

Complex Numbers

- Tell me something you know about complex numbers

$$z = x + iy = r e^{i\phi}$$

$$\bar{z} = x - iy = r e^{-i\phi}$$

\mathbb{C} is a vector space

← what dimension?!

normed: $|z|^2 = z\bar{z} = x^2 + y^2 = r^2$

algebra: \exists multiplication

non degenerate: $|z|=0 \Rightarrow z=0$

← careful!
correct def for
positive-definite
case only

composition: $|ab| = |a||b|$

division: $z \neq 0 \Rightarrow z^{-1} = \frac{\bar{z}}{|z|^2}$

• Arms

$$1, i, -i$$

$$3i$$

$$e^{i\pi}$$

$$z \text{ (your choice)}$$

$$e^{i\frac{\pi}{2}} z$$

$$iz$$

$$\bar{z}$$

Rotations

- Tell me something you know about rotations in 2 dimensions

composition (multiplication)
can undo (inverses)

↳ group!

Representations

Ⓘ angle $\phi \in \mathbb{S}^1$

Ⓙ matrix $M(\phi) \in \mathbb{R}^{2 \times 2}$

$$M(\phi) = \begin{pmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{pmatrix}$$

Ⓚ complex number $w(\phi) \in \mathbb{C}$

$$w(\phi) = e^{-i\phi}$$

Conditions : $v \in \mathbb{R}^2$, $|v|^2 = v^T v$, $|Mv| = |v|$
 $\Rightarrow \underline{M^T M = \mathbb{I}}$ ($\& \det M = 1$)

$$|wz| = |z| \Rightarrow \underline{|w| = 1}$$

Isomorphism

$$e^{-i\phi} = \cos \phi + i \sin \phi \longleftrightarrow M(\phi) = \cos \phi \mathbb{I} + \sin \phi \Omega$$

$$x + iy \longleftrightarrow \begin{pmatrix} x & -y \\ y & x \end{pmatrix}$$

Exponentiation

- what does e^{kx} mean?

$$e^{kx} = 1 + kx + \frac{(kx)^2}{2} + \dots$$

unique soln of $\frac{df}{dx} = kF$, $F(0) = 1$

But $M'(\phi) = \Omega M$ & $\Omega = M'(0) = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$

$w'(\phi) = iw$ & $i = w'(0)$

\Downarrow
 $\Omega^2 = -\mathbb{I}!$

$\Rightarrow w = e^{i\phi}$ ✓
 $M = e^{\Omega\phi}$

Orthogonal Groups

$$v \in \mathbb{R}^n, |v|^2 = v^T v, M \in \mathbb{R}^{2 \times 2}$$

want: $|Mv| = |v|$

$$\Rightarrow (Mv)^T (Mv) = v^T v$$

" $v^T M^T M v$

$$\Rightarrow \boxed{M^T M = \mathbb{I}} \quad \leftarrow O(n)$$

$$\boxed{SO(n) = \{M \in \mathbb{R}^{n \times n}; M^T M = \mathbb{I} \ \& \ \det M = 1\}}$$

group! manifold! \mapsto Lie group

Derivatives

Suppose $M(\phi) \in SO(n), M(0) = \mathbb{I}$

$$\Rightarrow M(\phi)^T M(\phi) = \mathbb{I}$$

$$\Rightarrow M'(\phi)^T M(\phi) + M(\phi)^T M'(\phi) = 0$$

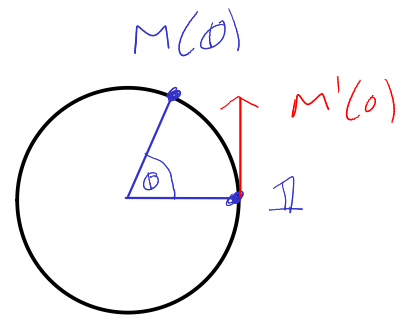
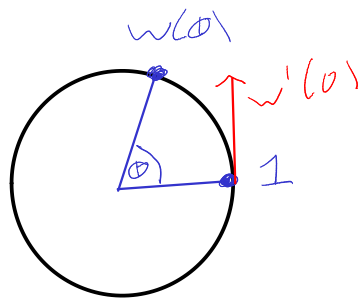
$$\Rightarrow M'(0)^T \mathbb{I} + \mathbb{I}^T M'(0) = 0$$

$$\mapsto \boxed{so(n) = \{A \in \mathbb{R}^{n \times n}; A^T + A = 0\}}$$

vector space! algebra! \mapsto Lie algebra

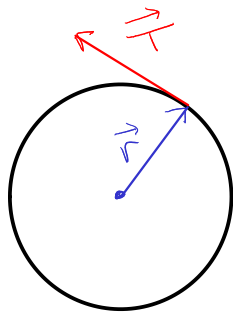
closed under commutator

SO(2)



Idea: Lie group = group with smooth parameters
 = group + manifold
Lie algebra = tangent vectors at identity

Vector fields



$$\begin{aligned} \vec{r} &= x\hat{x} + y\hat{y} \\ \vec{r}(\theta) &= \cos\theta\hat{x} + \sin\theta\hat{y} \\ \Rightarrow \vec{T}(\theta) &= \frac{d\vec{r}}{d\theta} = -\sin\theta\hat{x} + \cos\theta\hat{y} \\ &= -y\hat{x} + x\hat{y} \end{aligned}$$

Differential geometry

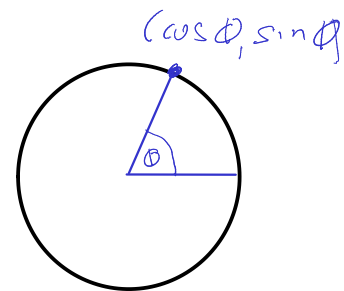
$$\vec{v}(f) = \vec{v} \cdot \nabla f$$

$$\therefore \vec{T}(f) = -y \frac{\partial f}{\partial x} + x \frac{\partial f}{\partial y}$$

$$\Rightarrow \vec{T} \leftrightarrow -y \partial_x + x \partial_y = \partial_\theta$$

vector fields are directional derivative operators

$$SO(2) \leftrightarrow M(\theta) = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \\ = e^{-\theta J}$$



$$J = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, \quad J^2 = -\mathbb{I} \\ = M'(0)$$

