# The Geometry of Relativity

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# Differential Geometry

#### Definition

A *topological manifold* is a second countable Housdorff space that is locally homeomorphic to Euclidean space. A *differentiable manifold* is a topological manifold equipped with an equivalence class of atlases whose transition maps are differentiable.

General Relativity  $\neq$  Differential Geometry

What math is needed for GR??

# Background

- Differential geometry course: Rick Schoen
- GR reading course: MTW
- GR course: Sachs-Wu
- designed and taught undergrad math course in GR: Schutz, d'Inverno, Wald, Taylor–Wheeler, Hartle
- designed and taught undergrad physics course in SR
- NSF-funded curricular work (math and physics) since 1996
- national expert in teaching 2nd-year calculus http://blogs.ams.org/matheducation

geometer, relativist, curriculum developer, education researcher Mathematics, Physics, PER, RUME

## Math vs. Physics

#### My math colleagues think I'm a physicsist.

My physics colleagues know better.

# Books



The Geometry of Special Relativity Tevian Dray A K Peters/CRC Press 2012 ISBN: 978-1-4665-1047-0 http://physics.oregonstate.edu/coursewikis/GSR



#### Differential Forms and the Geometry of General Relativity *Tevian Dray* A K Peters/CRC Press 2014 ISBN: 978-1-4665-1000-5 http://physics.oregonstate.edu/coursewikis/GDF http://physics.oregonstate.edu/coursewikis/GGR

Hyperbolic Trigonometry Applications

# Trigonometry



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# Length Contraction



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#### Paradoxes

A 20 foot pole is moving towards a 10 foot barn fast enough that the pole appears to be only 10 feet long. As soon as both ends of the pole are in the barn, slam the doors. How can a 20 foot pole fit into a 10 foot barn?



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#### **Relativistic** Mechanics

A pion of (rest) mass m and (relativistic) momentum  $p = \frac{3}{4}mc$ decays into 2 (massless) photons. One photon travels in the same direction as the original pion, and the other travels in the opposite direction. Find the energy of each photon.  $[E_1 = mc^2, E_2 = \frac{1}{4}mc^2]$ 





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### Addition Formulas

 $v = c \tanh \beta$ 

#### **Einstein Addition Formula:**

 $\tanh(\alpha + \beta) = \frac{\tanh \alpha + \tanh \beta}{1 + \tanh \alpha \tanh \beta} \qquad ("v + w" = \frac{v + w}{1 + vw/c^2})$ 

**Conservation of Energy-Momentum:** 

 $p = mc \sinh \alpha$  $E = mc^2 \cosh \alpha$ 

Moving Capacitor:

$$E'^{y} = C \cosh(\alpha + \beta) = E^{y} \cosh\beta - cB^{z} \sinh\beta$$
$$cB'^{z} = -C \sinh(\alpha + \beta) = cB^{z} \cosh\beta - E^{y} \sinh\beta$$

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## 3d spacetime diagrams



 $(v \Delta t)^2 - (c \Delta t)^2 = -(c \Delta t')^2$ 

http://relativity.geometryof.org/GSR/book/updates/3d

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# The Geometry of General Relativity





Doppler effect (SR) Cosmological redshift (GR)

Asymptotic structure

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# Line Elements



 $dr^2 + r^2 d\phi^2$   $d\theta^2 + \sin^2 \theta d\phi^2$   $d\beta^2 + \sinh^2 \beta d\phi^2$ 



# Vector Calculus





 $d\vec{r} = dx\,\hat{\imath} + dy\,\hat{\jmath} = dr\,\hat{r} + r\,d\phi\,\hat{\phi}$ 

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# Differential Forms in a Nutshell $(\mathbb{R}^3)$

Differential forms are integrands:  $(*^2 = 1)$ 

 $f = f \qquad (0-\text{form})$   $F = \vec{F} \cdot d\vec{r} \qquad (1-\text{form})$   $*F = \vec{F} \cdot d\vec{A} \qquad (2-\text{form})$   $*f = f \, dV \qquad (3-\text{form})$ 

Exterior derivative:  $(d^2 = 0)$ 

$$df = \vec{\nabla} f \cdot d\vec{r}$$
$$dF = \vec{\nabla} \times \vec{F} \cdot d\vec{A}$$
$$d*F = \vec{\nabla} \cdot \vec{F} \, dV$$
$$d*f = 0$$

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#### Maxwell's Equations

$ec  abla \cdot ec {oldsymbol E} = 4 \pi  ho$	
$ec{ abla}\cdotec{oldsymbol{B}}=0$	
$ec{ abla}  imes ec{m{ extbf{ ex} extbf{ extbf{ extbf{ extbf{ extbf{ extbf{ extbf{$	dE
$ec{ abla}  imes ec{m{B}} - \dot{ec{m{E}}} = 4\pi ec{m{J}}$	dB
$ec{ abla} \cdot ec{oldsymbol{J}} + \dot{ ho} = 0$	k
$ec{m{B}}=ec{ abla} imesec{m{A}}$	E
$ec{m{ extbf{E}}} = -ec{ abla} \Phi - \dot{ec{m{ extbf{A}}}}$	Ε

 $d * E = 4\pi * \rho$ d \* B = 0 $\dot{F} + *\dot{B} = 0$  $3 - *\dot{F} = 4\pi * I$  $*d*J + \dot{\rho} = 0$ 3 = \*dA $F = -d\Phi - \dot{A}$ 

