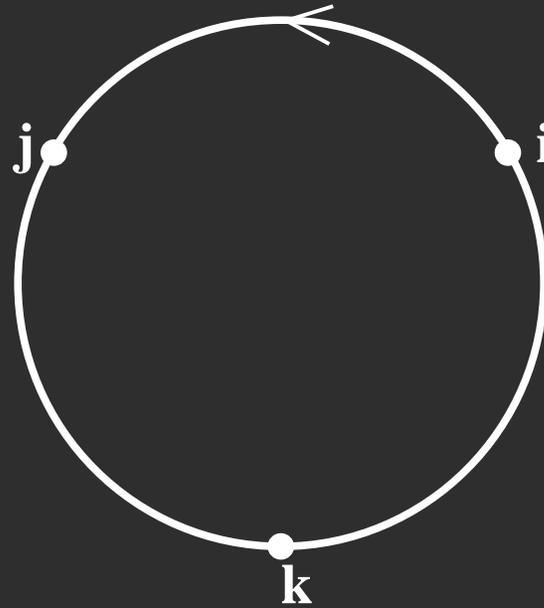
QUATERNIONS

$$k^2 = -1$$

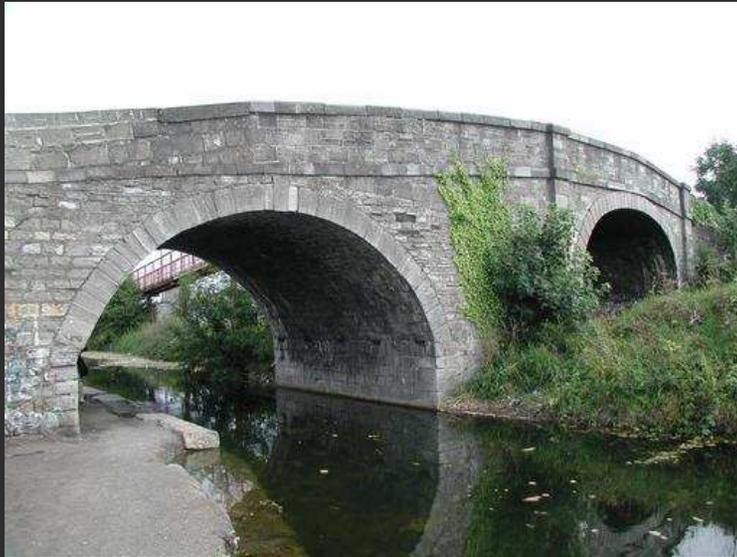
$$ij = +k$$

$$ji = -k$$

not commutative



THE DISCOVERY OF THE QUATERNIONS



Brougham Bridge (Dublin)



Here as he walked by
on the 16th of October 1843
Sir William Rowan Hamilton
in a flash of genius discovered
the fundamental formula for
quaternion multiplication
 $i^2 = j^2 = k^2 = ijk = -1$
& cut it on a stone of this bridge

VECTORS I

$$v = bi + cj + dk \longleftrightarrow \vec{v} = b\hat{i} + c\hat{j} + d\hat{k}$$

$$vw \longleftrightarrow -\vec{v} \cdot \vec{w} + \vec{v} \times \vec{w}$$

Dot product exists in any dimension

but

Cross product exists only in 3 and 7 dimensions

DIVISION ALGEBRAS

Real Numbers:

$$\mathbb{R}$$

Quaternions:

$$\begin{aligned}\mathbb{H} &= \mathbb{C} + \mathbb{C}j \\ q &= (a + bi) + (c + di)j\end{aligned}$$

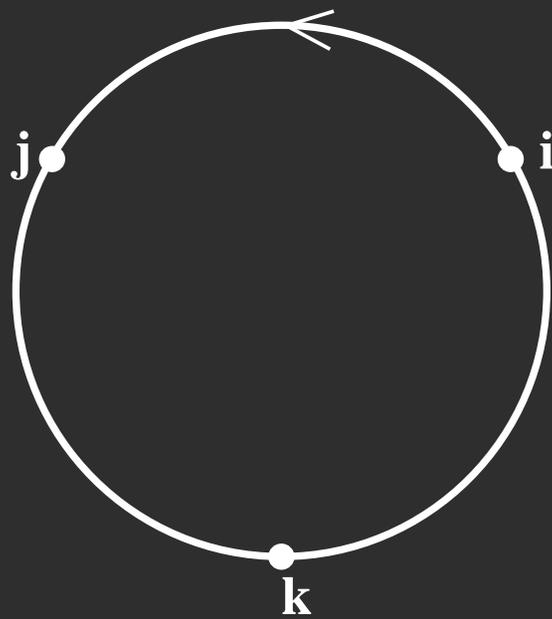
Complex Numbers:

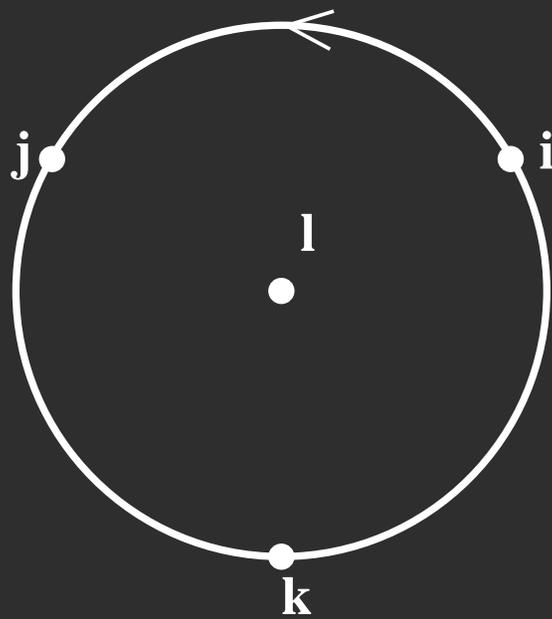
$$\begin{aligned}\mathbb{C} &= \mathbb{R} + \mathbb{R}i \\ z &= x + yi\end{aligned}$$

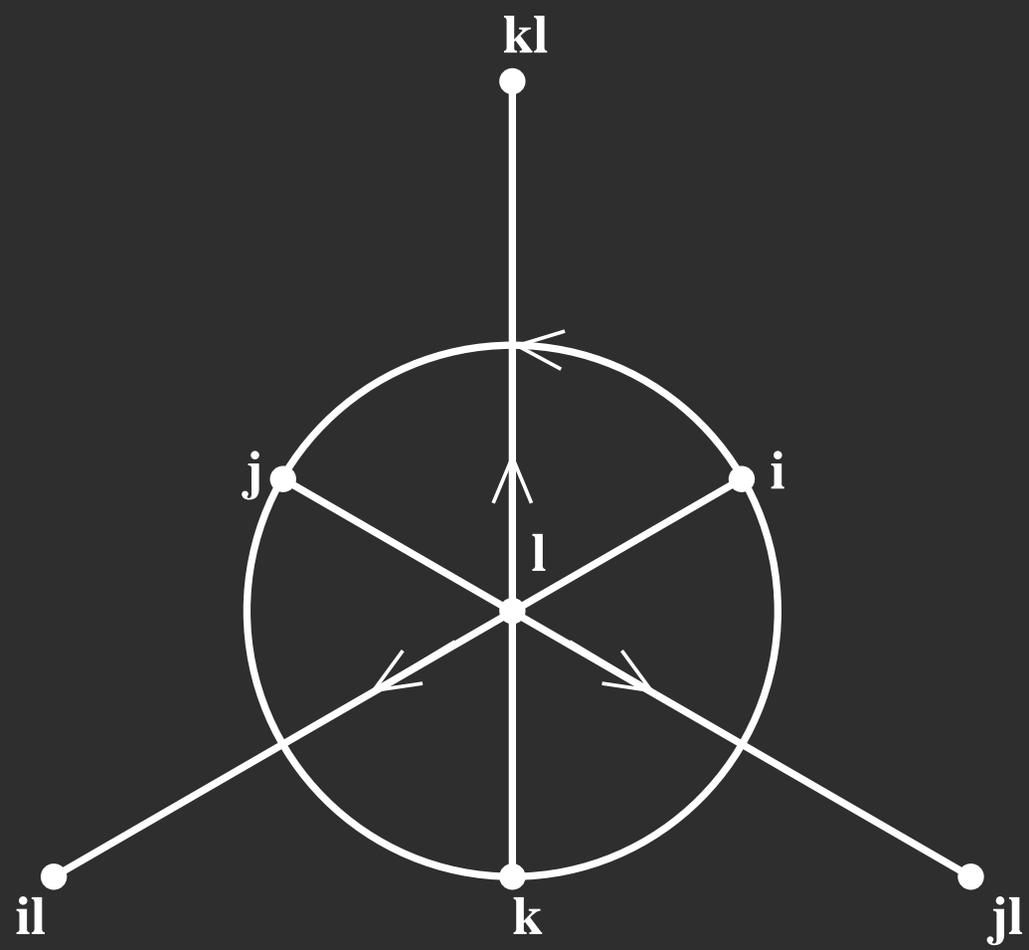
Octonions:

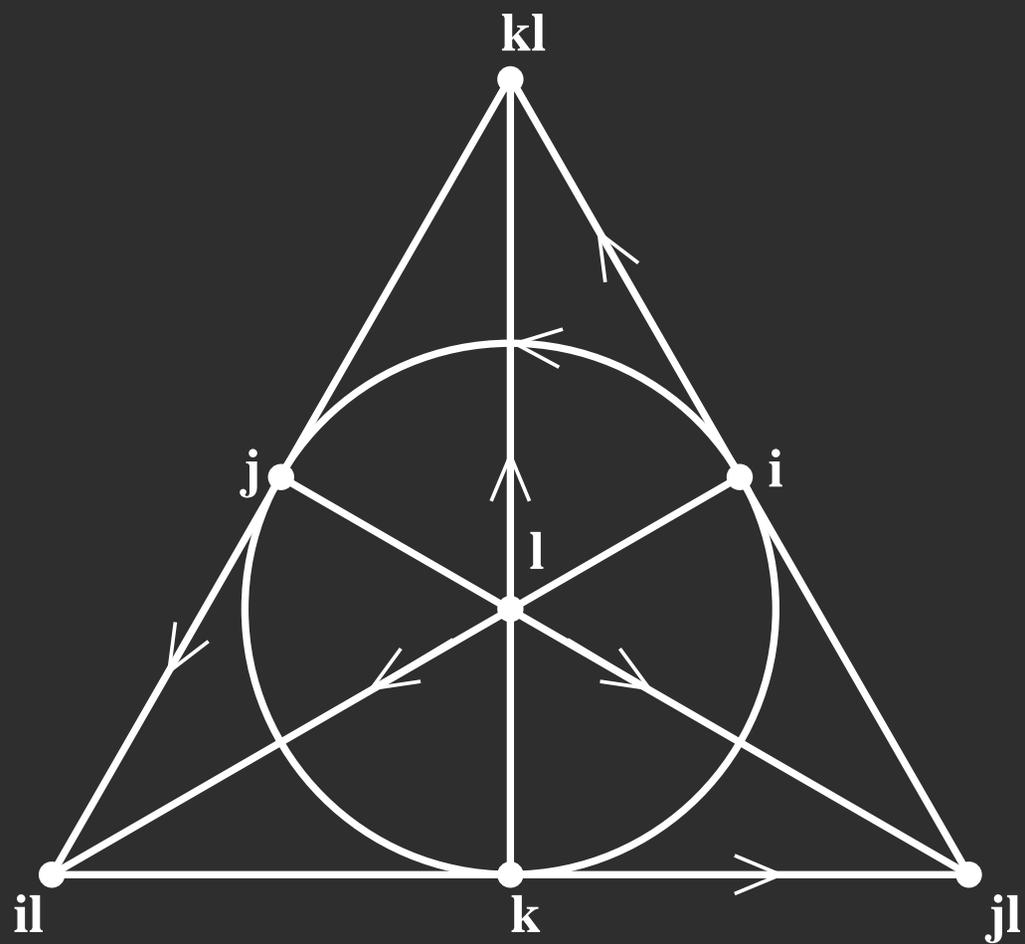
$$\mathbb{O} = \mathbb{H} + \mathbb{H}l$$

$$i^2 = j^2 = l^2 = -1$$







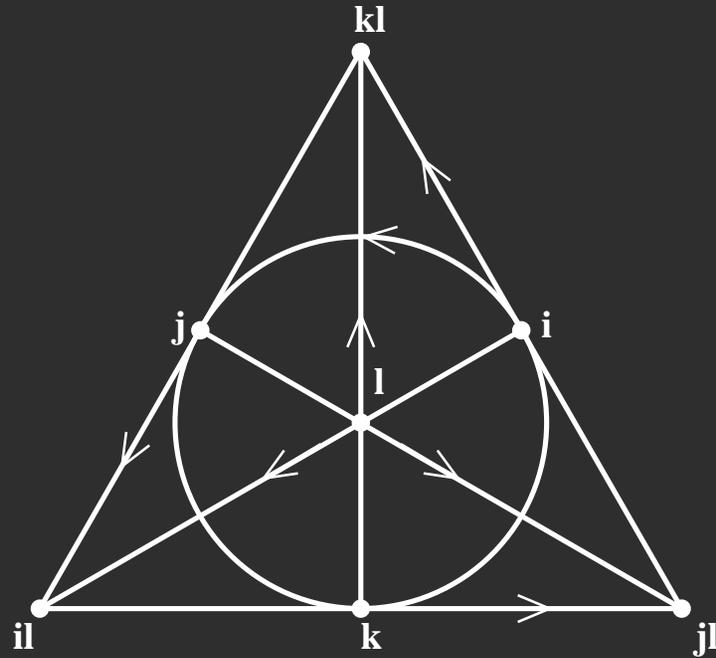


OCTONIONS

each line is quaternionic

$$(ij)l = +kl$$
$$i(jl) = -kl$$

not associative



THE DISCOVERY OF THE OCTONIONS



Brougham Bridge (Dublin)



John T. Graves (1843!)
Arthur Cayley (1845)
octaves, Cayley numbers

WHAT STILL WORKS?