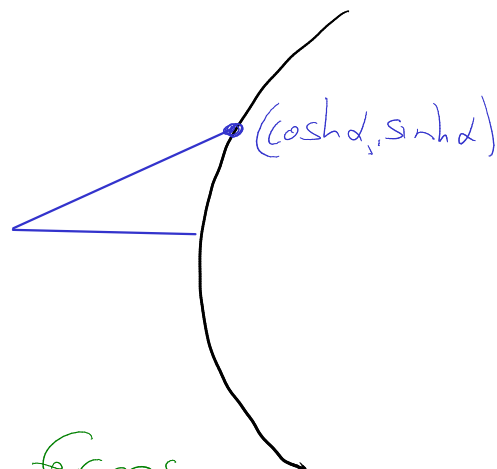
$$SO(1,1) \leftrightarrow P(\alpha) = \begin{pmatrix} \cosh \alpha & \sinh \alpha \\ \sinh \alpha & \cosh \alpha \end{pmatrix} \\ P'(0) = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \Lambda, \quad \Lambda^2 = +\mathbb{I} \\ \Rightarrow P(\alpha) = e^{-\Lambda \alpha}$$

Signature

M preserves $v^T v = v^T \mathbb{I} v$

P preserves $v^T \eta v$

$$\text{with } \eta = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$$



" $so(1,1)$ and $so(2)$ are real forms of (complex) $so(2)$ "

split complex numbers

"hyperbolic numbers"

$$L^2 = +1$$

$$\mathbb{C}' = \mathbb{R} \oplus \mathbb{R}L$$

compare: $\mathbb{C} = \mathbb{R} \oplus \mathbb{R}i$

$$\Rightarrow e^{L\alpha} = \cosh \alpha + L \sinh \alpha \leftrightarrow P(\alpha)$$

Rotations in 3d

- Tell me something you know about rotations in 3d

not commutative
group

Representations

$$\textcircled{\text{I}} R_z(\alpha) = \begin{pmatrix} \cos \alpha & -\sin \alpha & 0 \\ \sin \alpha & \cos \alpha & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$R_x(\alpha) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \alpha & -\sin \alpha \\ 0 & \sin \alpha & \cos \alpha \end{pmatrix}, \quad R_y(\alpha) = \begin{pmatrix} \cos \alpha & 0 & \sin \alpha \\ 0 & 1 & 0 \\ -\sin \alpha & 0 & \cos \alpha \end{pmatrix}$$

→ Euler angles

$$R(\theta, \phi, \psi) = R_z(\phi) R_y(\theta) R_z(\psi) \\ = [\text{mess}]$$

Derivatives

$$\dot{r}_m = \dot{R}_m = R'_m(0)$$

$$\textcircled{\text{II}} \begin{aligned} \vec{Z} &= x \partial_y - y \partial_x \\ \vec{X} &= y \partial_z - z \partial_y \\ \vec{Y} &= z \partial_x - x \partial_z \end{aligned}$$

Topology of $SO(3)$

Idea: Euler angles $\theta, \phi \leftrightarrow S^2$
 $\psi \leftrightarrow S^1$

$$\therefore SO(3) \approx S^2 \times S^1$$

Not quite: correct locally but
not globally

\mapsto Hopf bundle:

$$\begin{array}{c} S^3 \\ \downarrow S^1 \\ S^2 \end{array}$$

so perhaps $SO(3) \approx S^3$

Not quite: rotations around
antipodal points equivalent

$\therefore SO(3)$ is topologically S^3/\mathbb{Z}_2

so

$$SO(3) \cong \mathbb{R}P^3$$

Commutators

Goal: need product for Lie algebra

Matrices: $[A, B] = AB - BA$

$$\Rightarrow [r_x, r_y] = r_z, [r_y, r_z] = r_x, [r_z, r_x] = r_y$$

\Rightarrow so(3) closed under commutation ✓

Vector Fields

$$[\mathbb{X}, \mathbb{Y}](f) = \mathbb{X}(\mathbb{Y}(f)) - \mathbb{Y}(\mathbb{X}(f))$$

$$\Rightarrow [\mathbb{X}, \mathbb{Y}] = -\mathbb{Z}$$

↑ conventional sign,
can be eliminated
by change of basis

Properties: bilinear
antisymmetric
(Jacobi identity)

Check: $A^T = -A, B^T = -B$
(exercise) $\Rightarrow [A, B]^T = -[A, B]$

\Rightarrow so(n) closes
under commutation

Definitions

A Lie group is a group that is also a smooth manifold (surface), on which the group operations are smooth

↑
multiplication
and inverses

A Lie algebra is a vector space that is closed under a bilinear operation $[,]$ with:

$$[Y, X] = -[X, Y] \quad \text{antisymmetry}$$

$$[X, [Y, Z]] + [Y, [Z, X]] + [Z, [X, Y]] = 0$$

Jacobi identity
(automatic for matrices!)

Fact: elements of Lie group near \mathbb{I} look like $M(\alpha) = e^{A\alpha}$ for A in Lie algebra
matrix exponentiation if matrix Lie group

$\Rightarrow M(\alpha)$ is a 1-parameter family:

$$\begin{aligned} M(0) &= \mathbb{I} \\ M(\alpha + \beta) &= M(\alpha)M(\beta) \\ M(-\alpha) &= M(\alpha)^{-1} \end{aligned}$$

Vector Fields

Day 03

Recall

$$\vec{\nabla}(f) = \vec{v} \cdot \vec{\nabla} f$$

Ex:

$$\hat{x}(f) = \frac{\partial f}{\partial x}$$

so $\hat{x} \leftrightarrow \frac{\partial}{\partial x} = \partial_x$

$$(x\hat{y} - y\hat{x})(f) = x\frac{\partial f}{\partial y} - y\frac{\partial f}{\partial x} = \frac{\partial f}{\partial \theta}$$

