

Spinors

Recall: $so(3) \leftrightarrow 3 \times 3$ real, antisymmetric
acts on: 3-component real
column vectors
"vectors"

$su(2) \leftrightarrow 2 \times 2$ complex, antihermitian
acts on: 2-component complex
column vectors
"spinors"

$$\mathbb{R}^3 \leftrightarrow \langle \sigma_m \rangle \leftrightarrow \begin{pmatrix} z & x-iy \\ x+iy & -z \end{pmatrix} = \underline{\Sigma}$$

$$M \in SU(2) : \underline{\Sigma} \mapsto M \underline{\Sigma} M^\dagger$$

$$\Rightarrow A \in su(2) : \underline{\Sigma} \mapsto [A, \underline{\Sigma}]$$

\therefore $SU(2)$ is double cover of $SO(3)$

\therefore vectors and Lie group/algebra are matrices

$$\uparrow$$
$$\sigma_m$$

$$\uparrow$$
$$S_m = -i \sigma_m$$

But $\sigma_x \sigma_y = i \sigma_z$.

What other products are there?

$$\mathbb{1} = \sigma_m^2 ; \quad i \mathbb{1} = \sigma_x \sigma_y \sigma_z$$

Clifford Algebras

Given

- vector space V
- nondegenerate inner product g

\Rightarrow Clifford algebra $Cl(V)$:

add bilinear, associative product

$$\underline{x}^2 = |\underline{x}|^2 = g(\underline{x}, \underline{x})$$

$$\Leftrightarrow \{ \underline{x}, \underline{y} \} = 2g(\underline{x}, \underline{y})$$

anticommutator

$$\underline{x}\underline{y} + \underline{y}\underline{x}$$

symmetric

$$\underline{PF}: (\underline{x} + \underline{y})^2 = \cancel{\underline{x}^2} + \cancel{\underline{y}^2} + \{ \underline{x}, \underline{y} \}$$

$$\underline{PF}: g(\underline{x} + \underline{y}, \underline{x} + \underline{y}) = \cancel{|\underline{x}|^2} + \cancel{|\underline{y}|^2} + 2g(\underline{x}, \underline{y})$$

Orth onormal

Basis! $\{ \gamma_i, \gamma_j \} = \pm 2 \delta_{ij}$

gamma matrices!

signature!

$$V = \mathbb{R}^{p, q} \Rightarrow Cl(V) = Cl(p, q)$$

- ① $\langle \sigma_i \rangle \cong V \subset \text{Cl}(V)$ rank 1
- ② $\sigma_i^2 = \pm 1 \in \text{Cl}(V)$ rank 0
- ③ $\sigma_{i_1 \dots i_k} = \sigma_{i_1} \dots \sigma_{i_k}$ basis for rank k
if $1 \leq i_1 < \dots < i_k \leq n = p+q$

$\therefore \binom{n}{k}$ elements of rank k
& $\dim \text{Cl}(p+q)$ is n^2

Exterior Algebra

$\wedge(V)$: bilinear, associative product



$$v \wedge v = 0$$

$$w \wedge v = -v \wedge w$$

wedge

antisymmetric

\Rightarrow basis is $e_1, \dots, e_n = e_1 \wedge \dots \wedge e_n$

for each rank $k \leq n$

Same size

Extend to $\mathbb{1}$: $\mathbb{1} \wedge v = v$

rank 0

Again, $\forall \alpha \in \wedge(V)$ as rank - 1 elements

Ex : Differential forms \leftarrow vector fields

$$df(\underline{x}) = \underline{x}(f) = \sum^m e_m(f)$$

directional derivative

Clifford product

$$\therefore \text{define } a \vee b = a \wedge b + g(a, b)$$

vec

wedge

dot

Clifford Rotations

Recall: $\mathbb{R}^{p,q} \subset \mathcal{C}\ell(p,q)$

$$v = v^m \gamma_m$$

Assume first that σ_i, σ_j are spacelike

$$\begin{aligned} \therefore R_{ij}(d) : \quad & \sigma_i \longmapsto \sigma_i \cos d - \sigma_j \sin d \\ & \sigma_j \longmapsto \sigma_i \sin d + \sigma_j \cos d \\ & \sigma_k \longmapsto \sigma_k \quad k \neq i, j \end{aligned}$$

\curvearrowright
 $SO(p,q)$

$$\Rightarrow r_{ij} : \begin{array}{l} \sigma_i \longmapsto -\sigma_j \\ \sigma_j \longmapsto \sigma_i \\ \sigma_k \longmapsto 0 \end{array}$$

