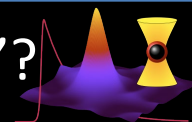


# Organic semiconductors

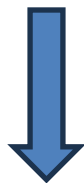
PH 673

Nanoscience and nanotechnology

November 10, 2025



Active layers are made of carbon-based molecules



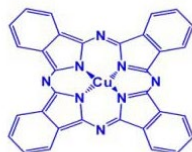
“small” molecules



(a)



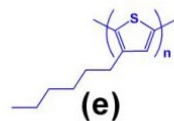
(d)



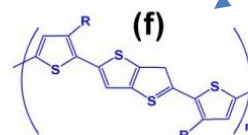
(b)



(c)

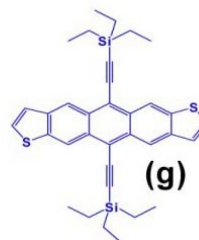


(e)

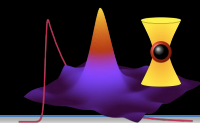


(f)

polymers



(g)



## Traditional silicon technology

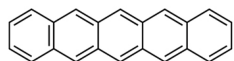
- High performance

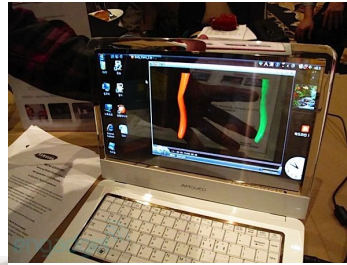
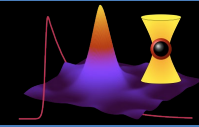
BUT

- Expensive fabrication and processing: high temperature, high pressure, special substrates
- Toxic byproducts of manufacturing
- Challenging to scale down to nanostructures

## Organic materials

- Lower performance  
BUT
- Solution processable
- Low cost
- Flexible, lightweight substrates
- Tunable properties
- Scalable to single-molecule level





[www.oled.com](http://www.oled.com)



<https://pid.samsungdisplay.com/en/digital-signage/transparent-displays>

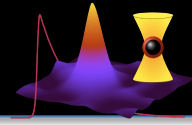


[www.oled-info.com](http://www.oled-info.com)

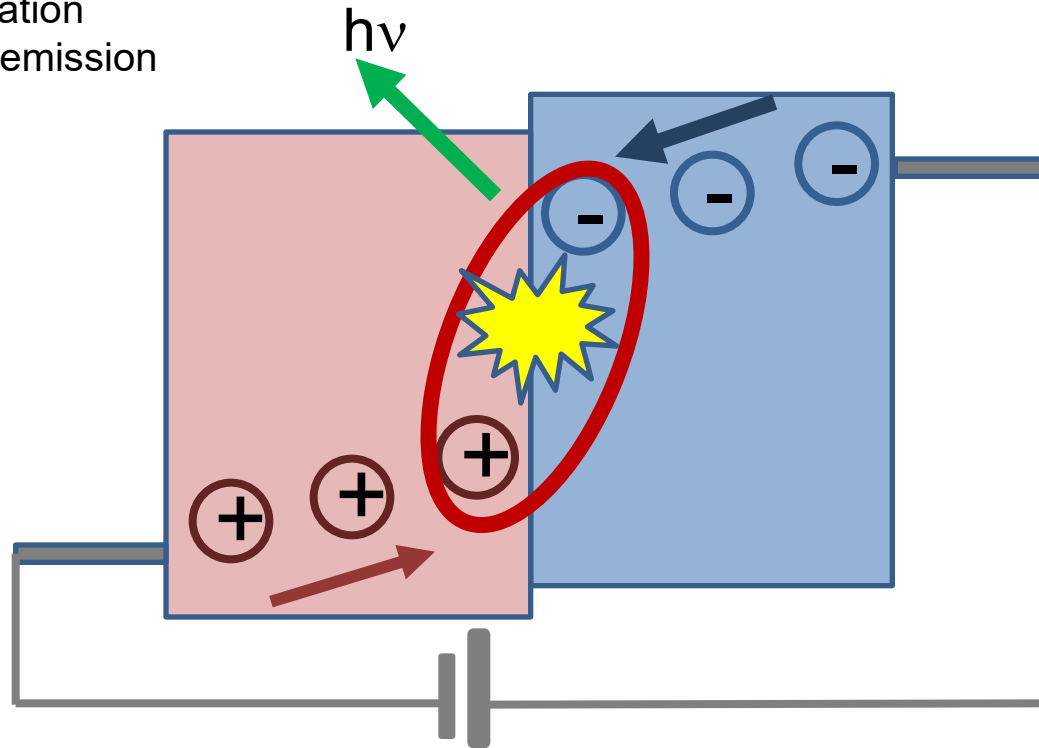
Current market (\$38B in 2021, \$73B in 2026):

- Displays (TVs, phones, watches, wearables, automotive)
- Lighting (low power consumption, brightness, color quality)

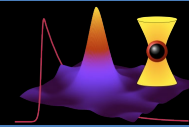
Samsung, LG, Sony, Apple, Nokia, Panasonic,...



1. Charge injection from electrodes
2. Charge (electron and hole) transport
3. Exciton (electron-hole) formation
4. Recombination with photon emission



- Need: molecules able to transport electrons and holes, as well as emit light of a particular frequency (wavelength)



Power conversion efficiency  
(% of solar energy converted into  
usable electricity):

Cost effectiveness for commercialization: ~10-15%

Single-junction Si: ~25%

Multi-junction Si: ~45%

**Organic solar cells:**

2000: 1%

2024: 20%

**2025: 26% (organic-perovskites)**

**Potentially:**

Combine Si with organic to boost to 35-38% in single junction !

Current market: \$100M in 2021, \$800M in 2030

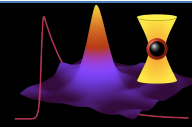


	<b>514 m<sup>2</sup></b> Total area
	<b>22.5 kWp</b> Total capacity
	<b>378</b> Total films

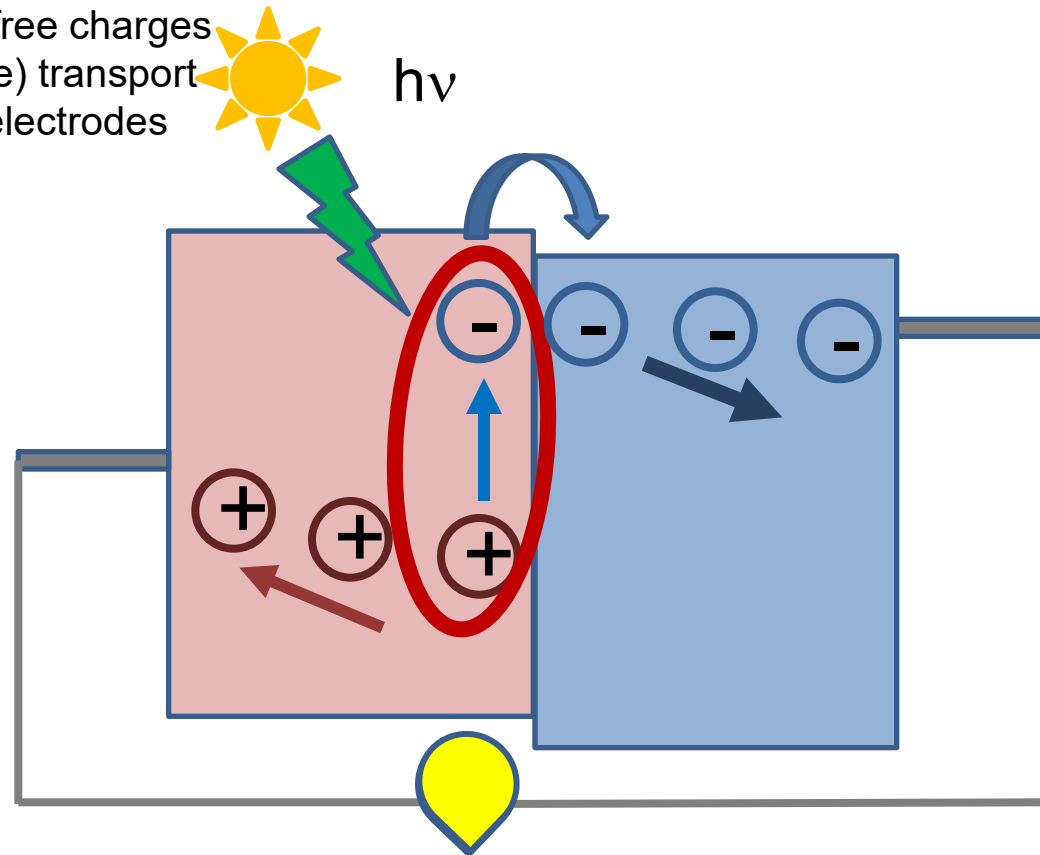
**La Rochelle, France**  
Installation location



<http://www.heliatek.com/>

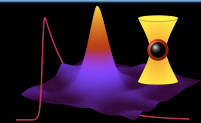


1. Sunlight photon absorption
2. Exciton (electron-hole) formation
3. Exciton dissociation into free charges
4. Charge (electron and hole) transport
5. Charge collection at the electrodes



- Need: molecules able to absorb sunlight, separate the electron-hole pair, and transport electrons and holes





Blue EQE > 40%  
UV EQE > 10%  
Near-IR EQE 10%

2025:  
Hall-effect  $\mu > 100 \text{ cm}^2/(\text{Vs})$

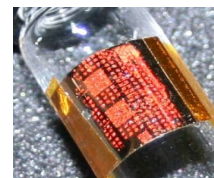
Vasilopoulou et al., *Nat. Comm.* **12**, 4868 (2021)  
Zhang et al., *Ang. Chem. Int. Ed.* **60**, 22241 (2021)  
Shahalizad et al., *Adv. Funct. Mat.* **31**, 2007119 (2021)



Ren et al., *Adv. Sci.* **8**, 2002418 (2021)

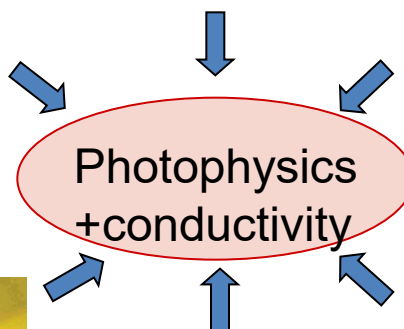
Photodetectors

OTFTs

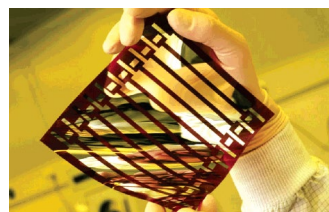


$\mu > 10 \text{ cm}^2/(\text{Vs})$

Paterson et al., *Adv. Mat.* **30**, 1801079 (2018)



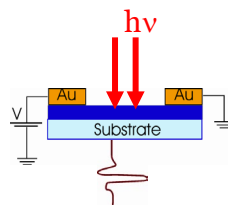
Solar cells



PCE ~ 20%

Gao et al., *Adv. Mat.* **34**, 2202089 (2022)

THz generation



Shin et al., *Adv. Sci.* **7**, 2001738 (2020)

Photorefractive devices



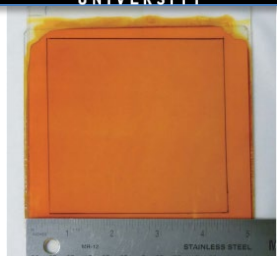
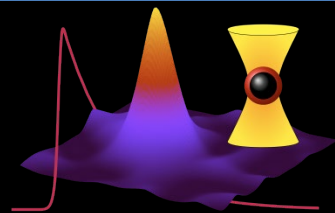
Blanche et al., *Materials* **14**, 5799 (2021)

**Other applications:** RFID tags, batteries, conductive ink, memory devices, sensors, smart cards, toys/games, ...

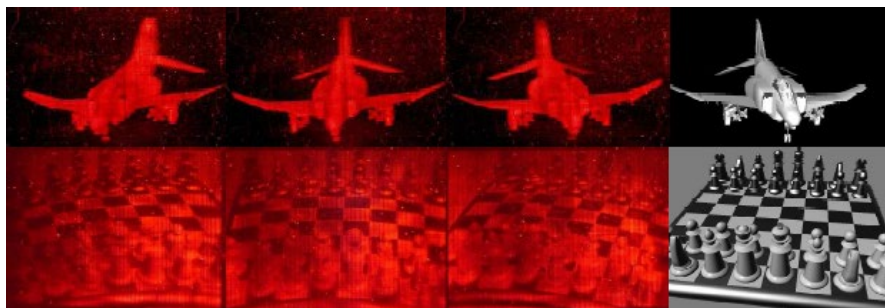
**Need to understand:** how molecules absorb, emit, and convert light into electric charge, conduct charge