Claim

$[\vec{v}, \vec{w}]$ is a vector field

Idea

$$v_x \hat{x} (v_y \hat{y}(f)) - v_y \hat{y} (v_x \hat{x}(f))$$

$$= v_x \frac{\partial}{\partial x} (v_y \frac{\partial f}{\partial y}) - v_y \frac{\partial}{\partial y} (v_x \frac{\partial f}{\partial x})$$

$$= v_x \frac{\partial v_y}{\partial x} \frac{\partial f}{\partial y} + v_x v_y \frac{\partial^2 f}{\partial x \partial y} - v_y \frac{\partial v_x}{\partial y} \frac{\partial f}{\partial x} - v_y v_x \frac{\partial^2 f}{\partial y \partial x}$$

$$= (v_x \frac{\partial v_y}{\partial x} - v_y \frac{\partial v_x}{\partial y}) f$$

$$\Rightarrow [v_x \hat{x}, v_y \hat{y}] = v_x \frac{\partial v_y}{\partial x} \hat{y} - v_y \frac{\partial v_x}{\partial y} \hat{x}$$

Rotations

$$\vec{Z} = x\partial_y - y\partial_x$$

$$\vec{X} = y\partial_z - z\partial_y$$

$$\vec{Y} = z\partial_x - x\partial_z$$

$$\Rightarrow [\vec{X}, \vec{Y}] = -\vec{Z} \dots$$

Jacobi Identity

Recall :
$$[X, [Y, Z]] + [Y, [Z, X]] + [Z, [X, Y]] = 0$$

Rearrange :

$$\textcircled{1} \quad [Z, [X, Y]] = [[Z, X], Y] + [X, [Z, Y]]$$

product rule!
(Lie bracket is a derivation)

$$\textcircled{2} \quad [[X, Y], Z] = [X, [Y, Z]] - [Y, [X, Z]]$$

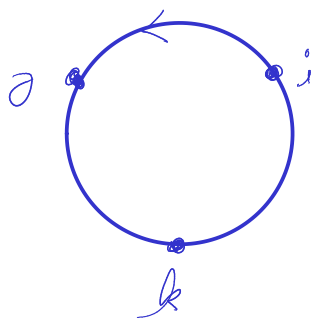
homomorphism!
(action of commutator is commutator of actions)

Quaternions

\mathbb{H} for Hamilton

$$q \in \mathbb{H} = \mathbb{R} \oplus \mathbb{R}i \oplus \mathbb{R}j \oplus \mathbb{R}k$$

$$i^2 = j^2 = k^2 = -1$$



$$ij = k = -ji \dots$$

- Show pictures of Broome Bridge (1853)
- $q \leftrightarrow$ Hamilton late 1800s
 $\vec{v} \leftrightarrow$ Gibbs
- Modern applications:
robotics, aeronautics, computer graphics

Correspondence

$$\vec{v} = x\hat{x} + y\hat{y} + z\hat{z} \leftrightarrow v = xi + yj + zk \in \text{Im } \mathbb{H}$$
$$"- \vec{v} \cdot \vec{w} + \vec{v} \times \vec{w}" \quad \leftrightarrow \quad vw$$

quaternionic product incorporates
both dot & cross products

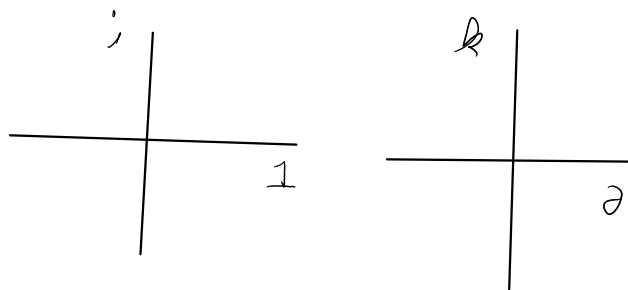
Quaternionic Rotations

$$q \in \mathbb{H} \Rightarrow q = a + bj \quad \text{with } \underline{a, b \in \mathbb{C}}$$

SWBQ: Find a, b in terms of w, x, y, z
with $q = w + ix + jy + kz$

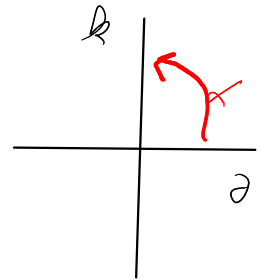
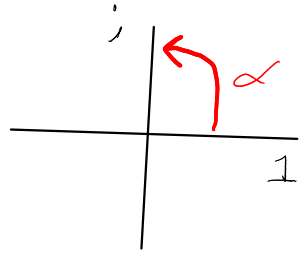
KA: Working with a partner, represent

1
k
i + j
q (your choice)
 $iq, e^{i\alpha} q$
 $qi, qe^{i\alpha}$

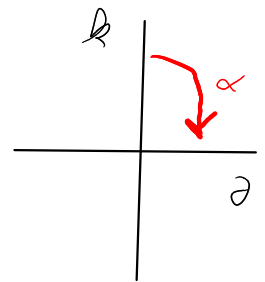
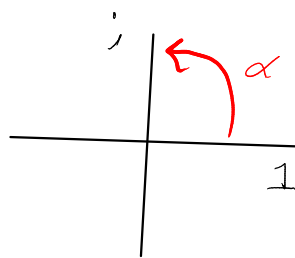


Idea: $iq = i(a+bi) = ia + ib^2$
 $qi = (a+bi)i = ia - b^2$

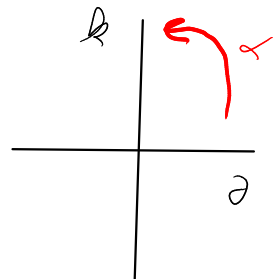
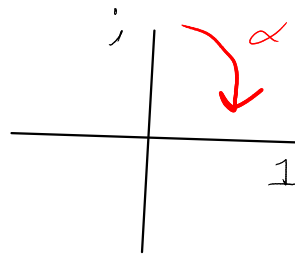
$\therefore e^{i\alpha} q :$



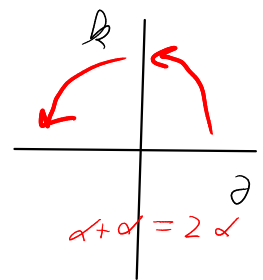
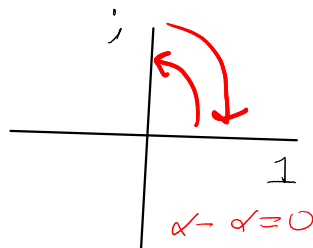
$q e^{i\alpha} :$



$q e^{-i\alpha} :$



$e^{i\alpha} q e^{-i\alpha} :$



$\therefore e^{\frac{i\alpha}{2}} q e^{-\frac{i\alpha}{2}} \leftrightarrow$ rotation by α in ab plane
 "about i axis in $\text{Im } \mathbb{H}$ "