\leftarrow If boost, use \sinh & \cosh (& drop minus sign)

Calculate: $[\sigma_i \sigma_j, \sigma_i] = -2 |\sigma_i|^2 \sigma_j$
 $[\sigma_i \sigma_j, \sigma_j] = +2 |\sigma_j|^2 \sigma_i$
 $[\sigma_i \sigma_j, \sigma_k] = 0$

$$\therefore r_{ij} \leftrightarrow \text{commute with } \frac{1}{2} \sigma_i \sigma_j$$

rank 0:	$\mathbb{1}$	scalars	\mathbb{R}
rank 1:	σ_i	vectors	$\mathbb{R}^{p,q}$
rank 2:	r_{ij}	rotations	$SO(p,q)$

Geometric Algebra

So $so(p, q) \subset \mathcal{C}\mathcal{O}(p, q)$.
What about $S\mathcal{O}(p, q)$?

Exponentiate!!

$$\begin{aligned} e^{\frac{1}{2}\sigma_i\sigma_j\alpha} &= \mathbb{1} + \frac{1}{2}\sigma_i\sigma_j\alpha + \frac{1}{4}(\sigma_i\sigma_j)^2\alpha^2 + \dots \\ &= \mathbb{1} \left(1 - |\sigma_i|^2|\sigma_j|^2\left(\frac{\alpha}{2}\right)^2 + \dots \right) \\ &\quad + \sigma_i\sigma_j \left(\frac{\alpha}{2} - |\sigma_i|^2|\sigma_j|^2\left(\frac{\alpha}{2}\right)^2 + \dots \right) \\ &= \begin{cases} \cos \frac{\alpha}{2} + \sigma_i\sigma_j \sin \frac{\alpha}{2} & (|\sigma_i|^2|\sigma_j|^2 = +1) \\ \cosh \frac{\alpha}{2} + \sigma_i\sigma_j \sinh \frac{\alpha}{2} & (|\sigma_i|^2|\sigma_j|^2 = -1) \end{cases} \\ &= R_{ij} \left(\frac{\alpha}{2} \right) \end{aligned}$$

$\therefore S\mathcal{O}(p, q) \subset \mathcal{C}\mathcal{O}(p, q) !$

Classification of Clifford Algebras

Idea: Find the smallest matrices that represent $Cl(p, q)$

- $Cl(p, q)$ is a real vector space
(can also classify over \mathbb{C})
- matrices can be real, complex, quaternionic

Vectors: copy of $\mathbb{R}^{p, q} \subset Cl(p, q)$
 $\{\sigma_m\}$

Spinors: column vectors acted on by $Cl(p, q)$

Special cases:

- smallest real matrices have same size as smallest complex matrices

↳ Majorana spinors

- "spinors" not irreducible rep of $so(p, q)$
($\Rightarrow \exists ! 2$ irreducible reps)

↳ Weyl spinors

Example : $\{\gamma_m\} = \text{orthonormal basis}$

$$\Rightarrow \begin{aligned} \gamma_m^2 &= +\mathbb{1} & (m=1 \dots p) \\ \gamma_m^2 &= -\mathbb{1} & (m=p+1 \dots p+q=n) \end{aligned}$$

$$\omega = \gamma_1 \dots \gamma_n$$

$$\begin{aligned} \Rightarrow \omega^2 &= (\gamma_1 \dots \gamma_n) (\gamma_1 \dots \gamma_n) \\ &= (-1)^{n(n-1)/2} \gamma_1^2 \dots \gamma_n^2 \\ &= (-1)^{n(n-1)/2} (-1)^q \mathbb{1} \end{aligned}$$

But : $n \frac{(n-1)}{2} + q = \frac{(p+q)(p+q-1) + 2q}{2}$

$$= \frac{(p+q)^2 + q - p}{2} = \frac{4pq + (p-q)^2 + q - p}{2}$$

depends only on $p-q$

$$= \cancel{2pq} + \frac{(p-q)(p-q-1)}{2}$$

even

$$\Rightarrow \begin{aligned} \omega^2 &= \mathbb{1} & (p-q \equiv 0, 1 \pmod{4}) \\ \omega^2 &= -\mathbb{1} & (p-q \equiv 2, 3 \pmod{4}) \end{aligned}$$