# 3D Holographic Displays

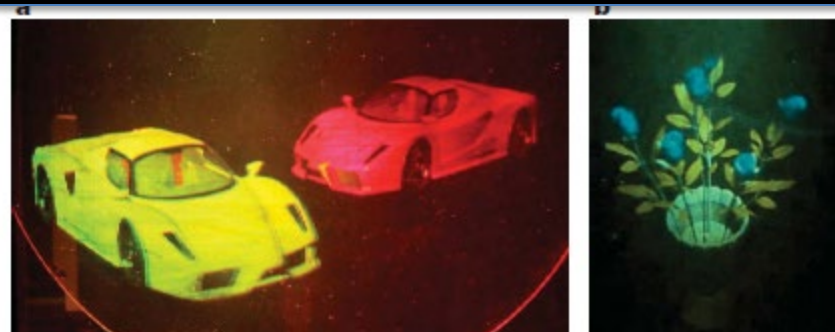


Photorefractive polymer film  
100  $\mu\text{m}$  thick

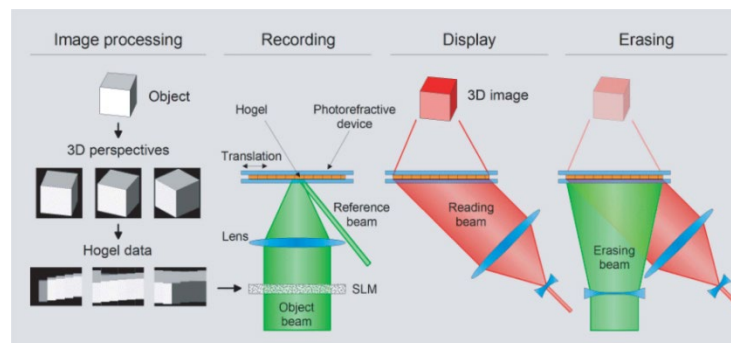


Tay et al., *Nature* **451**, 694 (2008)

P. Blanche et al., *Nature* **468**, 80  
(2010)



**Figure 3** | Pictures of coloured holograms. **a**, Hologram of two model cars recorded on a 12-inch-diameter photorefractive device in HPO geometry. **b**, Hologram of a vase and flowers.

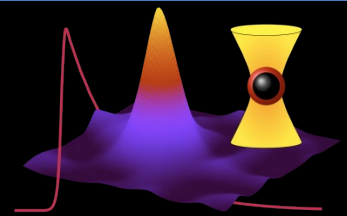


O. Ostroverkhova and W.E. Moerner, *Chem. Rev.* **104**, 3267 (2004)

**Tutorial:** O. Ostroverkhova, “Organic and polymeric photorefractive materials and devices,” in “*Introduction to organic electronic and optoelectronic materials and devices*” (S. Sun and L. Dalton, Eds.), CRC Press, 2008;  
[http://www.physics.oregonstate.edu/~ostroveo/research/resources/tutorials/Photorefractives tutorial/index.html](http://www.physics.oregonstate.edu/~ostroveo/research/resources/tutorials/Photorefractives%20tutorial/index.html)

# Optical properties

# Light-matter interactions: two-level system



state 2 —————  $E_2$        $\Psi_2(r, t) = \psi_2(r) \exp(-i \frac{E_2}{\hbar} t)$

state 1 —————  $E_1$        $\Psi_1(r, t) = \psi_1(r) \exp(-i \frac{E_1}{\hbar} t)$

Transition rate:

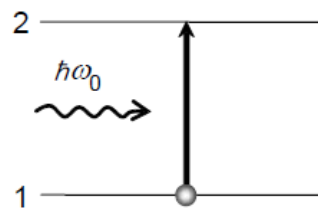
$$\Gamma_{12} = \frac{\pi}{\epsilon_0 \hbar^2} |\mu_{12}|^2 W(\omega_0)$$

Transition dipole moment:

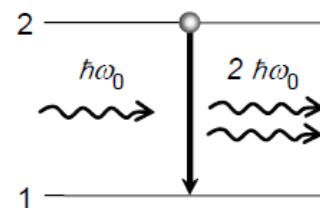
$$|\mu_{21}| = e \int \psi_1^* x \psi_2 dr$$

$$\Gamma_{12} = B_{12} W(\omega_0)$$

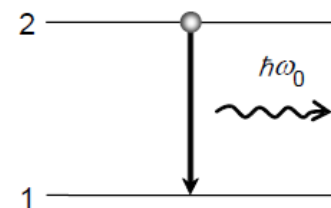
$$A = \frac{\hbar \omega^3}{\pi^2 c^3} B_{12}$$



absorption



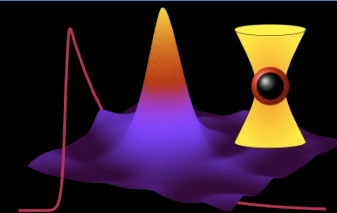
stimulated emission



spontaneous emission

A, B – Einstein coefficients

# Molecular absorption: particle-in-the box

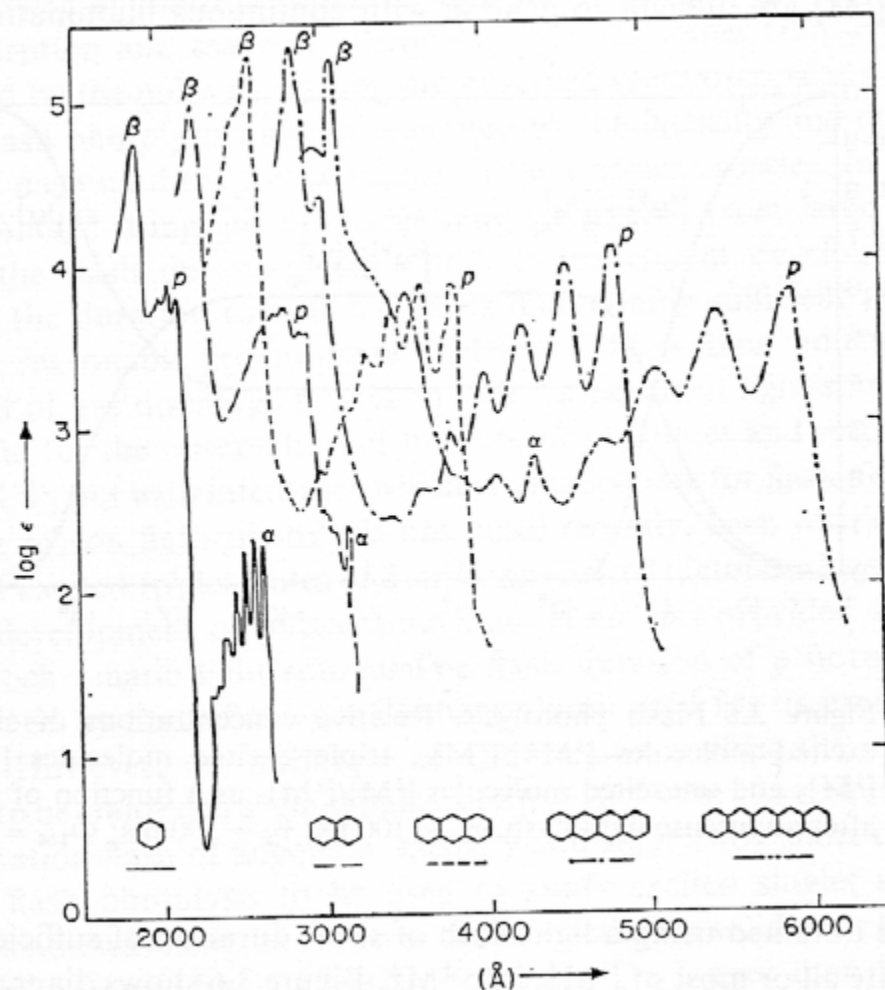


For conjugated molecules:

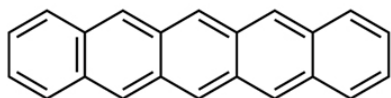
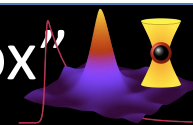
PFEO - perimeter free electron orbital model

$$E_n = \frac{n^2 h^2}{2ml^2}$$

$l$  - perimeter

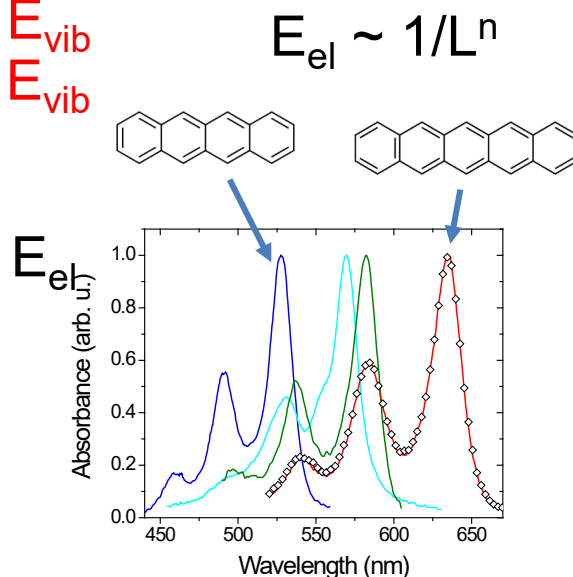
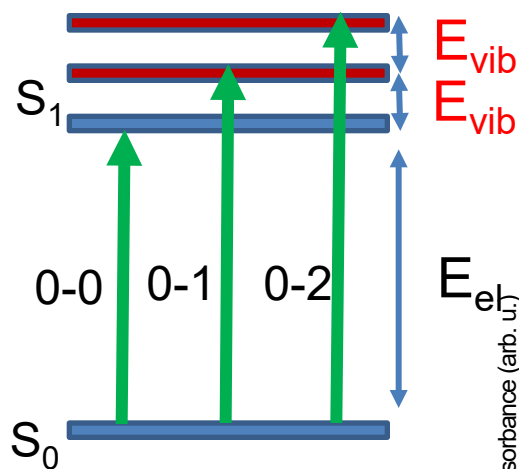