- $q \in \text{Im } \mathbb{O} \implies \bar{q} = -q$
- $q\bar{q} = |q|^2$
- $q^{-1} = \frac{\bar{q}}{|q|^2} \quad (q \neq 0)$
- $|pq| = |p||q|$
- $[p, p, q] = (pp)q - p(pq) = 0$

required for supersymmetry

alternativity

EXPONENTIAL FORM

$$\begin{aligned} p &= |p| e^{s\phi} \quad (s^2 = -1) \\ &= |p| (\cos \phi + s \sin \phi) \end{aligned}$$

$$e^{k\phi i} = i e^{-k\phi}$$

$$e^{k\phi i} e^{-k\phi} = i e^{-2k\phi}$$

ROTATIONS

$$x \in \mathbb{H} \longleftrightarrow \vec{x} \in \mathbb{R}^4$$

$$|x| \longleftrightarrow |\vec{x}|$$

$$|p| = 1 \implies |px| = |x|$$

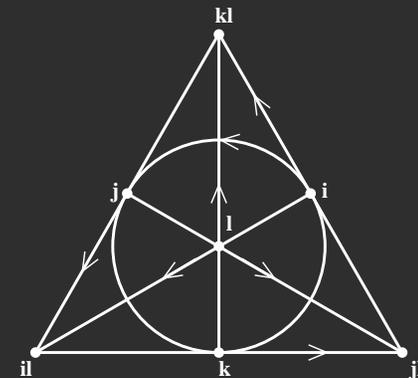
SO(3):

θ about k -axis: $e^{k\theta}x$

SO(7)?

$e^{k\theta}$ rotates 3 planes perpendicular to k ...

Need *Flips*



ROTATIONS

$$x \in \mathbb{O} \longleftrightarrow \vec{x} \in \mathbb{R}^8$$

$$|x| \longleftrightarrow |\vec{x}|$$

$$|p| = 1 \implies |px\bar{p}| = |x|$$

SO(7):

2θ about “ k -axis”: $e^{k\theta} x e^{-k\theta}$

flips (2θ in ij -plane):

$$(i \cos \theta + j \sin \theta)(ixi)(i \cos \theta + j \sin \theta)$$

nesting!

ROTATIONS

$$x \in \mathbb{O} \longleftrightarrow \vec{x} \in \mathbb{R}^8$$

$$|x| \longleftrightarrow |\vec{x}|$$

$$|p| = 1 \implies |pxp| = |x|$$

$SO(8)$:

flips OK

2θ in 1ℓ -plane: $e^{\ell\theta} x e^{\ell\theta}$

triality: pxp, px, xp

fails over \mathbb{H} !

AUTOMORPHISMS

$$|pxp^{-1}| = |x|$$

$SO(3)$:

over \mathbb{H} : $(pxp^{-1})(pyp^{-1}) = p(xy)p^{-1} \quad (\forall p)$

$G_2 \subset SO(7)$:

$$p = e^{\pi s/3} \quad (s^2 = -1)$$

$$\dim G_2 = 14$$

$$SU(3) \subset G_2$$

$\pi/3$ phase \longleftrightarrow quarks?

Fix 1;

Rotate i : 6 choices;

Rotate j : 5 choices;

k fixed;

Rotate ℓ : 3 choices;

Done.

VECTORS II

$$\mathbf{x} = \begin{pmatrix} t \\ x \\ y \\ z \end{pmatrix} \longleftrightarrow \mathbf{X} = \begin{pmatrix} t+z & x-iy \\ x+iy & t-z \end{pmatrix}$$
$$\mathbf{X}^\dagger = \overline{\mathbf{X}}^T = \mathbf{X}$$

$$-\det(\mathbf{X}) = -t^2 + x^2 + y^2 + z^2$$

- {vectors in (3+1)-dimensional spacetime}
 \longleftrightarrow { 2×2 complex Hermitian matrices}
- determinant \longleftrightarrow (Lorentzian) inner product
- $\mathbf{X} = tI + x\sigma_x + y\sigma_y + z\sigma_z$ (Pauli matrices)

LORENTZ TRANSFORMATIONS

Exploit (local) isomorphism:

$$SO(3, 1) \approx SL(2, \mathbb{C})$$

$$\boldsymbol{x}' = \boldsymbol{\Lambda} \boldsymbol{x} \quad \longleftrightarrow \quad \boldsymbol{X}' = \boldsymbol{M} \boldsymbol{X} \boldsymbol{M}^\dagger$$

$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \alpha & -\sin \alpha & 0 \\ 0 & \sin \alpha & \cos \alpha & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \longleftrightarrow \begin{pmatrix} e^{-\frac{i\alpha}{2}} & 0 \\ 0 & e^{\frac{i\alpha}{2}} \end{pmatrix}$$

$$\det(\boldsymbol{M}) = 1 \quad \implies \quad \det \boldsymbol{X}' = \det \boldsymbol{X}$$