correctly suggests that classification depends only on

$$p-q \pmod{8}$$

Spinorial Chessboard

Examples

$$\underline{\text{Cl}(1,0)} : \mathcal{V} = \langle \sigma_z \rangle$$

$$\Rightarrow \underline{\text{Cl}(1,0) = \mathbb{R} \oplus \mathbb{R}}$$

$$\underline{\text{Cl}(0,1)} : \mathcal{V} = \langle s_y \rangle$$

$$\Rightarrow \underline{\text{Cl}(0,1) = \mathbb{C}}$$

$$\underline{\text{Cl}(1,0) \text{ revisited}} : \mathcal{V} = \langle \sigma_x \rangle$$

$$\Rightarrow \underline{\text{Cl}(1,0) = \mathbb{C}'}$$

$$\underline{\text{Cl}(2,0)} : \mathcal{V} = \langle \sigma_x, \sigma_z \rangle \Rightarrow \omega = s_y$$
$$\Rightarrow \omega^2 = -\mathbb{I}$$
$$\Rightarrow \underline{\text{Cl}(2,0) = \mathbb{R}^{2 \times 2}}$$

$$\underline{\text{Cl}(1,1)} : \mathcal{V} = \langle \sigma_z, s_y \rangle \Rightarrow \omega = \sigma_z$$
$$\Rightarrow \omega^2 = +\mathbb{I}$$

swap
 $\sigma_x \leftrightarrow s_y!$

$$\Rightarrow \underline{\text{Cl}(2,0) = \mathbb{R}^{2 \times 2}}$$

$$\underline{\text{Cl}(0,2)} : \mathcal{V} = \langle s_x, s_y \rangle \Rightarrow \omega = s_z$$

$\uparrow \quad \uparrow \quad \uparrow$
 $i \quad j \quad k$

$$\therefore \underline{\text{Cl}(0,2) = \mathbb{H}} \quad !$$

Clifford Algebras over Division Algebras

Weyl representation of $Cl(3,1)$

$$\sigma_z = \left(\begin{array}{c|c} & 1 \\ \hline -1 & \end{array} \right) = -s_y \otimes \mathbb{1} = i\sigma_y \otimes \mathbb{1}$$

$$\sigma_m = \left(\begin{array}{c|c} & \sigma_m \\ \hline \sigma_m & \end{array} \right) = \sigma_x \otimes \sigma_m$$

\Rightarrow vectors (rank 1) are block off-diagonal

\Rightarrow adjoint (rank 2) is block diagonal

\Rightarrow 2 Weyl spinor representations $\left(\begin{array}{c} 1 \\ \hline 2 \end{array} \right)$

Step 1: $x\sigma_x + y\sigma_y = \begin{pmatrix} 0 & x-iy \\ \textcircled{x+iy} & 0 \end{pmatrix}$
 $\in \mathbb{C}$

\therefore generalize to \mathbb{H}, \mathbb{O} !

$$\mapsto \underline{x} = \begin{pmatrix} 0 & \bar{a} \\ a & 0 \end{pmatrix} \quad \underline{a} \in \mathbb{D}$$

$$\sigma_{\underline{x}} = \left(\begin{array}{c|c} 0 & \underline{x} \\ \hline \underline{x} & 0 \end{array} \right) \Rightarrow \{ \sigma_x, \sigma_{\underline{x}} \} \leftrightarrow \{ \underline{x}, \underline{y} \} \\ \leftrightarrow 2a \cdot b$$

$so(8)$!

Step 2! Can add δ_z, δ_t back in

$$\mapsto \underline{X} = \begin{pmatrix} t+z & a \\ a & t-z \end{pmatrix} \quad \underline{a} \in \mathbb{O}$$

$$\delta_{\underline{X}} = \left(\begin{array}{c|c} 0 & \underline{X} \\ \hline \underline{\tilde{X}} & 0 \end{array} \right)$$

$$\underline{\tilde{X}} = \underline{X} - \text{tr}(\underline{X}) \mathbb{1}$$

Aside: 2×2 characteristic equation:

$$\underline{X}^2 - \underline{X} \text{tr} \underline{X} + \det \underline{X} \mathbb{1} = 0$$

$$\Rightarrow \underline{X} \underline{\tilde{X}} = -\det \underline{X} \mathbb{1}$$

$$-\det \underline{X} = -t^2 + z^2 + |a|^2$$

$$\therefore \quad \mathbb{R} \leftrightarrow SO(2,1)$$

$$\mathbb{C} \leftrightarrow SO(3,1)$$

$$\mathbb{H} \leftrightarrow SO(5,1)$$

$$\mathbb{O} \leftrightarrow SO(9,1)$$

Classical supersymmetry only works
in 3, 4, 6, 10 spacetime dimensions!!