Can replace i by any unit element $u \in \text{Im } \mathbb{H}$, $|u|=1$

$\therefore e^{\frac{u\alpha}{2}} q e^{-\frac{u\alpha}{2}} \leftrightarrow$ rotation in $\text{Im } \mathbb{H}$
 by α about u axis

Homomorphism

Given: Lie algebra \mathfrak{g} with Lie bracket

$$[\bar{x}, \bar{y}]_{\mathfrak{g}}$$

$$\Rightarrow \mathfrak{g} \longrightarrow \text{Lin}(\mathfrak{g}, \mathfrak{g}) \longleftrightarrow \mathbb{R}^{n \times n}$$

(n = dim \mathfrak{g})

$$\bar{x} \xrightarrow{L} L_{\bar{x}}$$

where $L_{\bar{x}}: \mathfrak{g} \rightarrow \mathfrak{g}$

$$\bar{y} \mapsto [\bar{x}, \bar{y}]$$

$$P, Q \in \text{Lin}(\mathfrak{g}, \mathfrak{g}) \Rightarrow$$

$[P, Q]_L \in \text{Lin}(\mathfrak{g}, \mathfrak{g})$ where

$$[P, Q]_L : z \mapsto P(Q(z)) - Q(P(z))$$

Recall: Jacobi identity \Rightarrow

$$[[\bar{x}, \bar{y}], \bar{z}] = [\bar{x}, [\bar{y}, \bar{z}]] - [\bar{y}, [\bar{x}, \bar{z}]]$$

$$\Leftrightarrow L_{[\bar{x}, \bar{y}]} = [L_{\bar{x}}, L_{\bar{y}}]_L$$

$\therefore L$ is a Lie algebra homomorphism

Unitary Matrices

Recall: $SO(n) = \{M \in \mathbb{R}^{n \times n}; M^T M = \mathbb{1}, \det M = 1\}$
preserves $|v|^2 = v^T v$

Generalization: $SU(n) = \{M \in \mathbb{C}^{n \times n}; M^T M = \mathbb{1}, \det M = 1\}$
preserves $|v|^2 = v^T v$

Ex: $SU(2)$

$$S_y(\alpha) = \begin{pmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{pmatrix}$$

$$S_x(\alpha) = \begin{pmatrix} \cos \alpha & -i \sin \alpha \\ i \sin \alpha & \cos \alpha \end{pmatrix}$$

$$S_z(\alpha) = \begin{pmatrix} e^{-i\alpha} & 0 \\ 0 & e^{i\alpha} \end{pmatrix}$$

1-parameter families

Differentiate!

$$S_m = \dot{S}_m = S_m'(0)$$

$$\Rightarrow S_y = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

$$S_x = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$S_z = \begin{pmatrix} -i & 0 \\ 0 & i \end{pmatrix}$$

$$M^T M = \mathbb{1} \Rightarrow \dot{M}^T + M = 0 \Rightarrow M^T = -M$$

$$M = e^{A\alpha}$$

$$= \mathbb{1} + A\alpha + A^2\alpha^2 + \dots$$

suppose A is diagonalizable

$$\Rightarrow A = P D P^{-1}$$

$$\Rightarrow M = P e^{D\alpha} P^{-1}$$

$$\Rightarrow \det M = \det e^{D\alpha}$$

$$= e^{d_1\alpha} e^{d_2\alpha} \dots e^{d_n\alpha}$$

$$= e^{(d_1 + d_2 + \dots + d_n)\alpha}$$

$$= e^{\text{tr} A \alpha}$$

True in general:

$$\det e^A = e^{\text{tr} A}$$

$$\therefore \det M = 1$$

$$\Rightarrow \text{tr} \dot{M} = 0$$

Commutators

$$[S_x, S_y] = 2S_z$$

$$[S_y, S_z] = 2S_x$$

$$[S_z, S_x] = 2S_y$$

same as $SO(3)$!

$$(S_m \mapsto S_m/2)$$

$$\therefore \boxed{su(2) \cong so(3)}$$

Comparison with $SO(3)$

$$H_3^0(\mathbb{C}) = \left\{ P \in \mathbb{C}^{2 \times 2}; P^t = P, \text{tr} P = 0 \right\}$$

$$P = \begin{pmatrix} z & x-iy \\ x+iy & -z \end{pmatrix} \in \mathcal{V}$$

$$\Rightarrow MPM^t \in H_3^0(\mathbb{R}) \quad \text{for } M \in SU(2)$$

$$\text{But } \det P = -(x^2 + y^2 + z^2)$$

$\Rightarrow M$ preserves Euclidean inner product on \mathbb{R}^3

$\therefore P \mapsto MPM^t$ is in $SO(3)$!

But $\pm M$ induce the same rotation

$\therefore SU(2)$ is the double cover of $SO(3)$

$$\therefore SO(3) \approx \mathbb{R}P^3$$

$$SU(2) \approx S^3$$

SO(3,1)

$$\eta = \begin{pmatrix} -1 & & & \\ & 0 & & \\ & & 0 & \\ & & & 0 \end{pmatrix}, \quad v \in \mathbb{R}^4 \leftrightarrow \begin{pmatrix} t \\ x \\ y \\ z \end{pmatrix}$$

$$|v|^2 = v^T \eta v \leftrightarrow -t^2 + x^2 + y^2 + z^2$$

$$\Rightarrow SO(3,1) = \{ M \in \mathbb{R}^4 : |Mv| = |v| \text{ \& } \det M = 1 \}$$

$$\Rightarrow M^T \eta M = \eta$$

clearly, $SO(3) \subset SO(3,1)$

$$M \mapsto \left(\begin{array}{c|c} 1 & \\ \hline & M \end{array} \right)$$

But also 3 boosts

$$B_x(\alpha) = \left(\begin{array}{cc|c} \cosh \alpha & \sinh \alpha & 0 \\ \sinh \alpha & \cosh \alpha & 0 \\ \hline 0 & 0 & \mathbb{1} \end{array} \right)$$