ROTATIONS

$$M_{zx} = \begin{pmatrix} \cos \frac{\alpha}{2} & -\sin \frac{\alpha}{2} \\ \sin \frac{\alpha}{2} & \cos \frac{\alpha}{2} \end{pmatrix}$$

$$M_{xy} = \begin{pmatrix} e^{-\frac{i\alpha}{2}} & 0 \\ 0 & e^{\frac{i\alpha}{2}} \end{pmatrix}$$

$$M_{yz} = \begin{pmatrix} \cos \frac{\alpha}{2} & -i \sin \frac{\alpha}{2} \\ -i \sin \frac{\alpha}{2} & \cos \frac{\alpha}{2} \end{pmatrix}$$

$$i \longrightarrow j, k, \dots, l$$

ROTATIONS

Still missing: rotations in $\text{Im } \mathbb{K}$

$$\text{III: } \mathbf{M} = e^{k\theta} \mathbf{I}$$

⊙: flips!

$$\mathbf{X}' = \mathbf{M}_2(\mathbf{M}_1 \mathbf{X} \mathbf{M}_1^\dagger) \mathbf{M}_2^\dagger$$

$$\mathbf{M}_1 = i\mathbf{I} \quad \mathbf{M}_2 = (i \cos \theta + j \sin \theta) \mathbf{I}$$

WHICH DIMENSIONS?

$\mathbb{K} = \mathbb{R}, \mathbb{C}, \mathbb{H}, \mathbb{O} \longmapsto$

$$\mathbf{X} = \begin{pmatrix} p & \bar{a} \\ a & m \end{pmatrix} \quad (p, m \in \mathbb{R}; a \in \mathbb{K})$$

$\dim \mathbb{K} + 2 = 3, 4, 6, 10$ spacetime dimensions

supersymmetry

$$SO(5, 1) \approx SL(2, \mathbb{H})$$

$$SO(9, 1) \approx SL(2, \mathbb{O})$$

PENROSE SPINORS

$$v = \begin{pmatrix} c \\ \bar{b} \end{pmatrix}$$

$$vv^\dagger = \begin{pmatrix} |c|^2 & cb \\ \bar{b}\bar{c} & |b|^2 \end{pmatrix}$$

$$\det(vv^\dagger) = 0$$

$$(\text{spinor})^2 = \text{null vector}$$

Lorentz transformation:

$$v' = Mv$$

$$M(vv^\dagger)M^\dagger = (Mv)(Mv)^\dagger$$

compatibility

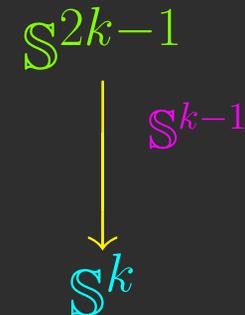
HOPF FIBRATIONS

$$v \in \mathbb{K}^2$$

$$v^\dagger v = 1 \implies v \in \mathbb{S}^{2k-1}$$

$$\text{tr}(vv^\dagger) = v^\dagger v = 1 \implies t = \text{const}$$

$$\det(vv^\dagger) = 0 \implies vv^\dagger \in \mathbb{S}^k$$



phase freedom:

$$\begin{pmatrix} p \\ q \end{pmatrix} \longmapsto \begin{pmatrix} \frac{p\bar{q}}{|q|} \\ |q| \end{pmatrix} e^{s\theta}$$

$$(s^2 = -1)$$

$$\mathbb{O}P^1 = \{(p, q) \sim (pq^{-1}\chi, \chi)\}$$

$$\mathbb{O}P^1 = \{X^2 = X, \text{tr } X = 1\}$$

WEYL EQUATION

- Massless, relativistic, spin $\frac{1}{2}$
- Momentum space

$$\tilde{\mathbf{P}}\psi = p^\mu \tilde{\boldsymbol{\sigma}}_\mu \psi = 0$$

$$\tilde{\mathbf{P}} = \mathbf{P} - (\text{tr } \mathbf{P}) \mathbf{I}$$

Pauli Matrices:

$$\begin{aligned} \boldsymbol{\sigma}_t &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} & \boldsymbol{\sigma}_z &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ \boldsymbol{\sigma}_x &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} & \boldsymbol{\sigma}_\ell &= \begin{pmatrix} 0 & -\ell \\ \ell & 0 \end{pmatrix} & \boldsymbol{\sigma}_k &= \begin{pmatrix} 0 & -k \\ k & 0 \end{pmatrix} \end{aligned}$$

SOLUTION

Weyl equation:

(3 of 4 string equations!)

$$-\tilde{P}\psi = 0 \quad \implies \det(P) = 0$$

One solution: (P, θ complex)

$$P = \pm \theta \theta^\dagger$$

$$\widetilde{\theta \theta^\dagger} \theta = (\theta \theta^\dagger - \theta^\dagger \theta) \theta = \theta \theta^\dagger \theta - \theta^\dagger \theta \theta = 0$$