Step 3 : $t\sigma_1 + z\sigma_2 \leftrightarrow \left(\begin{array}{c|c} t+z & \\ \hline & z-z \end{array} \right) \leftrightarrow \begin{pmatrix} t+z \\ z-z \end{pmatrix}$

$\mapsto tL\sigma_1 + z\sigma_2 \leftrightarrow \begin{pmatrix} z+z \\ -z+t \end{pmatrix} \in \mathbb{C}'!$

\therefore generalize to $\mathbb{H}^1, \mathbb{O}'!$

$\Sigma = \begin{pmatrix} A & 0 \\ 0 & -\bar{A} \end{pmatrix} \quad A \in \mathbb{C}' \Rightarrow \tilde{\Sigma} = \begin{pmatrix} \bar{A} & 0 \\ 0 & -A \end{pmatrix}$

$\sigma_{\Sigma} = \left(\begin{array}{c|c} \Sigma & \\ \hline \tilde{\Sigma} & \end{array} \right) \Rightarrow \{\sigma_{\Sigma}, \sigma_{\tilde{\Sigma}}\} \leftrightarrow 2A \cdot \bar{A}$

$SO(4,4)$

Step 4 : Do both!!

$\Sigma = \begin{pmatrix} A & \bar{a} \\ a & -\bar{A} \end{pmatrix}, \quad a \in \mathbb{O}, A \in \mathbb{O}'$

$\sigma_{\Sigma} = \left(\begin{array}{c|c} \Sigma & \\ \hline \tilde{\Sigma} & \end{array} \right) \quad (-\det \Sigma = |a|^2 + |A|^2)$

$\mapsto \underline{\underline{SO(12,4)}}$

The 2x2 Magic Square

Recall: $\underline{X} = \begin{pmatrix} A & \bar{a} \\ a & -\bar{A} \end{pmatrix} = x^m \Gamma_m$

$\Rightarrow \underline{X} \in \mathbb{R}^{p,q}$

$a \in \mathbb{K} \leftrightarrow (k, 0)$

$A \in \mathbb{K}' \leftrightarrow (\frac{k'}{2}, \frac{k'}{2})$

$p = k + \frac{k'}{2}, q = \frac{k'}{2}$

$\mapsto \Gamma_m = \begin{pmatrix} 0 & \Gamma_m \\ \tilde{\Gamma}_m & 0 \end{pmatrix}$ generates $Cl(p, q)$

$\therefore \Gamma_m \Gamma_n$ generates $so(p, q)$

two copies!

$\therefore \Gamma_m \tilde{\Gamma}_n$ generates $so(p, q)$

↑
anti-Hermitian

$\mapsto \underline{so(p, q)} \cong \underline{su(2, 1\mathbb{K}' \oplus \mathbb{K})}$

	\mathbb{R}	\mathbb{C}	\mathbb{H}	\mathbb{O}	\mathbb{K}
\mathbb{R}'	$so(2)$	$so(3)$	$so(5)$	$so(9)$	$su(2, \mathbb{K})$
\mathbb{C}'	$so(2, 1)$	$so(3, 1)$	$so(5, 1)$	$so(9, 1)$	$sl(2, \mathbb{K})$
\mathbb{H}'	$so(3, 2)$	$so(4, 2)$	$so(6, 2)$	$so(10, 2)$	$sp(2, \mathbb{K})$
\mathbb{O}'	$so(5, 4)$	$so(6, 4)$	$so(8, 4)$	$so(12, 4)$?

This is half-split case. Also have double-split & compact

Spinors

Recall: $so(p, q)$ acts on vectors \mathbb{R} as

$$\mathbb{R} \mapsto [\Gamma_m \Gamma_n, \mathbb{R}] \\ = \Gamma_m(\Gamma_n \mathbb{R}) - \mathbb{R}(\Gamma_m \Gamma_n)$$

Spinors are just 2-component column vectors
over $\mathbb{K}' \otimes \mathbb{K}$!