- Larger molecule – more red-shifted spectrum
- $E_{\min} \sim 1/l$

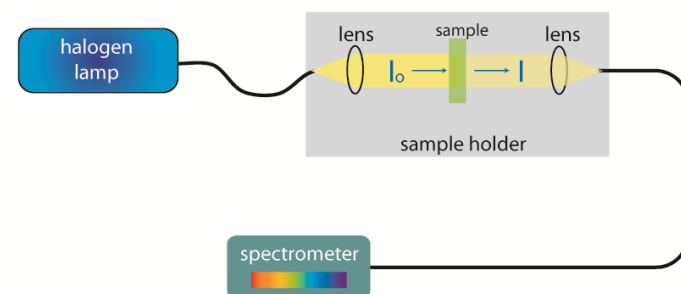


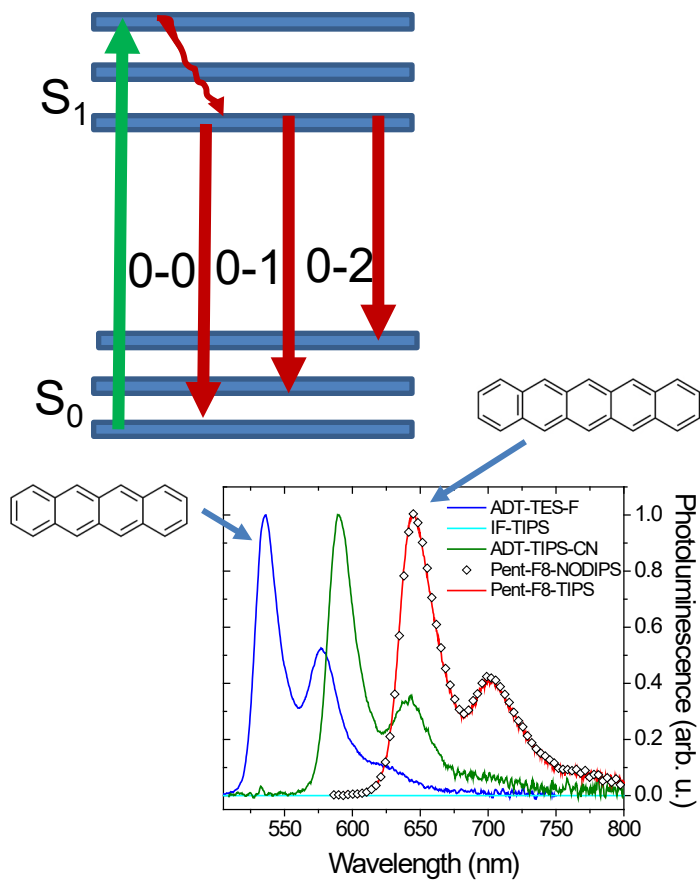
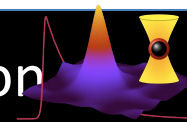
- Need molecules with conjugated core (alternating single/double bonds)
- Electron delocalized over the molecule
- Size of the conjugated core = size of the box
- Larger molecule = lower energies = longer wavelengths

$$\text{Total Energy} = E_{\text{el}} + E_{\text{vib}} + E_{\text{rot}}$$

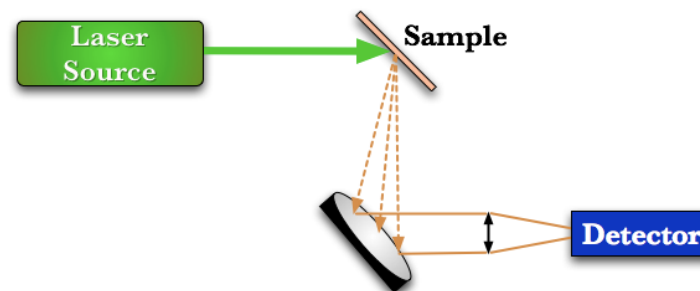


- Optical absorption spectroscopy

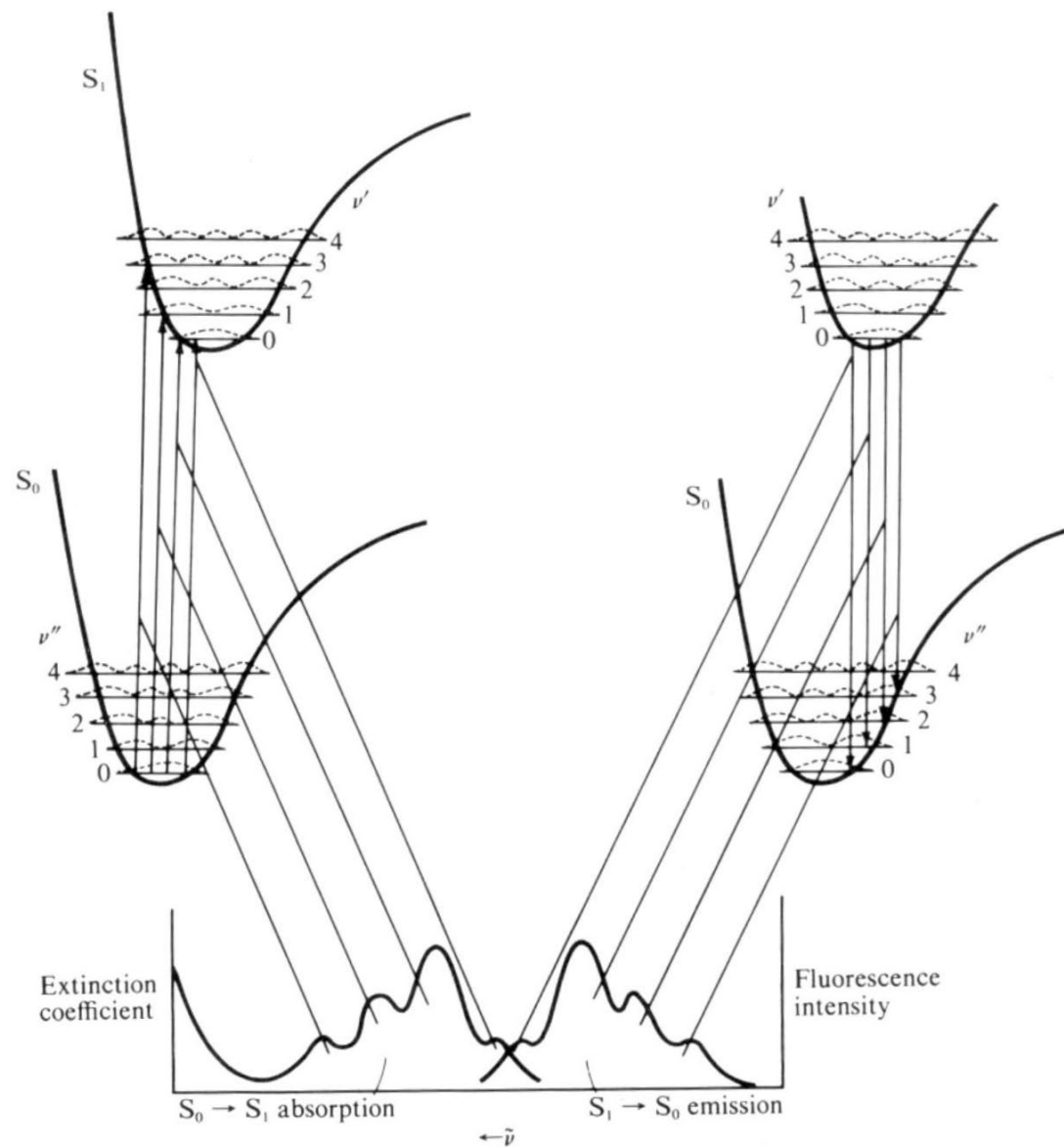




- Fluorescence (or photoluminescence) spectroscopy



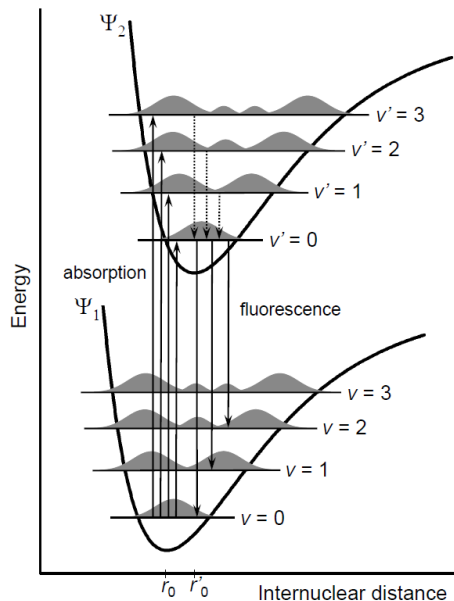
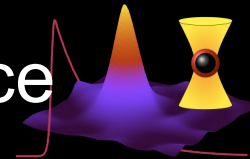
- Measure spectra and timing of each emission event



Simplified potential energy curves with vibrational probability functions showing how a mirror-image relationship can arise between the electronic absorption and emission bands. (From Kearwell and Wilkinson 1969, p. 108)



# Absorption and fluorescence



$$\Psi(\mathbf{r}, \mathbf{R}) = \psi(\mathbf{r}) \chi(\mathbf{R})$$

Electronic part

Nuclear part

- Absorption and emission  $\sim |\mu_{12}|^2$
- Absorption and emission are mirror images
- Vibronic progression

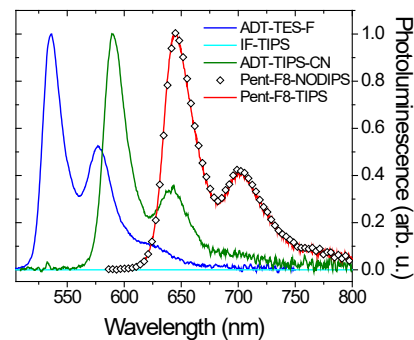
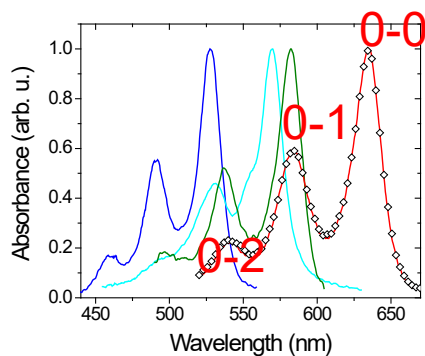
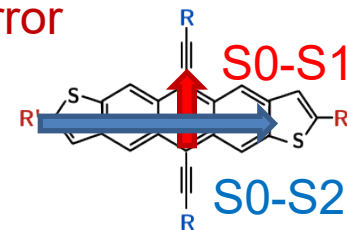
Transition dipole moment:

Absorption

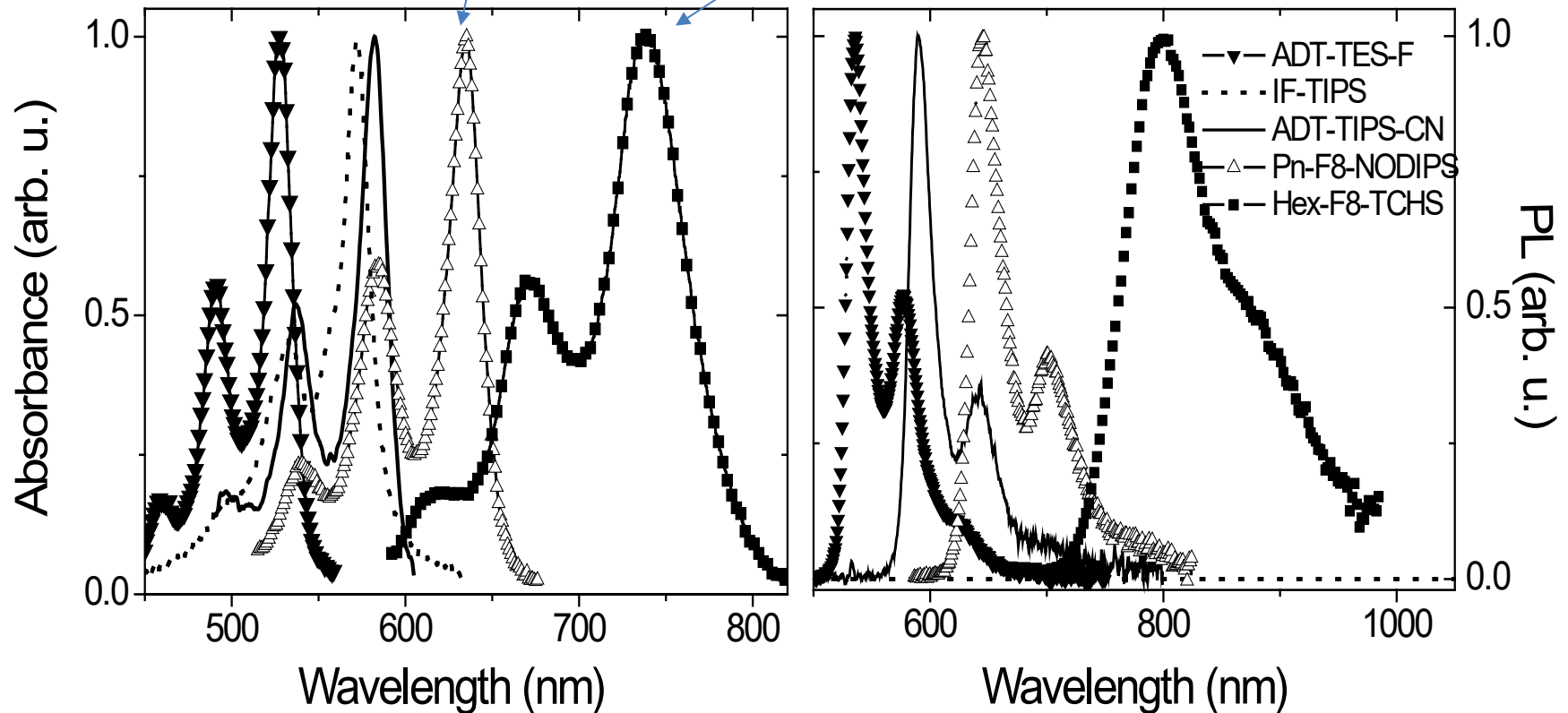
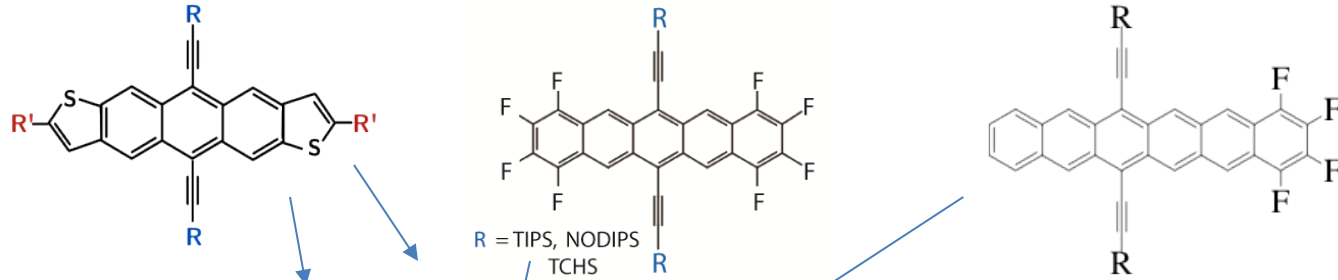
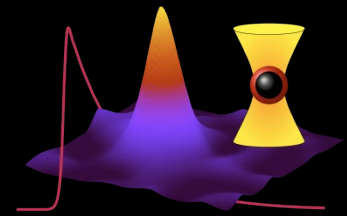
Emission