General solution: ($\xi \in \mathbb{O}$)

$$\psi = \theta \xi$$

P, ψ quaternionic

PHASE FREEDOM

$$\mathbf{P} = \begin{pmatrix} |c|^2 & cb \\ \bar{b}\bar{c} & |b|^2 \end{pmatrix} = \theta\theta^\dagger = (\theta e^{s\phi})(\theta e^{s\phi})^\dagger$$

$$\theta = \begin{pmatrix} |c| \\ \bar{b}\bar{c} \\ |c| \end{pmatrix}$$

$$\implies \tilde{\mathbf{P}} = \begin{pmatrix} -|b|^2 & cb \\ \bar{b}\bar{c} & -|c|^2 \end{pmatrix} \implies \psi = \theta\xi = \begin{pmatrix} |c| \\ \bar{b}\bar{c} \\ |c| \end{pmatrix} r e^{s\alpha}$$

- Phase freedom is supersymmetry.
- Solutions are quaternionic (only 2 directions: $\bar{b}\bar{c}, s$).

DIRAC EQUATION

Gamma matrices: (\mathbb{C} Weyl representation)

$$\gamma_\mu = \begin{pmatrix} 0 & \tilde{\sigma}_\mu \\ \sigma_\mu & 0 \end{pmatrix}$$

Dirac equation: (momentum space)

$$(p^\mu \gamma_\mu - m) \Psi = 0$$

Weyl equation: (\mathbb{H} Penrose spinors)

$$\Psi = \begin{pmatrix} \theta \\ \eta \end{pmatrix} \longleftrightarrow \psi = \eta + \sigma_k \theta$$
$$(p^\mu \tilde{\sigma}_\mu - m \tilde{\sigma}_k) \psi = 0$$

COMPARISON

4×4 complex:

$$0 = (\gamma_t \gamma_\mu p^\mu - m \gamma_t) \Psi$$

2×2 quaternionic:

$$\begin{aligned} 0 &= (p^t \sigma_t - p^\alpha \sigma_\alpha - m \sigma_k) \psi \\ &= -\tilde{P} \psi \end{aligned}$$

Isomorphism: $(\mathbb{H}^2 \approx \mathbb{C}^4)$

$$\begin{pmatrix} c - kb \\ d + ka \end{pmatrix} \longleftrightarrow \begin{pmatrix} a \\ b \\ c \\ d \end{pmatrix}$$

DIMENSIONAL REDUCTION

$$SO(3, 1) \approx SL(2, \mathbb{C}) \subset SL(2, \mathbb{O}) \approx SO(9, 1)$$

Projection: $(\mathbb{O} \rightarrow \mathbb{C})$

$$\pi(p) = \frac{1}{2}(p + \ell p \bar{\ell})$$

Determinant: $\det(P) = 0 \implies$

$$\det(\pi(P)) = m^2$$

Mass Term:

$$P = \pi(P) + m \sigma_k \quad \sigma_k = \begin{pmatrix} 0 & -k \\ k & 0 \end{pmatrix}$$

$$P = \begin{pmatrix} p^t + p^z & p^x - \ell p^y - km \\ p^x + \ell p^y + km & p^t - p^z \end{pmatrix}$$

SPIN

Finite rotation:

$$R_z = \begin{pmatrix} e^{l\frac{\theta}{2}} & 0 \\ 0 & e^{-l\frac{\theta}{2}} \end{pmatrix}$$

Infinitesimal rotation:

$$L_z = \left. \frac{dR_z}{d\theta} \right|_{\theta=0} = \frac{1}{2} \begin{pmatrix} l & 0 \\ 0 & -l \end{pmatrix}$$

Right self-adjoint operator:

$$\hat{L}_z \psi := (L_z \psi) \bar{\ell}$$

Right eigenvalue problem:

$$\hat{L}_z \psi = \psi \lambda$$

ANGULAR MOMENTUM REVISITED

$$L_x = \frac{1}{2} \begin{pmatrix} 0 & \ell \\ \ell & 0 \end{pmatrix} \quad L_y = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$
$$L_z = \frac{1}{2} \begin{pmatrix} \ell & 0 \\ 0 & -\ell \end{pmatrix} \quad \hat{L}_\mu \psi := -(L_\mu \psi) \ell$$

$$\psi = e_\uparrow = \begin{pmatrix} 1 \\ k \end{pmatrix} \implies$$

$$\hat{L}_z \psi = \psi \frac{1}{2} \quad \hat{L}_x \psi = -\psi \frac{k}{2} \quad \hat{L}_y \psi = -\psi \frac{k\ell}{2}$$

Simultaneous eigenvector!