$\therefore \chi = \left(\begin{array}{c|c} \mathbb{R} & \theta \\ \hline -\theta^\dagger & \phi \end{array} \right)$ combines vector \mathbb{R} with
spinor θ and scalar ϕ

Furthermore $m = \left(\begin{array}{c|c} m & 0 \\ \hline 0 & i \end{array} \right)$ acts as

$$m \chi m^{-1} = \left(\begin{array}{c|c} m \mathbb{R} m^{-1} & m \theta \\ \hline -\theta^\dagger m^\dagger & \phi \end{array} \right)$$

\leftarrow vector
 \leftarrow spinor
 \leftarrow scalar

works well for representations in
first two rows - but not otherwise

\mathbb{R} has only \mathbb{K} on off diagonal;

θ has $\mathbb{K}' \otimes \mathbb{K}$,

\therefore doesn't close

3 x 3 Magic Square

Idea: $\mathfrak{su}(3, \mathbb{K}' \oplus \mathbb{K})$

\therefore start with $\sum_{aA} = \begin{pmatrix} 0 & aA & 0 \\ -\bar{aA} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$

& $\mathbb{I}_{aA}, \mathbb{Z}_{aA}$ cyclic

- How many?
- How many products?

Answer: over \mathbb{H} , need 3 extra diagonal elements as before

But over \mathbb{O} only need 14, not 21

$$\mathfrak{so}(\mathbb{I}_m \mathbb{H}) = \mathfrak{so}(3)$$

$$\mathfrak{so}(\mathbb{I}_m \mathbb{O}) = \mathfrak{so}(7)$$

but

$$\text{Aut } \mathbb{H} = \text{SO}(3)$$

$$\text{Aut } \mathbb{O} = \mathfrak{g}_2 \subset \text{SO}(7)$$

Automorphisms

want!

$$\Phi(xy) = \Phi(x)\Phi(y)$$

$$\Rightarrow \underline{\Phi(1) = 1}$$

\mathbb{R} : Done! Nothing to do!

\mathbb{A} : Only automorphism is $i \mapsto -i$

\mathbb{H} : $x \mapsto qxq^{-1}$

$$\begin{aligned} \Rightarrow xy &\mapsto qxq^{-1}qyq^{-1} \\ &= qxyq^{-1} \end{aligned}$$

$\therefore \text{Aut } \mathbb{H} \leftrightarrow \underline{\text{unit quaternions}}$
 $\leftrightarrow \text{SO}(3)$

$$\text{Aut } \mathbb{H} = \text{SO}(3)$$

$\underline{\Phi}$: $(qxq^{-1})(qyq^{-1}) \neq q(xy)q^{-1}$!

G_2

Count Aut(\mathcal{O}):

① $i \mapsto j$	6 choices	$\mathbb{S}^6 \subset \mathbb{R}^7$
② $j \mapsto l$	5 choices	$\mathbb{S}^5 \subset \mathbb{R}^6$
③ $k \mapsto jl$	0 choices!	
④ $l \mapsto i$	3 choices	$\mathbb{S}^3 \subset \mathbb{R}^4$

$$\therefore \underline{\dim \text{Aut}(\mathcal{O}) = 14}$$

Idea: Rotate $ij \Rightarrow$ must also rotate il and for k kl

$$k = ij = (jl)(il) = l(kl)$$

Fact: $i \mapsto j \mapsto -i$
& $jl \mapsto -il \mapsto -jl$

opposite
orientations

works

Count: 7 planes \times 2 indpt rotations
 $= 14$ ✓

Albert Algebra

Recall: $\Sigma = \begin{pmatrix} tL+z & \bar{a} \\ a & tL-z \end{pmatrix} \in H_2(\mathbb{O})$
Hermitian

$\Rightarrow -\det \Sigma = -z^2 + \bar{z}^2 + |a|^2$

$SO(K+1, 1)$

\cong
 $sl(2, \mathbb{K})$

\cong

$su(2, \mathbb{C} \otimes \mathbb{K})$

Try 3x3:

$\alpha = \begin{pmatrix} p & \bar{a} & c \\ a & m & \bar{b} \\ \bar{c} & b & n \end{pmatrix} \in H_3(\mathbb{O})$

Jordan algebra!

$$(\alpha \circ a y) \circ \alpha^2 = \alpha \circ (a y \circ \alpha^2)$$

needed for observables in quantum mechanics
via density matrices " $\sum |\psi\rangle\langle\psi|$ "

But: need commutative product.

\therefore Jordan product

$$\alpha \circ a y = \frac{1}{2}(\alpha a y + a y \alpha)$$

Fact: all Jordan algebras except Albert algebra
arise from Jordan product on associative
algebra.