$$\mu_{12} = \int \chi_{2v'}^* \chi_{10} dR \int \psi_2^* \mu \psi_1 dr$$

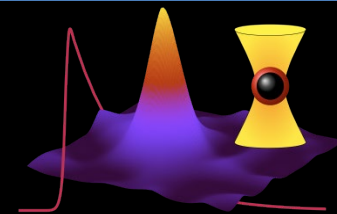
$$\mu_{21} = \int \chi_{1v}^* \chi_{20} dR \int \psi_1^* \mu \psi_2 dr$$



# Optical properties: solutions



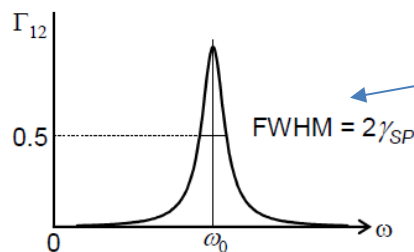
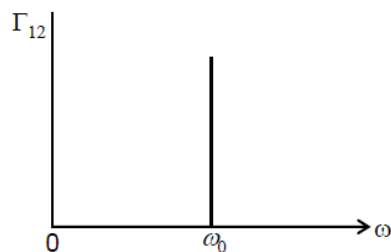
# Linewidth



$$\Gamma_{12} = \frac{\pi}{2\hbar^2} E_0^2 |\mu_{12}|^2 \delta(\omega_0 - \omega)$$



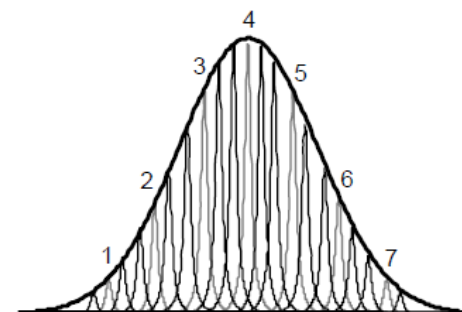
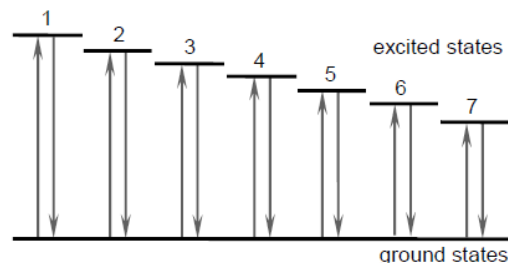
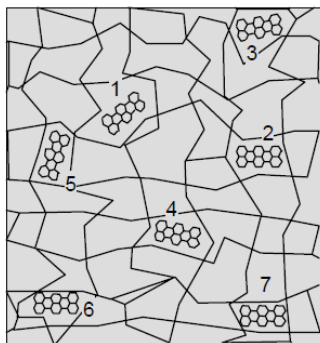
$$\Gamma_{12} = \frac{\pi}{2\hbar^2} E_0^2 |\mu_{12}|^2 \frac{\gamma_{SP} / \pi}{(\omega_0 - \omega)^2 + \gamma_{SP}^2}$$



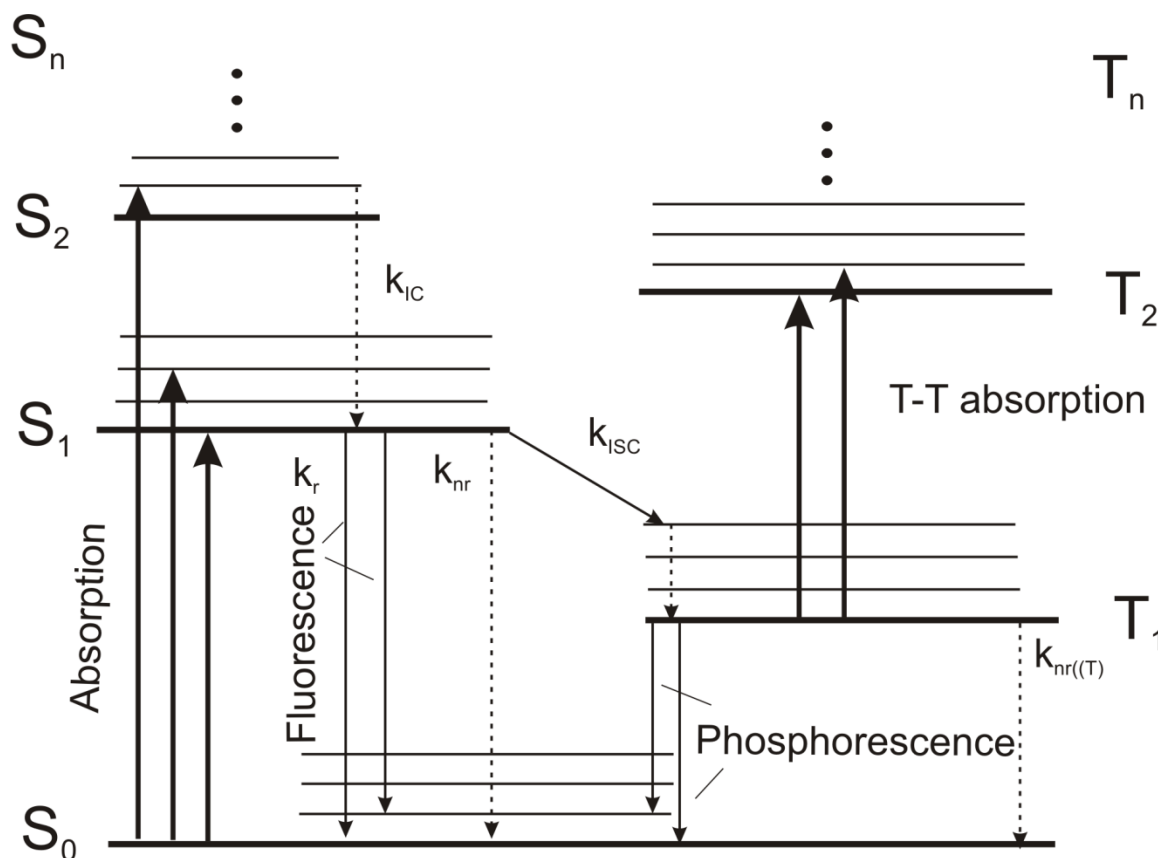
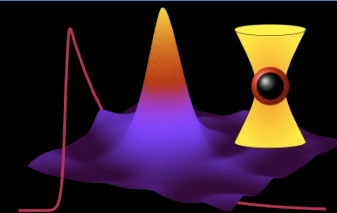
Lorentzian

- Homogeneous broadening: Lorentzian (interactions with phonons)
- Inhomogeneous broadening: Gaussian (mechanical strain, electrostatic interactions)

Inhomogeneous broadening



# Jablonski diagram



Energy:

$$E = E_e + E_{vib} + E_{rot}$$

$$\psi_i = \psi_i(r) \chi_i(\sigma)$$

↑  
Spin

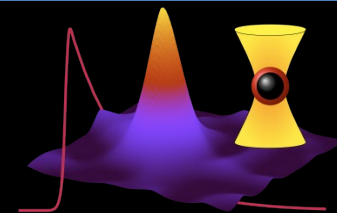
$S_i$  – singlet states  
(spin  $S = 0$ )

$T_i$  – triplet states  
(spin  $S = 1$ )

$$\mu_{12} = \int \psi_2^*(r) \boldsymbol{\mu} \psi_1(r) dr \int \chi_2^*(\sigma) \chi_1(\sigma) d\sigma$$

← Zero for S-T transitions  
(forbidden transitions)

# Franck-Condon principle



$$\Psi(\mathbf{r}, \mathbf{R}) = \psi(\mathbf{r}) \chi(\mathbf{R})$$

Electronic part

Nuclear part

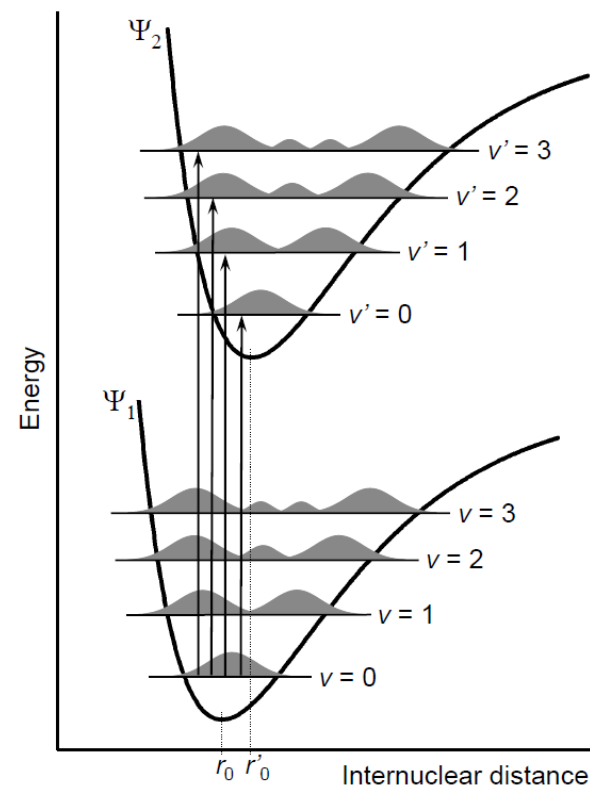
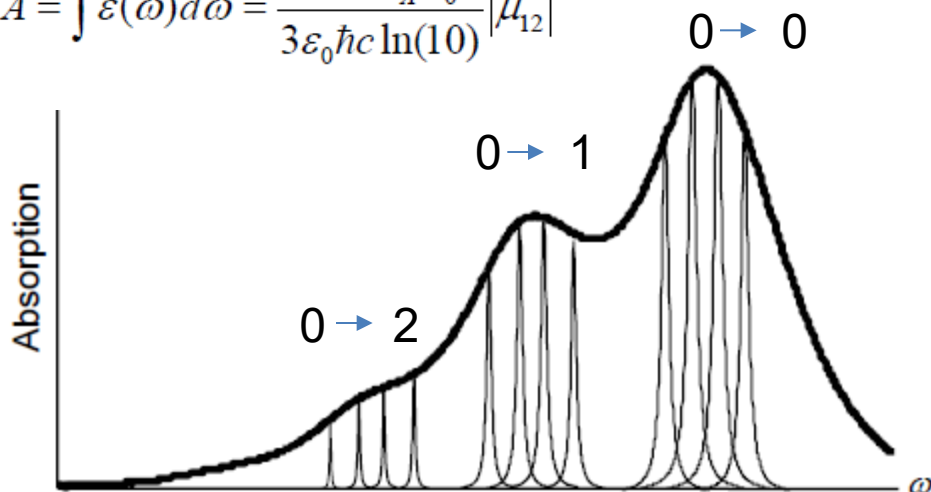
$$\mu_{12} = \int \chi_2^* \chi_1 dR \int \psi_2^* \mu \psi_1 d\mathbf{r}$$

$$\longrightarrow S(\chi_2, \chi_1) = \left| \int \chi_2^* \chi_1 dR \right|^2$$

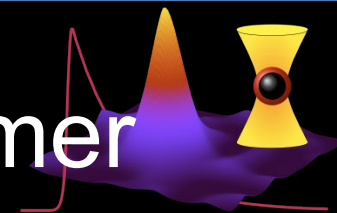
Frank-Condon factor

- Electronic transitions much faster than nuclear vibrations - most probable transition occurs between states with the same nuclear coordinates (vertical trans

$$A = \int \varepsilon(\omega) d\omega = \frac{2\pi^2 N_A \omega_0}{3\varepsilon_0 \hbar c \ln(10)} |\mu_{12}|^2$$