Warning: $\det(-M) = \det M$
 $\Rightarrow SO(3,1)$ disconnected

Idea: \circ vs $)$ $($

$SL(2, \mathbb{C})$

$$H_3(\mathbb{C}) = \{Q \in \mathbb{C}^{2 \times 2} : Q^\dagger = Q\}$$

$$\Rightarrow Q = \begin{pmatrix} t+z & x-iy \\ x+iy & t-z \end{pmatrix}$$

$$\Rightarrow \det Q = -(-t^2 + x^2 + y^2 + z^2)$$

$$SL(2, \mathbb{C}) = \{M \in \mathbb{C}^{2 \times 2} : \det M = 1\}$$

$$M \in SL(2, \mathbb{C}) \Rightarrow \det M Q M^\dagger = \det M$$

$\Rightarrow SL(2, \mathbb{C})$ is the double cover
of $SO(3, 1)$

(really of its component
connected to \mathbb{I})

$sl(2, \mathbb{C})$ spanned by

S_m

$$T_m = i S_m$$

Pauli
matrices \nearrow

Components

Let $\{e_m\}$ be a basis for Lie algebra \mathfrak{g}

$$\Rightarrow [e_m, e_n] = C_{mn}^a e_a$$

↑
structure constants
for \mathfrak{g}

$$\Rightarrow [[e_m, e_n], e_p] = C_{mn}^a C_{ap}^b e_b$$

|| (by Jacobi identity)

$$\begin{aligned} \Rightarrow [e_m, [e_n, e_p]] - [e_n, [e_m, e_p]] \\ = (C_{mq}^b C_{np}^q - C_{nq}^b C_{mp}^q) e_b \end{aligned}$$

\therefore the matrices $C_m = (C_{mq}^b)$ satisfy

$$C_m C_n - C_n C_m = C_{mn}^a C_a$$

and \therefore reproduce the Lie algebra!
(since structure constants are the same)

"adjoint representation" $e_m \mapsto C_m$

Representations

Ⓘ math

$$\rho: \mathfrak{g} \rightarrow \mathfrak{gl}(V)$$

Lie alg *linear endomorphisms*
vector space

$$[\bar{x}, \bar{y}] \mapsto \rho(\bar{x})\rho(\bar{y}) - \rho(\bar{y})\rho(\bar{x})$$

Lie algebra
homomorphism

matrix multiplication
("composition")

Ex: adjoint: $V = \mathfrak{g}$ "g acting on itself"

$$\text{ad}: \mathfrak{g} \rightarrow \mathfrak{g}$$
$$\bar{x} \mapsto \text{ad}_{\bar{x}}$$

$$\text{ad}_{\bar{x}}(\bar{y}) = [\bar{x}, \bar{y}]$$

$$\begin{aligned} \Rightarrow \text{ad}_{[\bar{x}, \bar{y}]} \bar{z} &= [[\bar{x}, \bar{y}], \bar{z}] \\ &= [\bar{x}, [\bar{y}, \bar{z}]] - [\bar{y}, [\bar{x}, \bar{z}]] \\ &= \text{ad}_{\bar{x}} \text{ad}_{\bar{y}} \bar{z} - \text{ad}_{\bar{y}} \text{ad}_{\bar{x}} \bar{z} \end{aligned}$$

$$\Rightarrow \text{ad}_{[\bar{x}, \bar{y}]} = [\text{ad}_{\bar{x}}, \text{ad}_{\bar{y}}]$$

Ⓙ physics representation \leftrightarrow matrices $\rho(\mathfrak{g})$

Ⓚ particle physics & representation thy: representation \leftrightarrow vector space V

Killing Form

\exists natural inner product on \mathfrak{g} !

$$B(\mathfrak{X}, \mathfrak{Y}) = \text{tr}(\text{ad}_{\mathfrak{X}}, \text{ad}_{\mathfrak{Y}})$$

In practice, $B = \text{tr}(\mathfrak{X}\mathfrak{Y})$ (unique up to overall scale)
usually

Properties :

- symmetric

$$B(\mathfrak{Y}, \mathfrak{X}) = B(\mathfrak{X}, \mathfrak{Y})$$

- bilinear

Def : \mathfrak{g} semisimple \Leftrightarrow B nondegenerate

\mathfrak{g} simple \Leftrightarrow no proper ideals

$$\nexists 0 \neq \mathfrak{h} \subseteq \mathfrak{g} : [\mathfrak{g}, \mathfrak{h}] \subseteq \mathfrak{h}$$

Fact : \mathfrak{g} semisimple $\Leftrightarrow \mathfrak{g} = \mathfrak{g}_1 \oplus \dots \oplus \mathfrak{g}_k$
with \mathfrak{g}_m simple

Ex : $\mathfrak{so}(2)$

$$\text{ad}_{\mathfrak{X}} \mathfrak{Y} = [\mathfrak{X}, \mathfrak{Y}]$$

but $[\mathfrak{X}, \mathfrak{Y}] \equiv 0$!

$$\Rightarrow \text{ad}_{\mathfrak{X}} = 0$$

$$\Rightarrow B \equiv 0$$

degenerate!

Basis

e_p

$$\Rightarrow \text{ad}_{e_p}(e_q) = [e_p, e_q] = C^a_{pq} e_a$$

$$\Rightarrow e_q \mapsto C^a_{pq} e_a$$

$$\Rightarrow \begin{matrix} q \\ \left(\begin{array}{c} \vdots \\ 1 \\ \vdots \end{array} \right) \end{matrix} \mapsto \begin{matrix} C^1_{pq} \\ \vdots \\ C^n_{pq} \end{matrix}$$

$$\Rightarrow \text{ad}_{e_p} = \begin{pmatrix} C^1_{pq} \\ \vdots \\ C^n_{pq} \end{pmatrix} = \begin{pmatrix} C^a_{pq} \end{pmatrix} = C_p$$

$$\begin{aligned} \therefore B(e_p, e_q) &= \text{tr}(C_p C_q) \\ &= \text{tr}(C^a_{pk} C^k_{qb}) = C^m_{pk} C^k_{qm} \end{aligned}$$

Similar argument shows

$$B([X, Y], Z) = B([Y, Z], X)$$

$$\Rightarrow B([Z, X], Y) + B(X, [Z, Y]) = 0$$

$\Rightarrow B$ invariant under G !