"exceptional Jordan algebras"

Properties

$$\textcircled{1} \quad \mathcal{X}^2 = \mathcal{X} \circ \mathcal{X}$$

$$\textcircled{2} \quad \mathcal{X}^3 = \mathcal{X}^2 \circ \mathcal{X} = \mathcal{X} \circ \mathcal{X}^2 \quad \text{by definition}$$

not defined otherwise)

Freudenthal product:

$$\mathcal{X} * \mathcal{Y} = \mathcal{X} \circ \mathcal{Y} - \frac{1}{2} (\mathcal{X} \operatorname{tr} \mathcal{Y} + \mathcal{Y} \operatorname{tr} \mathcal{X}) + \frac{1}{2} (\operatorname{tr} \mathcal{X} \operatorname{tr} \mathcal{Y} - \operatorname{tr} (\mathcal{X} \circ \mathcal{Y}))$$

$\Rightarrow \mathcal{X}$ satisfies characteristic equation

$$\Leftrightarrow \det \mathcal{X} = \frac{1}{3} \operatorname{tr} (\mathcal{X} * \mathcal{X} \circ \mathcal{X})$$

Idea: Consider $v, w \in \mathbb{R}^3$

$$\Rightarrow \mathcal{X} = vv^t, \mathcal{Y} = ww^t \in H_3(\mathbb{R})$$

$$\Rightarrow \begin{aligned} \operatorname{tr} (\mathcal{X} \circ \mathcal{Y}) &= (v \cdot w)^2 \\ \mathcal{X} * \mathcal{Y} &= (v \times w)(v \times w)^t \end{aligned}$$

"dot & cross products!"

F_4 and E_6

$$F_4 = \text{SU}(3, \mathbb{O})$$

\cap

$$E_6 = \text{SL}(3, \mathbb{O})$$

preserves trace
of $\alpha \in H_3(\mathbb{O})$

preserves determinant

Count: $\underline{f_4}$: \mathbb{O}^3 on off diagonal
 $\leftrightarrow 24$

$\text{so}(8)$ on 2×2 diagonal

$\leftrightarrow 28$

Nothing more!!

$$\therefore \underline{\dim f_4 = 52}$$

e_6 : $(\mathbb{C}' \otimes \mathbb{O})^3$ on off diagonal
 $\leftrightarrow 48$

$\text{so}(8)$ on diagonal over \mathbb{O}

$\leftrightarrow 28$

λ_3, λ_8 on \mathbb{C}' diagonal

$\leftrightarrow 2$

$$\therefore \underline{\dim e_6 = 78}$$

Decompositions

Given any subalgebra $h \subset g$, h acts on h^\perp

(nondgenerate)

Ex: $so(3) = so(2) \oplus \mathbb{R}^2$

Annotations:
- $so(2)$: Lie alg
- $so(2)$: Lie subalg
- \mathbb{R}^2 : representation of Lie subalg
- S_z (pointing to $so(2)$)
- S_x, S_y (pointing to \mathbb{R}^2)

Orthogonal

$$so(p+1) = so(p) \oplus \mathbb{R}^p$$

$$so(p+q) = so(p) \oplus so(q) \oplus \mathbb{R}^{pq}$$

Unitary

$$su(p+1) = su(p) \oplus \mathbb{C}^p \oplus u(1)$$

$$su(p+q) = su(p) \oplus su(q) \oplus \mathbb{C}^{pq} \oplus u(1)$$

E₈

Fact: smallest representation of e_8 is $e_{8,24}$.

Problem: How to act on nested elements?!

Instead: Decompose!

$$M \in SU(2, K' \otimes K), \theta \in (K' \otimes K)^2$$

$$m = \left(\begin{array}{c|c} M & \theta \\ \hline -\theta^\dagger & \alpha \end{array} \right) \in SU(3, K \otimes K')$$

$\therefore M \in SO\left(K + \frac{K'}{2}, \frac{K'}{2}\right)$ adjoint

\oplus spinor of $SO\left(K + \frac{K'}{2}, \frac{K'}{2}\right)$

α ?? — but not needed over \oplus

$$\therefore e_{8(24)} = SO(12, 4) \oplus \mathbb{R}^{128}$$

↑
contains

↑
Majorana-Weyl
spinor of $SO(12, 4)$

Standard model \rightarrow $SU(3) \oplus SU(2) \oplus U(1)$
 $\oplus SO(3, 1)$

↑
Lorentz