# Two interacting molecules: dimer



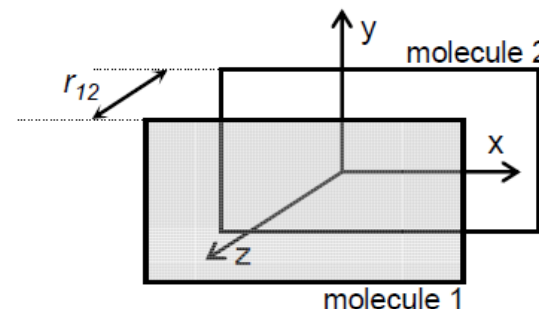
$$\Psi_G = \psi_1 \psi_2$$

Ground state wavefunction

$$H = H_1 + H_2 + V_{12}$$



Intermolecular interaction



$$E_G = \iint \psi_1 \psi_2 H \psi_1 \psi_2 dr_1 dr_2 = E_1 + E_2 + \iint \psi_1 \psi_2 V_{12} \psi_1 \psi_2 dr_1 dr_2 = E_1 + E_2 + D_G$$



Ground state energy

- Shift in the ground state energy due to intermolecular interactions

$$\Psi_E = a \psi_1^u \psi_2 + b \psi_1 \psi_2^u$$

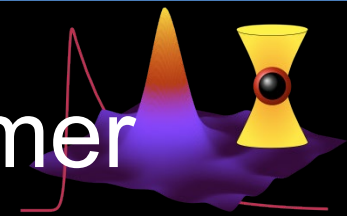
Excited state wavefunction

$$H \Psi_E = E_E \Psi_E \quad \text{or} \quad H(a \psi_1^u \psi_2 + b \psi_1 \psi_2^u) = E_E (a \psi_1^u \psi_2 + b \psi_1 \psi_2^u)$$



Excited state energy

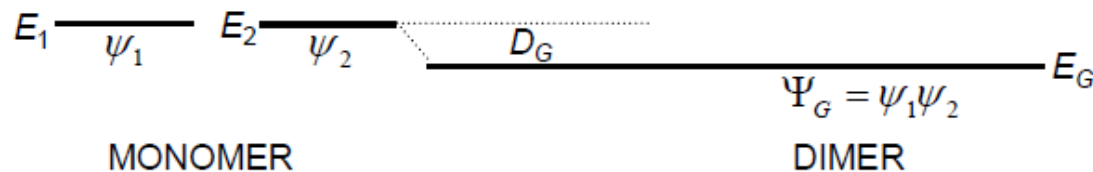
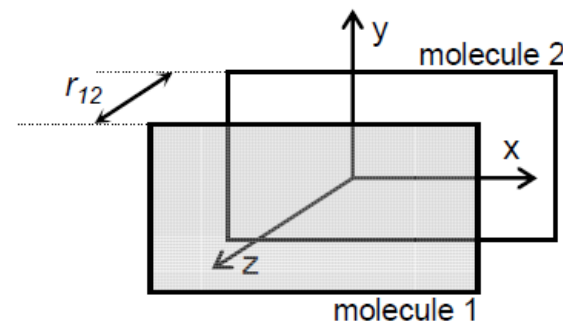
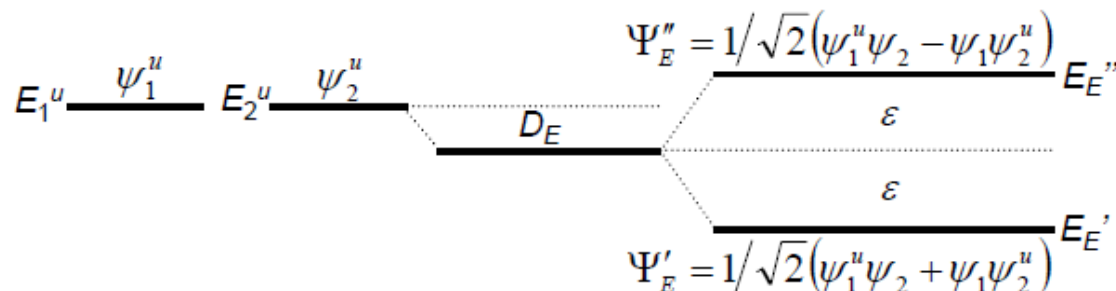
# Two interacting molecules: dimer



$$E'_E = E_1^u + E_2 + D_E + \varepsilon$$

$$V_{12} = \frac{e^2}{4\pi\epsilon_0 r_{12}^3} \sum_{i,j} x_1^i x_2^j \quad \longrightarrow \quad \varepsilon = \frac{\mu_1 \mu_2}{4\pi\epsilon_0 r_{12}^3}$$

$$E''_E = E_1^u + E_2 + D_E - \varepsilon$$

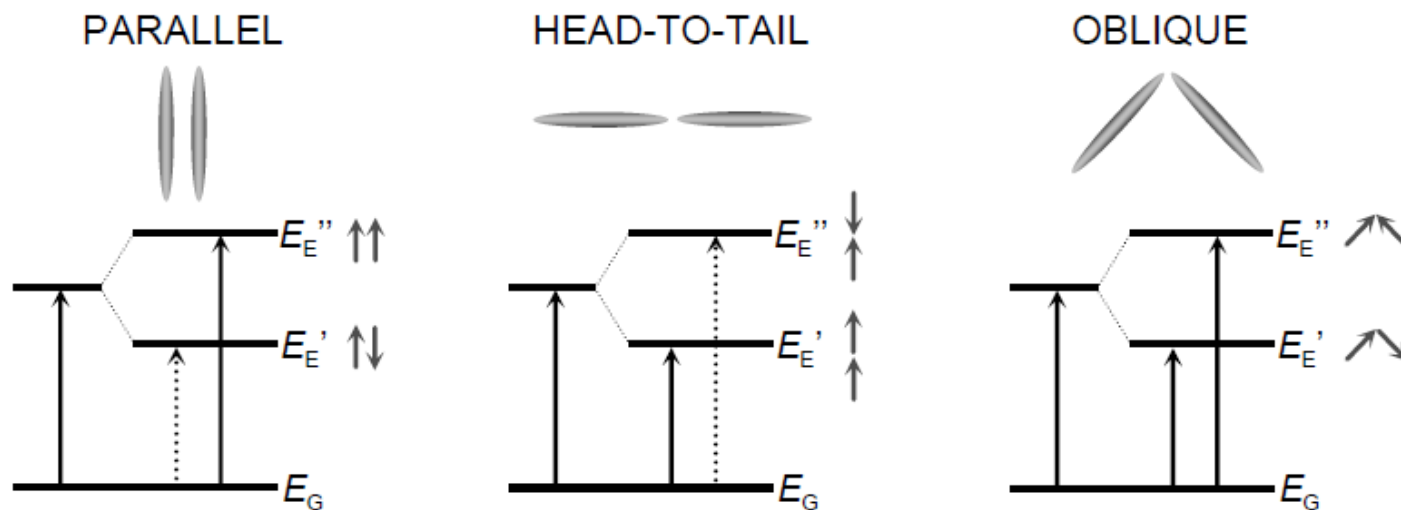
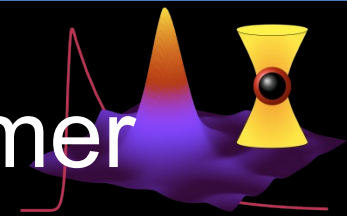


$$E_E - E_G = \Delta E = \Delta E_{monomer} + \Delta D \pm \varepsilon$$

- Shift and splitting of the excited state energy and of energy of transitions due to intermolecular interactions

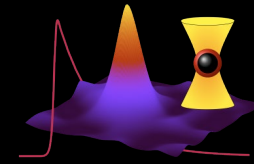


# Two interacting molecules: dimer



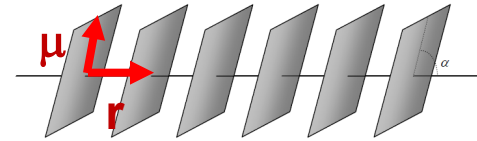
- Allowed or disallowed transitions depending on the orientation of molecular transition dipole moments

# Many interacting molecules: aggregates



$$\Psi_G = \psi_1 \psi_2 \psi_3 \dots \psi_N = \prod_{n=1}^N \psi_n \quad \leftarrow \text{Ground state}$$

$$\Phi_a = \psi_1 \psi_2 \psi_3 \dots \psi_a^u \dots \psi_N = \psi_a^u \prod_{\substack{n=1 \\ n \neq a}}^N \psi_n \quad \leftarrow \text{Excited state}$$



$$\Psi_E^k = \sum_{a=1}^N C_{ak} \Phi_a$$

$$\varepsilon_{a,a+1} = \int \Phi_a V_{a,a+1} \Phi_{a+1} dr$$

$$E_E^k = E_{E,a} + 2 \left( \frac{N-1}{N} \right) \cos \left( \frac{2\pi k}{N} \right) \varepsilon_{a,a+1}$$

Nearest-neighbor interactions

$$V_{12} = \frac{e^2}{4\pi\epsilon_0 r_{12}^3} \sum_{i,j} x_1^i x_2^j$$

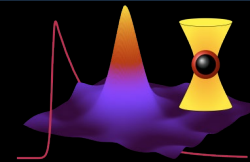
$$\text{Energy: } E_E^k = E_{E,a} + \overbrace{2 \left( \frac{N-1}{N} \right) \cos \left( \frac{2\pi k}{N} \right) \left( \frac{\mu^2}{4\pi\epsilon_0 r^3} \right) (1 - 3 \cos^2 \alpha)}^{J_0}$$

- Delocalized Frenkel excitons as a result of intermolecular coupling
- Properties depend on the transition dipole moment and the angle between  $\mu$  and  $r$
- $J_0$  can be  $<0$  (J-aggregates) or  $>0$  (H-aggregates)



- Shifts in the spectra
- Line broadening
- Intensity redistribution

# Many interacting molecules: charge transfer excitons

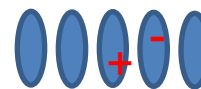


$$\Psi_G = \psi_1 \psi_2 \psi_3 \dots \psi_N = \prod_{n=1}^N \psi_n \quad \leftarrow \text{Ground state}$$

$$\Phi_a = \psi_1 \psi_2 \psi_3 \dots \psi_a^u \dots \psi_N = \psi_a^u \prod_{n \neq a}^N \psi_n \quad \leftarrow \text{Frenkel excitons (comprise neutral states)}$$



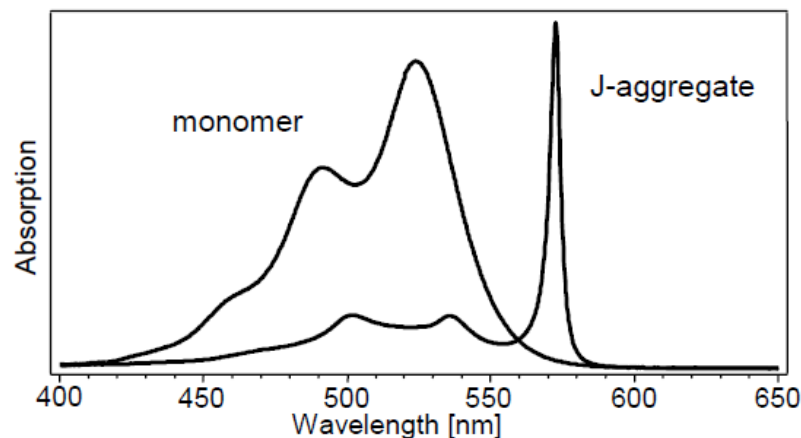
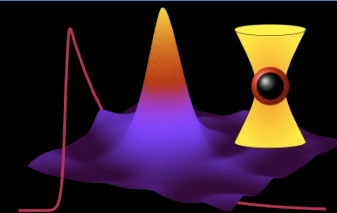
$$\Phi_{CT} = \psi_a^+ \psi_{a+1}^- \prod_{\substack{n=1 \\ n \neq a, a+1}}^N \psi_n \quad \leftarrow \text{CT excitons (comprise ionic states)}$$



$$\Psi_E^k = \sum_{a=1}^N C_{ak} \Phi_a \quad (\text{or } \Phi_{CT})$$

- Delocalized Frenkel and CT excitons also interact with each other
- Such interaction determines important properties: singlet fission, charge generation, emission, etc.
- Subject of debate
- Systematic studies in molecular crystals are needed