G acts on $\mathfrak{g} : X \mapsto M X M^{-1}$

Differentiate: $X \mapsto A X - X A = [A, X]$

$$\Rightarrow B(X, Y) \mapsto 0$$

i.e. B unchanged by M

$$\Rightarrow B(M X M^{-1}, M Y M^{-1}) = B(X, Y)$$

Example $so(5,1)$

- what size matrices are in $so(5,1)$?
 - what is the dimension of $so(5,1)$?
 - what size is matrix representation of Killing form?
 - How many boosts & how many rotations?
-

Example: $so(4)$

Basis 1: $\Gamma_{xy}, \Gamma_{yz}, \Gamma_{zx}, \Gamma_{wz}, \Gamma_{wx}, \Gamma_{wy}$

Basis 2: $\Gamma_{xy} \pm \Gamma_{wz}$
 $\Gamma_{yz} \pm \Gamma_{wx}$
 $\Gamma_{zx} \pm \Gamma_{wy}$

$$so(4) = su(2) \oplus su(2)$$

not simple!!
(only such example)

Representations of $su(2)$

Day 06

- Idea:
- Start with $\mathcal{V}_z \leftarrow$ real eigenvalues!
 - Find eigenvectors \leftarrow in $(\text{complex})/su(2)$!

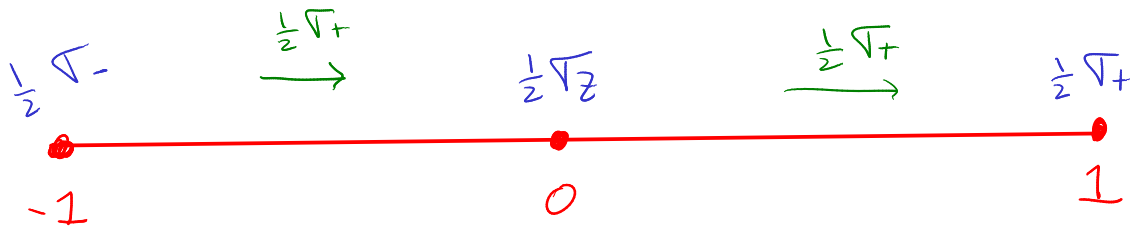
Recall: $[S_x, S_y] = 2S_z \dots$
 $\Rightarrow [\frac{1}{2}S_x, \frac{1}{2}S_y] = 2\frac{S_z}{4} = \frac{1}{2}S_z \dots$

Furthermore, S_z is anti-Hermitian \rightarrow imaginary eigenvalues

\therefore consider $\frac{1}{2}\mathcal{V}_z = \frac{1}{2}S_z \in su(2) \otimes \mathbb{C}$

& $\frac{1}{2}(\underbrace{\mathcal{V}_x \pm i\mathcal{V}_y}_{\mathcal{V}_{\pm}}) = \frac{1}{2}(\mathcal{V}_x \mp S_y)$

\Rightarrow
$$\left[\frac{1}{2}\mathcal{V}_z, \frac{1}{2}\mathcal{V}_{\pm} \right] = \pm \frac{1}{2}\mathcal{V}_{\pm}$$
$$\left[\frac{1}{2}\mathcal{V}_+, \frac{1}{2}\mathcal{V}_- \right] = 2\left(\frac{1}{2}\mathcal{V}_z\right)$$



Adjoint rep!
Root diagram!

$\frac{1}{2}\mathcal{V}_{\pm}$ are raising/lowering operators

$\{ \mathcal{V}_z, \mathcal{V}_x, \mathcal{V}_y \}$ span $sl(2, \mathbb{R})$!
(real form of $su(2)$)

Other Representations

$$\rho: \mathfrak{su}(2) \mapsto \mathfrak{gl}(V)$$

$$L_z = \rho\left(\frac{1}{2}\sigma_z\right)$$

$$L_{\pm} = \rho\left(\frac{1}{2}\sigma_{\pm}\right)$$

$$\Rightarrow \begin{cases} [L_z, L_{\pm}] = \pm L_{\pm} \\ [L_+, L_-] = 2L_z \end{cases}$$

Suppose L_z diagonalizable (true if semi-simple)

Choose basis of eigenvectors $\{w, \dots\}$

$$\Rightarrow L_z w = \lambda w$$

$$\begin{aligned} \Rightarrow L_z L_{\pm} w &= [L_z, L_{\pm}] w + L_{\pm} L_z w \\ &= \pm L_{\pm} w + \lambda L_{\pm} w \\ &= (\lambda \pm 1) L_{\pm} w \end{aligned}$$

For \bar{V} to be irreducible (no proper subreps)

\Rightarrow basis must be $\{L_{\pm}^n w\}$

For \bar{V} to be finite, must have largest eigenvalue

$$\begin{aligned} & \mapsto L_+ w = 0 \quad (\lambda \text{ largest}) \\ \Rightarrow & L_+ L_- w = [L_+, L_-] w + L_- L_+ w \\ & = 2L_z w = 2\lambda w \\ \Rightarrow & L_+ L_- L_- w = [L_+, L_-] L_- w + L_- L_+ L_- w \\ & = 2L_z L_- w + 2\lambda L_- w \\ & = 2(2\lambda - 1)L_- w \\ & \vdots \\ \Rightarrow & L_+ L_-^k w = (2k\lambda - k(k-1)) L_-^{k-1} w \end{aligned}$$

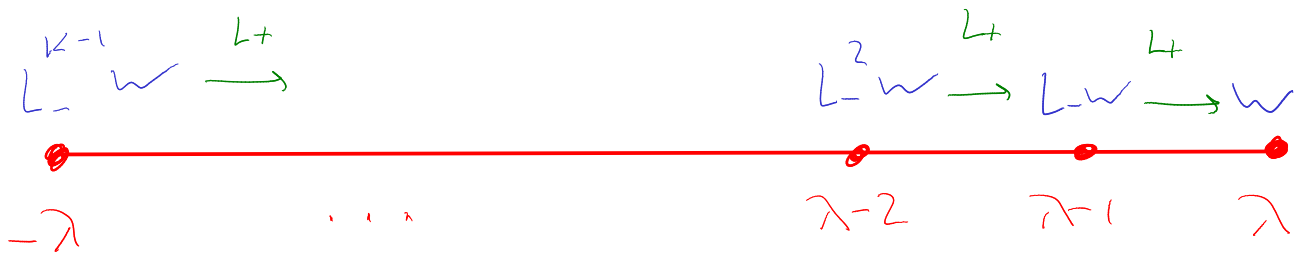
But finite $\Rightarrow L_-^k w = 0$ for some k

$$\Rightarrow 2k\lambda - k(k-1) = 0$$

$$\Rightarrow \lambda = \frac{k-1}{2} \in \frac{\mathbb{Z}}{2}! \quad \text{"spin"}$$

$$\& \Rightarrow k = 2\lambda + 1$$

$$\therefore \lambda, \lambda-1, \dots, \lambda - (2\lambda) = -\lambda$$



Weight diagram of any irrep of $su(2)$

su(3)