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### Maxwell's Equations II

 $\vec{\nabla} \cdot \vec{E} = 4\pi\rho$  $\vec{\nabla} \cdot \vec{B} = 0$  $\vec{\nabla} \times \vec{\pmb{E}} + \dot{\vec{\pmb{B}}} = 0$  $\vec{\nabla} \times \vec{B} - \dot{\vec{E}} = 4\pi \vec{I}$  $\vec{\nabla} \cdot \vec{J} + \dot{
ho} = 0$  $\vec{B} = \vec{\nabla} \times \vec{A}$  $\vec{E} = -\vec{\nabla}\Phi - \dot{\vec{A}}$ 

$$F = \hat{E} \wedge dt + \hat{*}\hat{B}$$
$$*F = \hat{B} \wedge dt - \hat{*}\hat{E}$$
$$A = \hat{A} - \Phi dt$$

$$F = dA$$
  
 $d*F = 4\pi *J$ 

$$\implies dF = 0$$
$$d*J = 0$$

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## Geodesic Equation

Orthonormal basis:

$$d\vec{r} = \sigma^i \,\hat{\mathbf{e}}_i \quad (\Longrightarrow ds^2 = d\vec{r} \cdot d\vec{r})$$

Connection:

$$egin{aligned} & \omega_{ij} = \hat{\mathbf{e}}_i \cdot d\hat{\mathbf{e}}_j \ d\sigma^i + \omega^i{}_j \wedge \sigma^j = 0 \ & \omega_{ij} + \omega_{ji} = 0 \end{aligned}$$

Geodesics:

$$\vec{v} d\lambda = d\vec{r}$$
  
 $\dot{\vec{v}} = 0$ 

Symmetry:

 $d\vec{X} \cdot d\vec{r} = 0$  $\implies \vec{X} \cdot \vec{v} = \text{const}$ 

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# Example: Polar Coordinates

Symmetry:

$$ds^{2} = dr^{2} + r^{2} d\phi^{2}$$
$$\implies d\vec{r} = dr \hat{r} + r d\phi \hat{\phi}$$
$$\implies r \hat{\phi} \text{ is Killing}$$

**Idea**: 
$$df = \vec{\nabla} f \cdot d\vec{r} \implies r \hat{\phi} \cdot \vec{\nabla} f = \frac{\partial f}{\partial \phi} \implies r \hat{\phi} = \frac{\partial}{\partial \phi}$$
  
**Check**:  $d(r \hat{\phi}) = dr \hat{\phi} + r d\hat{\phi} = dr \hat{\phi} - r d\phi \hat{r} \perp d\vec{r}$ 

**Geodesic Equation:** 

$$\vec{\nu} = \dot{r}\,\hat{r} + r\,\dot{\phi}\,\hat{\phi} \Longrightarrow r\,\hat{\phi}\cdot\vec{v} = r^2\,\dot{\phi} = \ell$$
$$\implies 1 = \dot{r}^2 + r^2\,\dot{\phi}^2 = \dot{r}^2 + \frac{\ell^2}{r^2}$$

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#### Einstein's Equation

Curvature:

$$\Omega^{i}_{j} = \boldsymbol{d}\omega^{i}_{j} + \omega^{i}_{k} \wedge \omega^{k}_{j}$$

Einstein tensor:

$$\gamma^{i} = -\frac{1}{2} \Omega_{jk} \wedge *(\sigma^{i} \wedge \sigma^{j} \wedge \sigma^{k})$$
$$G^{i} = *\gamma^{i} = G^{i}{}_{j} \sigma^{j}$$
$$\vec{G} = G^{i} \hat{e}_{i} = G^{i}{}_{j} \sigma^{j} \hat{e}_{i}$$
$$\Longrightarrow d * \vec{G} = 0$$

Field equation:

$$\vec{\boldsymbol{G}} + \Lambda \, d\, \vec{\boldsymbol{r}} = 8\pi\, \vec{\boldsymbol{T}}$$

(vector valued 1-forms, not tensors)

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# Stress-Energy Tensor

$$d\vec{r} = \sigma^a \,\hat{e}_a$$

Vector-valued 1-form:

$$\vec{T} = T^a{}_b \sigma^b \, \hat{e}_a$$

3-form:

$$\tau^a = *T^a$$

**Conservation:** 

$$d\left(\tau^{a}\,\hat{\boldsymbol{e}}_{a}\right)=0$$

$$*d*\vec{T}=0$$

#### What about Tensors?

What tensors are needed to do GR?

Metric? Use  $d\vec{r}$ ! (vector-valued 1-form!) Curvature? Riemann tensor is really a 2-form. (Cartan!) Ricci? Einstein? Stress-Energy? Vector-valued 1-forms!

> ∃ only 1 essential symmetric tensor in GR! Killing eq:  $d\vec{X} \cdot d\vec{r} = 0$ Students understand line elements...

$$ds^2 = d\vec{r} \cdot d\vec{r}$$

# Topic Order

#### Examples First!

- Schwarzschild geometry can be analyzed using vector calculus.
- Rain coordinates! (Painlevé-Gullstrand; freely falling)

#### Geodesics:

- EBH: Principle of Extremal Aging
- Hartle: variational principle (Lagrangian mechanics?)
- TD: "differential forms without differential forms"

# Choices

#### Language:

- Mathematicians: invariant objects (no indices)
- Physicists: components (indices)
- Relativists: abstract index notation ("indices without indices")
- Cartan: curvature without tensors

 $\therefore$  use differential forms?

#### **Coordinates:**

- Mathematicians: coordinate basis (usually)
- Physicists: calculate in coordinates; interpret in orthonormal basis
- Equivalence problem: 79310 coordinate components reduce to 8690

 $\therefore$  use orthonormal frames?  $(d\vec{r}?)$ 

# SUMMARY



http://relativity.geometryof.org/GSR http://relativity.geometryof.org/GDF http://relativity.geometryof.org/GGR



- Special relativity is hyperbolic trigonometry!
- General relativity can be described without tensors!
- BUT: Need vector-valued differential forms...

#### THE END