# H- and J-aggregates

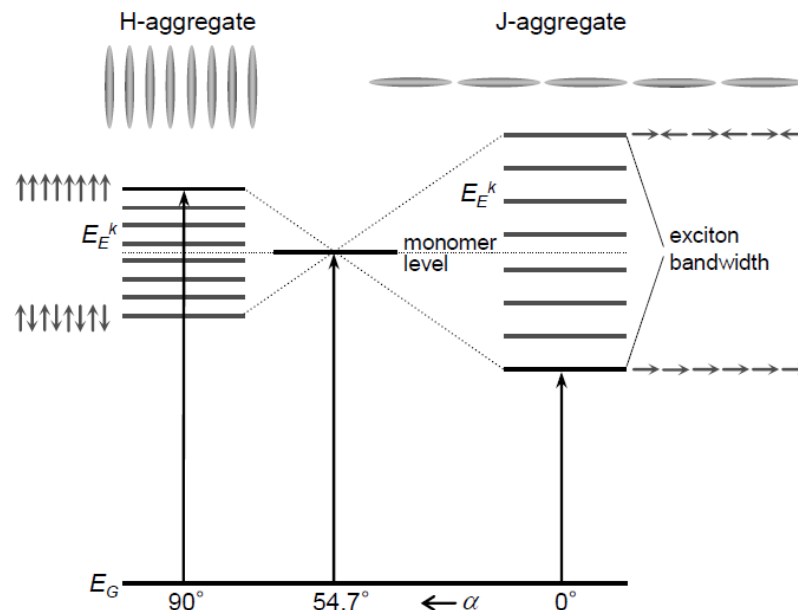


## J-aggregate:

- red-shifted absorption
- strong 0-0 PL emission,  $I_{0-0}/I_{0-1} \sim L_{\text{coh}}$

## H-aggregate:

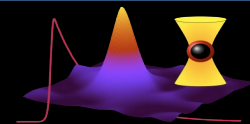
- blue-shifted absorption
- forbidden 0-0 PL emission
- For both, mirror symmetry is not expected anymore



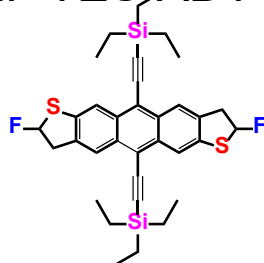
$$\text{bandwidth} = 4 \left( \frac{N-1}{N} \right) \cos \left( \frac{2\pi k}{N} \right) \epsilon_{a,a+1}$$

## Realistic systems:

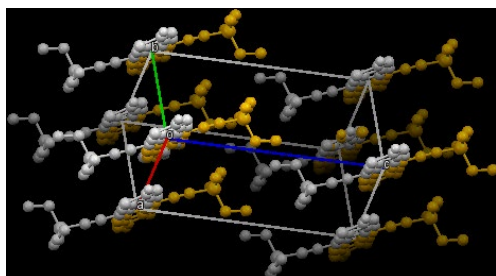
- not easy to assign
- can be a H-J mixture (e.g.  $\pi$ -stacked polymers)



diF TES-ADT



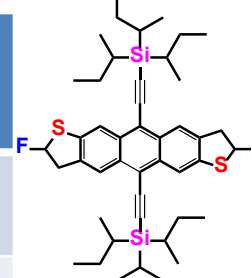
2D “brick-work



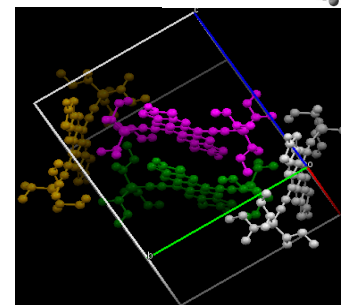
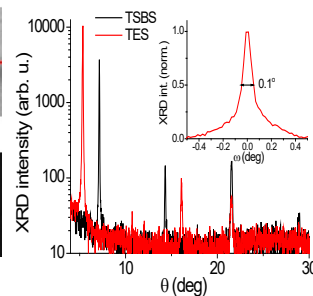
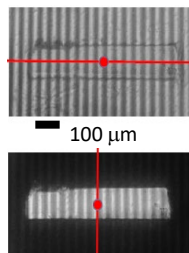
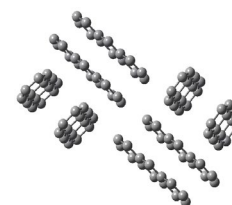
High-performance organic semiconductors

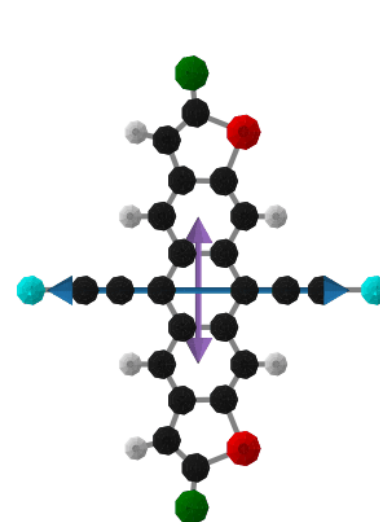
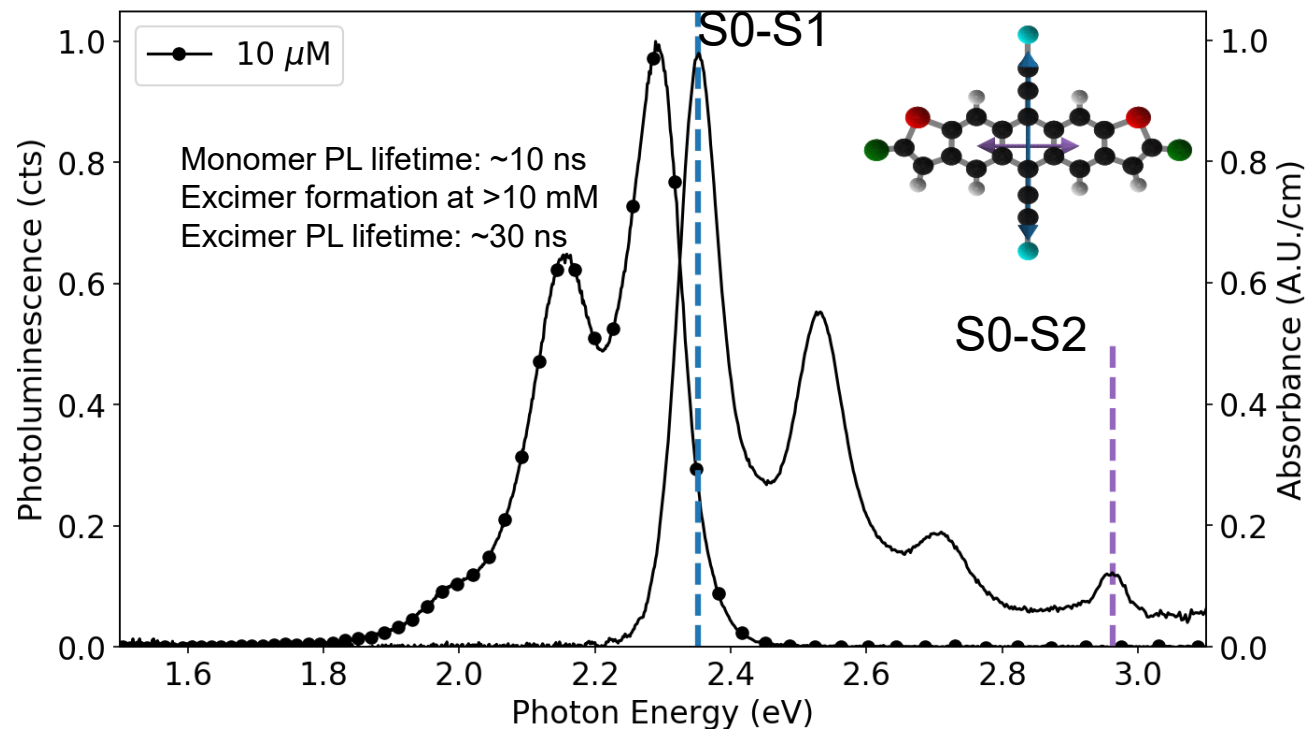
	a, Å	b, Å	c, Å	$\alpha$ , deg	$\beta$ , deg	$\gamma$ , deg
diF TES-ADT	7.1	7.2	16.6	97.5	91.4	107.5
diF TSBS-ADT	15.1	16.4	18.2	90	103	90

diF TSBS-ADT

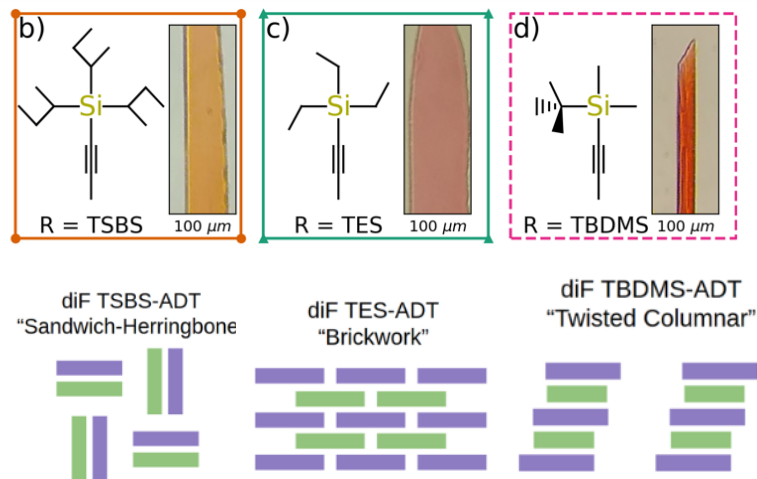
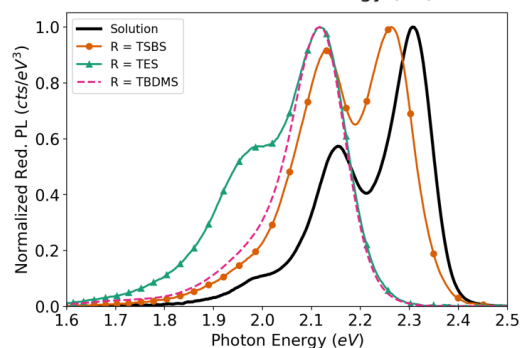
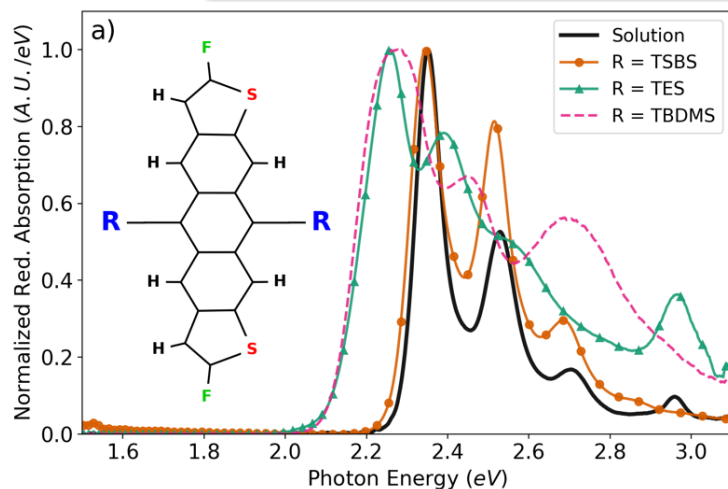
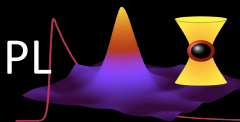