Gell-Mann matrices:

$$\lambda_1 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad \lambda_2 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad \lambda_3 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\lambda_4 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}, \quad \lambda_5 = \begin{pmatrix} 0 & 0 & i \\ 0 & 0 & 0 \\ -i & 0 & 0 \end{pmatrix}$$

$$\lambda_6 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, \quad \lambda_7 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix}, \quad \lambda_8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$

basis for $su(3) \otimes \mathbb{C}$

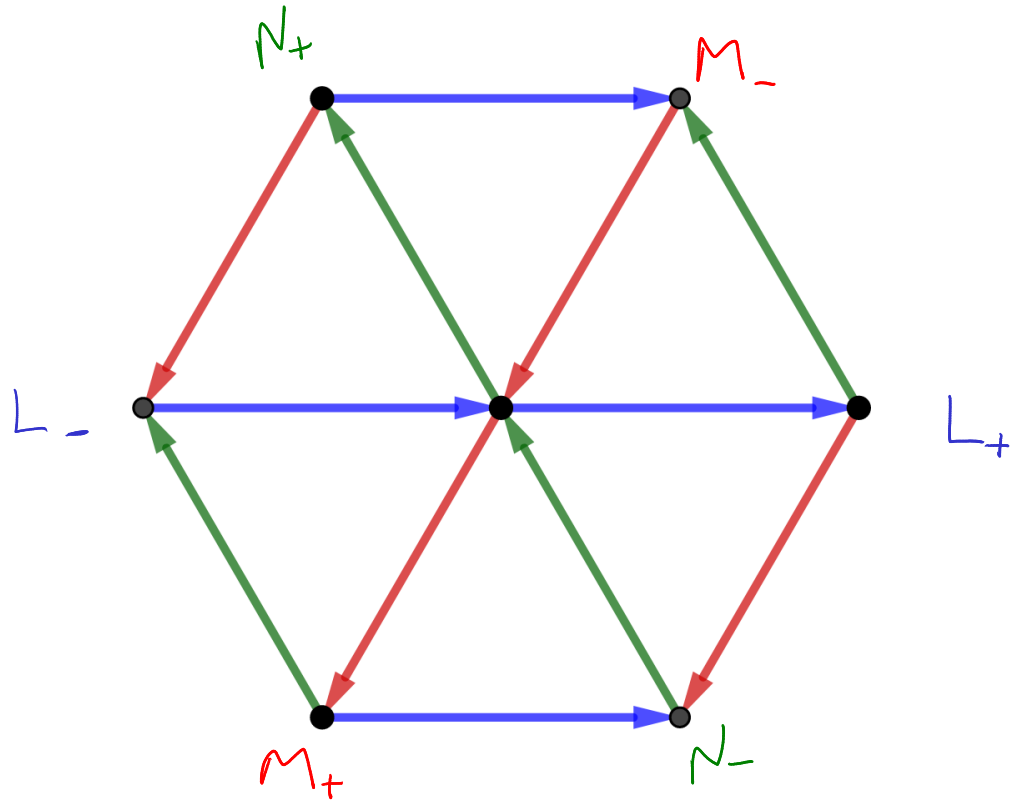
($\mu_m = -i\lambda_m$ is basis for $su(3)$)

$$\longmapsto sl(3, \mathbb{R}) \leftrightarrow \lambda_1, \mu_2, \lambda_3, \lambda_4, \mu_5, \lambda_6, \mu_7, \lambda_8$$

\Rightarrow simultaneous eigenvectors of λ_3, λ_8 are

		λ_3	λ_8
$2L_{\pm}$	$\lambda_1 \pm i\lambda_2$	± 2	0
$2M_{\pm}$	$\lambda_4 \pm i\lambda_5$	∓ 1	$\mp \sqrt{3}$
$2N_{\pm}$	$\lambda_7 \pm i\lambda_8$	∓ 1	$\pm \sqrt{3}$
	λ_3	0	0
	λ_8	0	0

Root Diagram of $su(3)$



origin is 2-d vector space
spanned by λ_3, λ_8

Weight Diagram of $su(3)$ rep

Consider defining rep of $su(3)$:

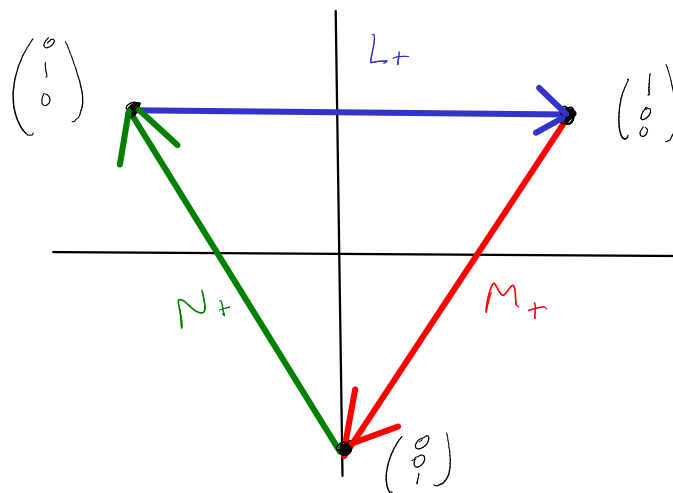
$$\left\{ A \in \mathbb{C}^{3 \times 3} : A^\dagger + A = 0, \text{tr} A = 0 \right\}$$