(only 1 *real* eigenvalue)

LEPTONS

 ψ $P = \psi\psi^\dagger$

$$e_\uparrow = \begin{pmatrix} 1 \\ k \end{pmatrix}$$

$$e_\uparrow e_\uparrow^\dagger = \begin{pmatrix} 1 & -k \\ k & 1 \end{pmatrix}$$

$$e_\downarrow = \begin{pmatrix} -k \\ 1 \end{pmatrix}$$

$$e_\downarrow e_\downarrow^\dagger = \begin{pmatrix} 1 & -k \\ k & 1 \end{pmatrix}$$

$$\nu_z = \begin{pmatrix} 0 \\ k \end{pmatrix}$$

$$\nu_z \nu_z^\dagger = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\nu_{-z} = \begin{pmatrix} k \\ 0 \end{pmatrix}$$

$$\nu_{-z} \nu_{-z}^\dagger = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

How Many Quaternionic Spaces?

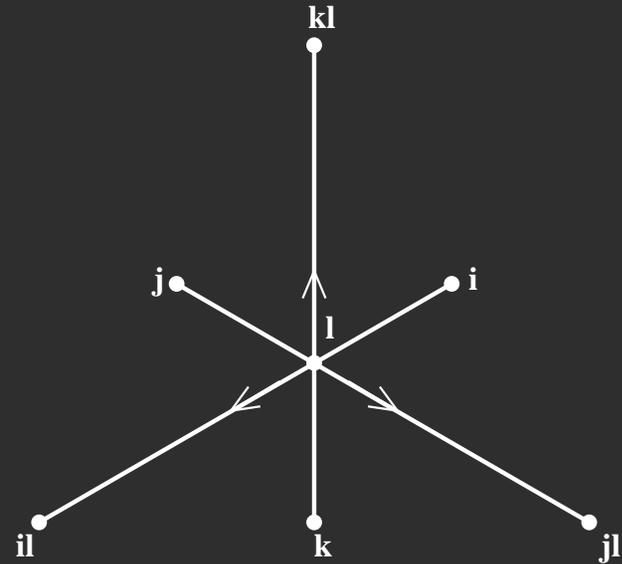
Dimensional Reduction:

$$\implies \ell \in \mathbb{H}$$

Orthogonality:

$$(\mathbb{H}_1 \cap \mathbb{H}_2 = \mathbb{C})$$

$$\longmapsto i, j, k$$



Answer: 3!

LEPTONS

$$e_{\uparrow} = \begin{pmatrix} 1 \\ k \end{pmatrix} \quad e_{\uparrow} e_{\uparrow}^{\dagger} = \begin{pmatrix} 1 & -k \\ k & 1 \end{pmatrix}$$

$$e_{\downarrow} = \begin{pmatrix} -k \\ 1 \end{pmatrix} \quad e_{\downarrow} e_{\downarrow}^{\dagger} = \begin{pmatrix} 1 & -k \\ k & 1 \end{pmatrix}$$

$$\nu_z = \begin{pmatrix} 0 \\ k \end{pmatrix} \quad \nu_z \nu_z^{\dagger} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\nu_{-z} = \begin{pmatrix} k \\ 0 \end{pmatrix} \quad \nu_{-z} \nu_{-z}^{\dagger} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$\emptyset_z = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \emptyset_z \emptyset_z^{\dagger} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

WHAT NEXT?

Have:

3 generations of leptons!

Neutrinos have just one helicity!

What about \mathcal{O}_z ?

Want:

- interactions
- quarks/color ($SU(3)$!)
- charge

JORDAN ALGEBRAS

$$\mathcal{X} = \begin{pmatrix} p & a & \bar{c} \\ \bar{a} & m & b \\ c & \bar{b} & n \end{pmatrix}$$

$$\mathcal{X} \circ \mathcal{Y} = \frac{1}{2}(\mathcal{X}\mathcal{Y} + \mathcal{Y}\mathcal{X})$$

$$\begin{aligned} \mathcal{X} * \mathcal{Y} &= \mathcal{X} \circ \mathcal{Y} - \frac{1}{2}(\mathcal{X} \operatorname{tr}(\mathcal{Y}) + \mathcal{Y} \operatorname{tr}(\mathcal{X})) \\ &\quad + \frac{1}{2}(\operatorname{tr}(\mathcal{X}) \operatorname{tr}(\mathcal{Y}) - \operatorname{tr}(\mathcal{X} \circ \mathcal{Y})) \mathcal{I} \end{aligned}$$