1D “sandwich-herringbone





- Absorption and photoluminescence spectra behave as expected from “isolated” molecules



- Absorption and PL shifts and oscillator strength redistribution due to intermolecular interactions in crystals
- Absorption strongly polarization-dependent
- Analyze polarization and temperature dependence, in conjunction with TD-DFT to understand intermolecular interactions



- Frenkel/CT exciton mixing
- Packing-dependent exciton nature

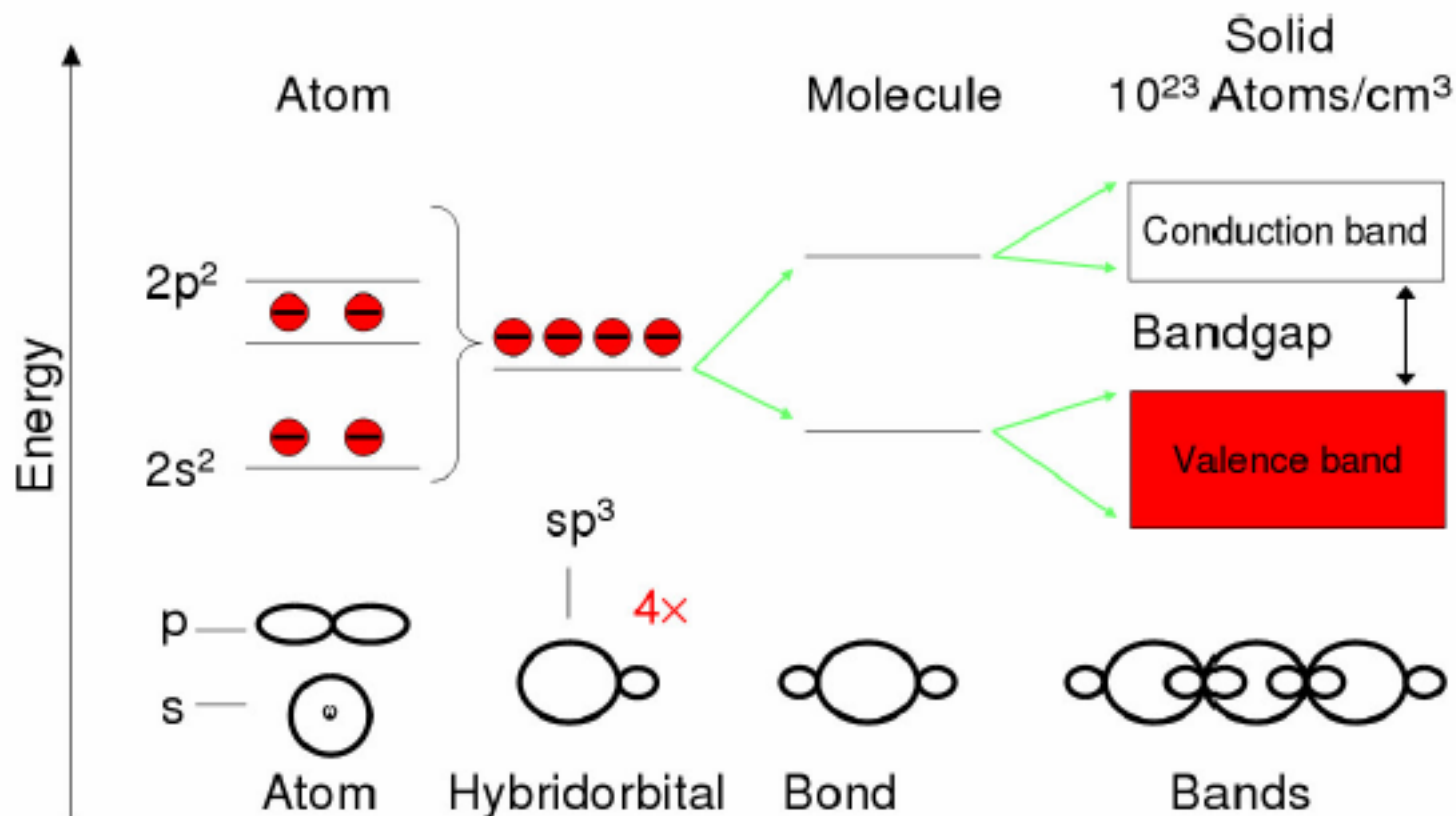
J. Van Schenck, G. Mayonado, J. E. Anthony, M. Graham, O. Ostroverkhova, "Molecular packing-dependent exciton dynamics in functionalized anthradithiophene derivatives: from solutions to crystals", *J. Chem. Phys.* **153**, 164715 (2020) (Featured Article)



# Organic vs inorganic semiconductors: electronic properties

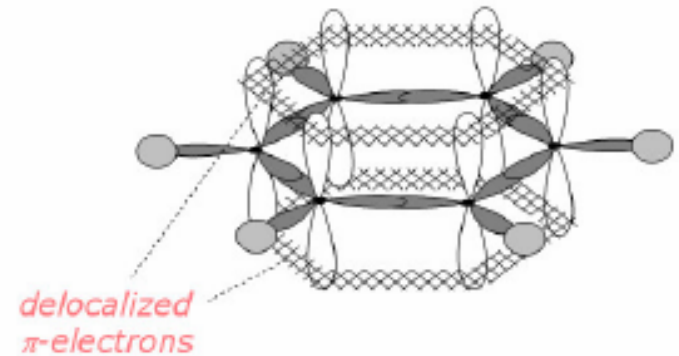
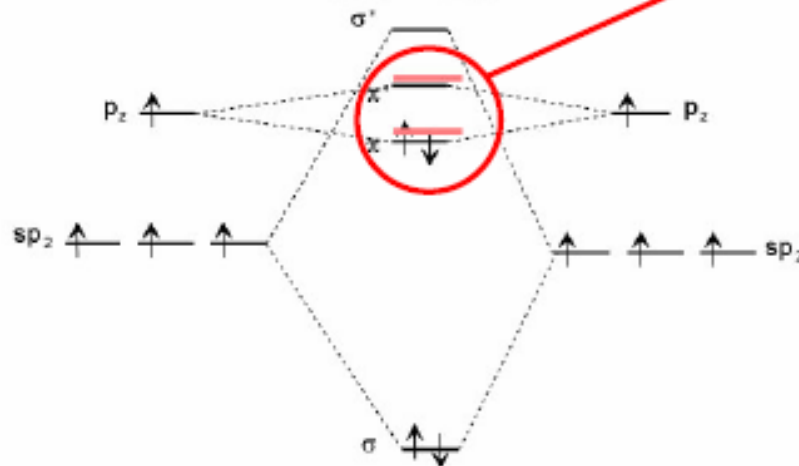
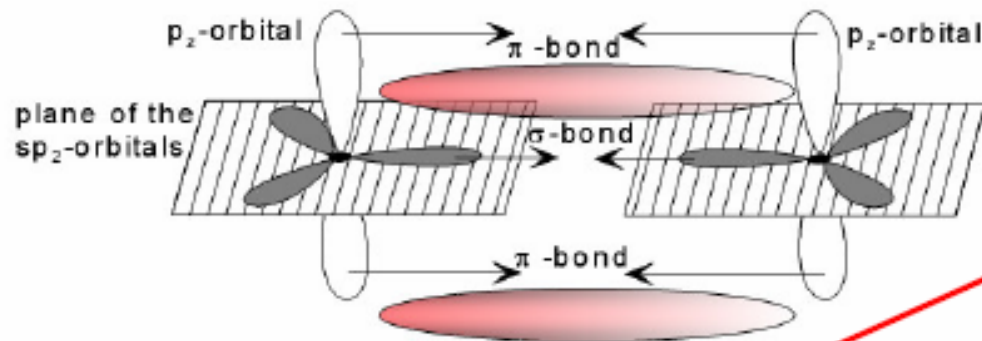
# The carbon as fundamental element of organic materials

From atomic bonds to solid bands through intermolecular connectivity



# The carbon as fundamental element of organic materials

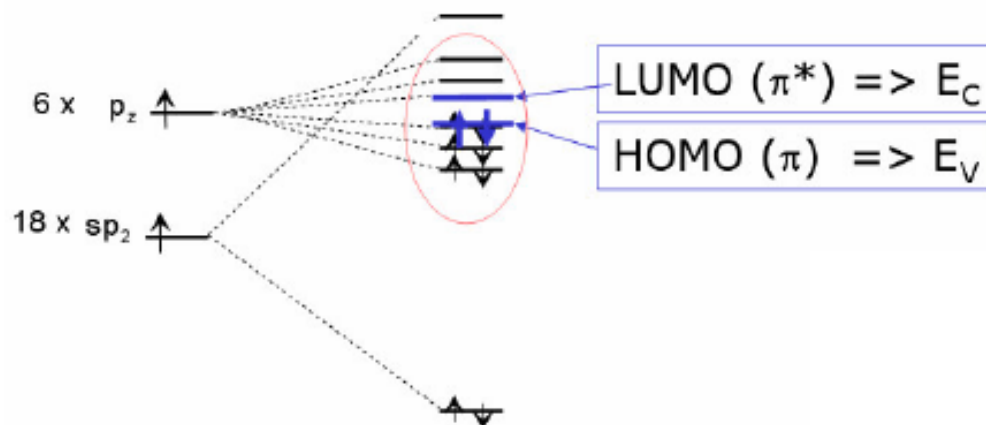
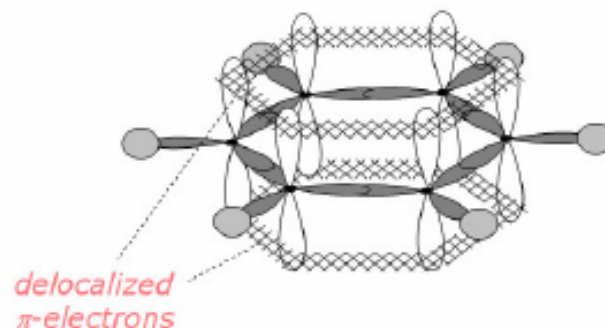
## Sp<sub>2</sub>-hybridised Carbon:



- Spatially extended electron system
- Low electronic binding energies
- Larger polarizability
- Good coupling between molecules in solid state
- Larger reactivity

## The carbon as fundamental element of organic materials

- Saturated carbon – SP3 hybridization. Large energy gaps between filled and empty states.
- Unsaturated carbon – SP2 and SP hybridization.  $\pi$  states closer in energy, and more extended in space.
- Small  $\pi$ - $\pi$  overlap between molecules create narrow bands, favoring electron mobility



- Spatially extended electron system
- Low electronic binding energies
- Larger polarizability
- Good coupling between molecules in solid state
- Larger reactivity

# Non-covalent connectivity

How can neutral and non-polar molecular units form a crystal?  
Non-covalent interactions are responsible of crystal structure, governing the shape and size of the unit cell, as well as the functionality.

Types of interactions between two neutral molecules:

A) Both molecules have dipolar moment: dipolar, H-bonding



Aka Keesom interaction

$$V = \frac{-1}{4\pi\epsilon_0} \frac{2p_1p_2}{r^3}$$

B) One molecule has dipolar moment, the other doesn't: Induction



Aka Debye interaction

$$V = \frac{-1}{(4\pi\epsilon_0)^2} \frac{2p_1^2\alpha}{r^6}$$

C) None of the molecules has a dipolar moment: van der Waals



Aka London interaction

$$V = \frac{-1}{(4\pi\epsilon_0)^2} \frac{A\alpha^2}{r^6}$$



The Nobel Prize in Physics 1910 was awarded to Johannes Diderik van der Waals *"for his work on the equation of state for gases and liquids"*.