acting on \mathbb{C}^3

As before, complexify & choose 2 commuting elements

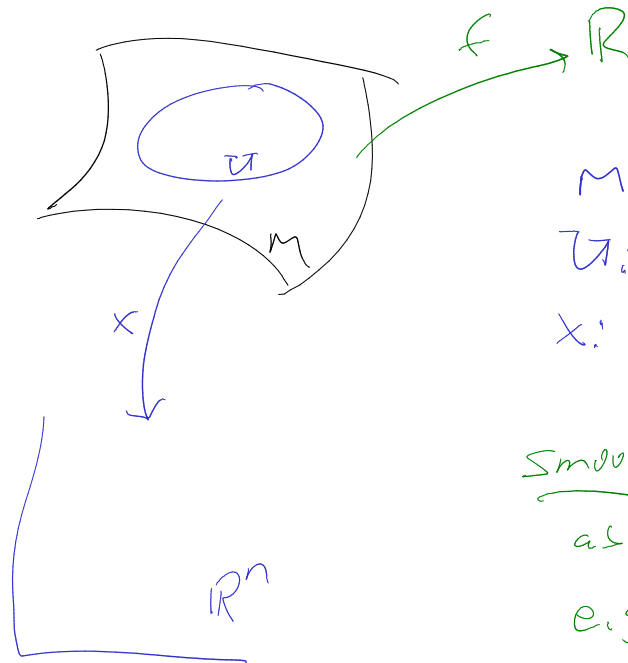
Now, find eigenvectors & eigenvalues

	λ_3	λ_8
$\begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$	1	$1/\sqrt{3}$
$\begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$	-1	$1/\sqrt{3}$
$\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$	0	$-2/\sqrt{3}$



Differential Geometry

Idea:



M : topological space

U : open subset

x : coordinates ("chart")

smooth means smooth
as maps on \mathbb{R}^n

e.g. $f \circ x^{-1} : \mathbb{R}^n \rightarrow \mathbb{R}$

vector fields \Leftrightarrow directional derivatives
(w.r.t. coordinates)

Left-Invariant Vector Fields

Day 07

Idea: Given Lie group G

\mapsto Lie algebra \mathfrak{g} is $T_{\mathbb{1}}G$

\parallel
tangent space
at $\mathbb{1} \in G$

But $M \in G \Rightarrow M: G \rightarrow G$

$\mathbb{1} \mapsto M$

"push-forward"
"differential" $\Rightarrow M_*: T_{\mathbb{1}}G \rightarrow T_M G$

\therefore get vector field
from any vector $A \in \mathfrak{g}$

$$\mathbb{X}|_M = M_* A$$

$$\Rightarrow M_* \mathbb{X} = \mathbb{X}$$

left-invariant
vector field

"nice" curves

\Rightarrow integral curves through $\mathbb{1}$ of
(left-invariant) vector fields satisfy

$$M(\alpha) = e^{A\alpha}$$

1-parameter
family

Properties of Killing form

Recall : • $B(\underline{Y}, \underline{X}) = B(\underline{X}, \underline{Y})$ Symmetric

• $B([\underline{X}, \underline{Y}], \underline{Z}) = B([\underline{Y}, \underline{Z}], \underline{X})$

Cyclic

$$\Rightarrow B([\underline{Z}, \underline{X}], \underline{Y}) + B(\underline{X}, [\underline{Z}, \underline{Y}]) = 0$$

invariant under G

$\Rightarrow B$ is constant on
left-invariant
vector fields!

Connection

Given a metric B on a Lie group G
 $\exists!$ Levi-Civita connection, that is, a
 derivative operator ∇_Z on vector fields:

- $\nabla_Z f = Z(f)$ (directional derivative)
- $\nabla_Z (B(X, Y)) = B(\nabla_Z X, Y) + B(X, \nabla_Z Y)$
 $= Z(B(X, Y))$ (metric compatible)
- $\nabla_X Y - \nabla_Y X = [X, Y]$ (torsion free)

But $B(X, Y) = \text{const}$ (for left-invariant vector fields)!

$$\begin{aligned}
 \Rightarrow B(\nabla_X Y, Z) &= B([X, Y], Z) + B(\nabla_Y X, Z) \\
 &= B([X, Y], Z) - B(X, \nabla_Y Z) \\
 &= B([X, Y], Z) - B(X, [Y, Z]) - B(X, \nabla_Z Y) \\
 &= B([X, Y], Z) - B(X, [Y, Z]) + B(\nabla_Z X, Y) \\
 &= B([X, Y], Z) - B([Y, Z], X) + B([Z, X], Y) \\
 &\quad - B(\nabla_X Z, Y) \\
 &= B([X, Y], Z) - \cancel{B([Y, Z], X)} + \cancel{B([Z, X], Y)} \\
 &\quad \text{(cyclic!)} \quad + B(Z, \nabla_X Y)
 \end{aligned}$$

$$\Rightarrow 2 B(\nabla_X Y, Z) = B([X, Y], Z) \quad \nabla Z$$

$$\Rightarrow \nabla_X Y = \frac{1}{2} [X, Y]$$

for left-invariant
vector fields

Curvature

Given a connection ∇_Z , the curvature operator is

$$R(X, Y)Z = \nabla_X \nabla_Y Z - \nabla_Y \nabla_X Z - \nabla_{[X, Y]} Z$$

Here, $\nabla_X Y = \frac{1}{2} [X, Y]$, so

$$R(X, Y)Z = \frac{1}{4} [X, [Y, Z]] - \frac{1}{4} [Y, [X, Z]]$$

$$- \frac{1}{2} [[X, Y], Z]$$

Jacobi
identity

$$\Rightarrow = \frac{1}{4} [[X, Y], Z] - \frac{1}{2} [[X, Y], Z]$$

$$\text{so } R(X, Y)Z = -\frac{1}{4} [[X, Y], Z]$$

for left-invariant vector fields

Ex: $so(3)$, basis $e_a = \{e_x, e_y, e_z\}$

$$\Rightarrow R(u, v)w = R^a_{bcd} w^b u^c v^d e_a$$

$$\Rightarrow R(e_x, e_y)e_z = 0$$

$$R(e_x, e_y)e_x = -\frac{1}{4} e_y$$

components of
Riemann tensor

$$\Rightarrow R^a_{zxy} = 0, R^y_{xxy} = -\frac{1}{4} \Rightarrow R^xy_{xy} = +\frac{1}{4}$$

$$\Rightarrow R^x_x = \frac{1}{2} = R^y_y = R^z_z \Rightarrow R = \frac{3}{2} = \text{const} \checkmark$$

Ricci tensor: $R^a_b = R^{ma}_{mb}$

Ricci scalar: $R = R^m_m = \text{tr Ric}$