JORDAN ALGEBRAS

$u, v, w \in \mathbb{R}^3 \implies$

$$\begin{aligned}2uu^\dagger \circ vv^\dagger &= (u \cdot v)(uv^\dagger + vu^\dagger) \\ \text{tr}(uu^\dagger \circ vv^\dagger) &= (u \cdot v)^2 \\ 2uu^\dagger * vv^\dagger &= (u \times v)(u \times v)^\dagger\end{aligned}$$

$$\mathcal{X} \circ \mathcal{Y} = \frac{1}{2}(\mathcal{X}\mathcal{Y} + \mathcal{Y}\mathcal{X})$$

$$\begin{aligned}\mathcal{X} * \mathcal{Y} &= \mathcal{X} \circ \mathcal{Y} - \frac{1}{2}(\mathcal{X} \text{tr}(\mathcal{Y}) + \mathcal{Y} \text{tr}(\mathcal{X})) \\ &\quad + \frac{1}{2}(\text{tr}(\mathcal{X}) \text{tr}(\mathcal{Y}) - \text{tr}(\mathcal{X} \circ \mathcal{Y})) \mathcal{I}\end{aligned}$$

JORDAN ALGEBRAS

Exceptional quantum mechanics:

(Jordan, von Neumann, Wigner)

$$(\mathcal{X} \circ \mathcal{Y}) \circ \mathcal{X}^2 = \mathcal{X} \circ (\mathcal{Y} \circ \mathcal{X}^2)$$

$$\mathcal{X} \circ \mathcal{Y} = \frac{1}{2}(\mathcal{X}\mathcal{Y} + \mathcal{Y}\mathcal{X})$$

$$\begin{aligned} \mathcal{X} * \mathcal{Y} &= \mathcal{X} \circ \mathcal{Y} - \frac{1}{2}(\mathcal{X} \operatorname{tr}(\mathcal{Y}) + \mathcal{Y} \operatorname{tr}(\mathcal{X})) \\ &\quad + \frac{1}{2}(\operatorname{tr}(\mathcal{X}) \operatorname{tr}(\mathcal{Y}) - \operatorname{tr}(\mathcal{X} \circ \mathcal{Y})) \mathcal{I} \end{aligned}$$

MORE ROTATION GROUPS

$$\begin{aligned}\sqrt{3}t &= p + m + n \\ 2\sqrt{3}w &= p + m - 2n \\ 2z &= p - m\end{aligned}$$

***SO*(27):**

$$\text{tr}(\mathcal{X} \circ \mathcal{X}) = 2(|a|^2 + |b|^2 + |c|^2 + w^2 + z^2) + t^2$$

***SO*(26, 1):**

$$-\text{tr}(\mathcal{X} * \mathcal{X}) = |a|^2 + |b|^2 + |c|^2 + w^2 + z^2 - t^2$$

EXCEPTIONAL GROUPS

F_4 : “ $SU(3, \mathbb{O})$ ”

$$(\mathcal{M}\mathcal{X}\mathcal{M}^\dagger) \circ (\mathcal{M}\mathcal{Y}\mathcal{M}^\dagger) = \mathcal{M}(\mathcal{X} \circ \mathcal{Y})\mathcal{M}^\dagger$$

E_6 : “ $SL(3, \mathbb{O})$ ”

$$\det \mathcal{X} = \frac{1}{3} \operatorname{tr} ((\mathcal{X} * \mathcal{X}) \circ \mathcal{X})$$

$$SO(3, 1) \times U(1) \times SU(2) \times SU(3) \subset E_6$$

SUPERSPINORS

$$\mathcal{X} = \begin{pmatrix} \mathbf{X} & \theta \\ \theta^\dagger & n \end{pmatrix} \quad \mathcal{M}\mathcal{X}\mathcal{M}^\dagger = \begin{pmatrix} \mathbf{M}\mathbf{X}\mathbf{M}^\dagger & \mathbf{M}\theta \\ \theta^\dagger\mathbf{M}^\dagger & n \end{pmatrix}$$
$$\mathcal{M} = \begin{pmatrix} \mathbf{M} & 0 \\ 0 & 1 \end{pmatrix}$$

Cayley/Moufang plane:

$$\begin{aligned} \mathbb{O}P^2 &= \{\mathcal{X}^2 = \mathcal{X}, \operatorname{tr} \mathcal{X} = 1\} \\ &= \{\mathcal{X} * \mathcal{X} = 0, \operatorname{tr} \mathcal{X} = 1\} \\ &= \{\mathcal{X} = \psi\psi^\dagger, \psi^\dagger\psi = 1\} \end{aligned}$$

$$(\psi \in \mathbb{H}^3)$$

DIRAC EQUATION II

$$\mathcal{P} = \begin{pmatrix} \mathbf{P} & \theta\xi \\ \bar{\xi}\theta^\dagger & |\xi|^2 \end{pmatrix}$$

$$\mathcal{P} * \mathcal{P} = 0 \implies \tilde{\mathbf{P}}\theta = 0$$
$$\mathcal{P} = \psi\psi^\dagger$$

quaternionic!

Furthermore:

$$\mathcal{X}^\dagger = \mathcal{X} \implies \mathcal{X} = \sum_{n=1}^3 \psi_n \psi_n^\dagger$$

leptons, mesons, baryons?

Life is complex.

It has real and imaginary parts.

Life is octonionic...

THE END

Start

Close

Exit