## Non-covalent connectivity

Real potential comes from addition of the attractive intermolecular forces and a general repulsive term accounting for steric molecular repulsion: this is repulsion between the nuclei and electron clouds of each molecule

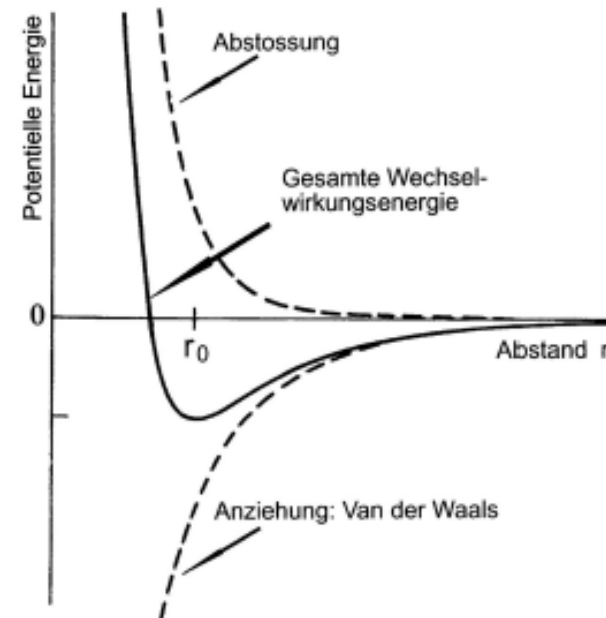
$$V_{rep.} = \frac{C}{r^{12}}$$

### Lennard Jones potential:

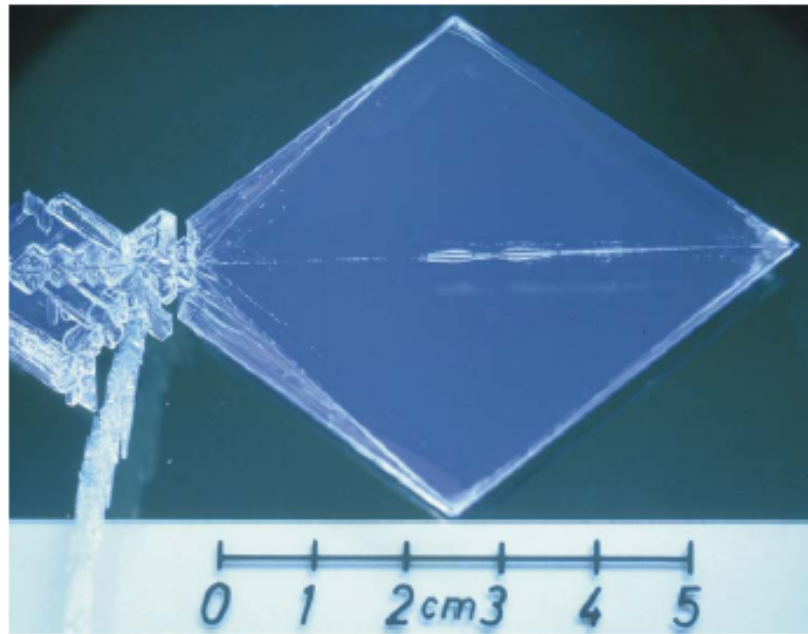
- Inter-molecular distance in the order of 0.3 nm
- Intermolecular bond energy weak, typically less than 10 kcal/mol (0.1 eV).

$$V = \frac{C}{r^{12}} - \frac{D}{r^6}$$

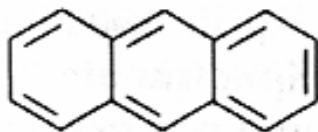
$$V = 0 \quad \text{für} \quad r > r_0, \quad V = \infty \quad \text{für} \quad r \rightarrow 0$$



# Molecular Crystals

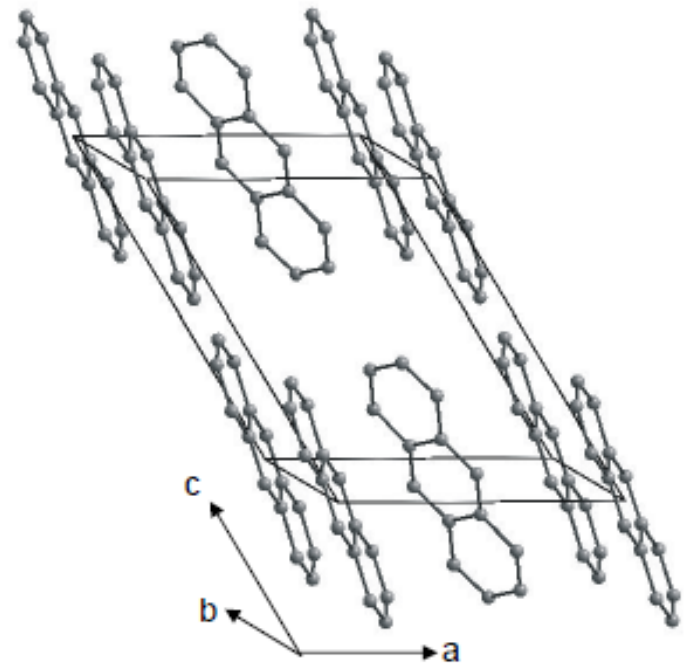


Lattice Constants:  $a \sim 8.6 \text{ \AA}$   
 $b \sim 6.0 \text{ \AA}$   
 $c \sim 11.2 \text{ \AA}$   
 $\beta \sim 125^\circ$



Anthracene  
( $C_{14}H_{10}$ )

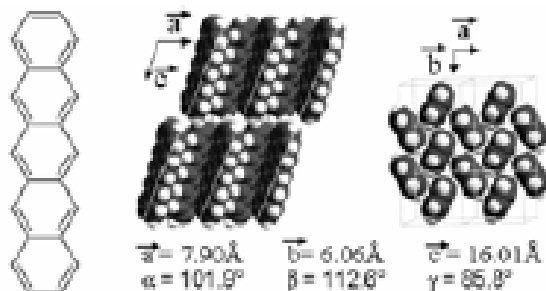
→ van der Waals Bonding





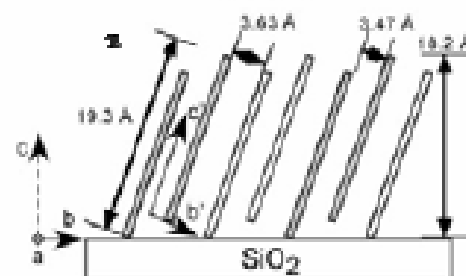
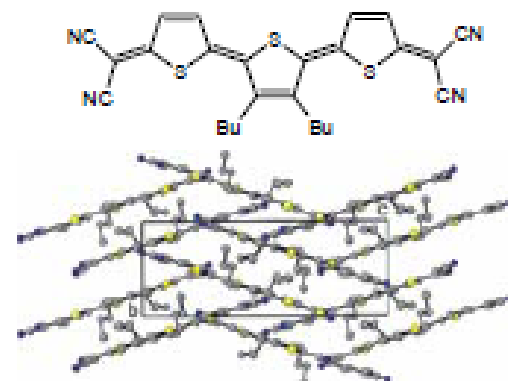
# Organic Semiconductors for Thin Film Transistors

## Pentacene: State of the Art Organic Semiconductor



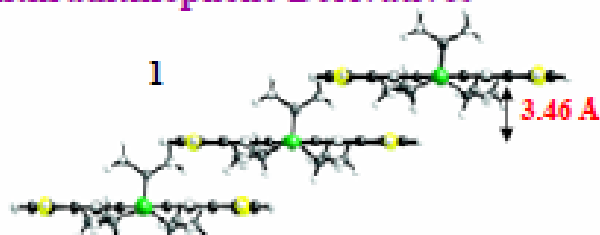
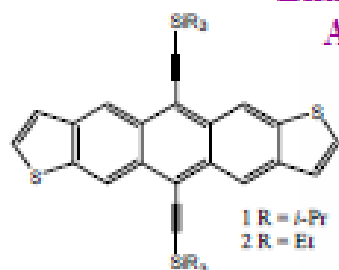
- On flat, inert oxide or polymeric dielectrics, pentacene molecules stand nearly vertical on the substrate
- Highest Hole Mobility in Organic Semiconductors:  $3\text{--}5 \text{ cm}^2/\text{Vs}$

## $\pi$ -Stacking Quinoidal Terthiophene



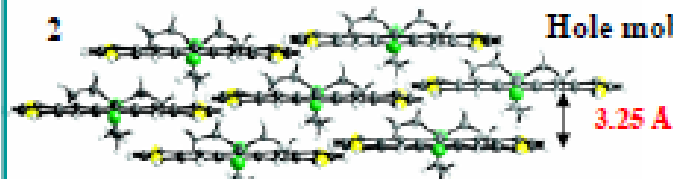
- High Electron Mobility:  $\mu_e \sim 0.2 \text{ cm}^2/\text{Vs}$
- Ambipolar Charge Transport

## Enhanced $\pi$ -Stacking Interactions in Anthradithiophene Derivatives



1-D slipped-stack arrangement

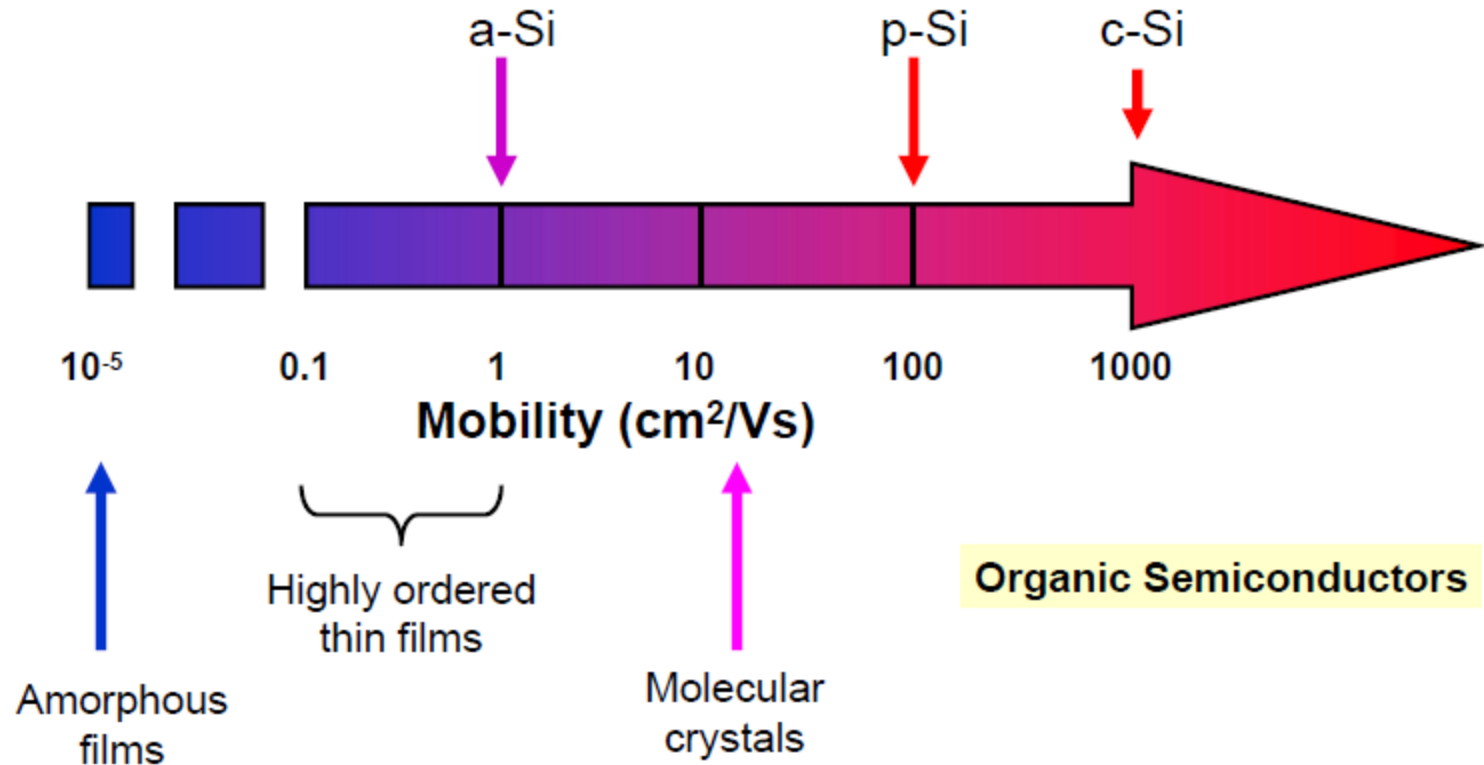
Hole mobility of  $<10^{-4} \text{ cm}^2/\text{Vs}$



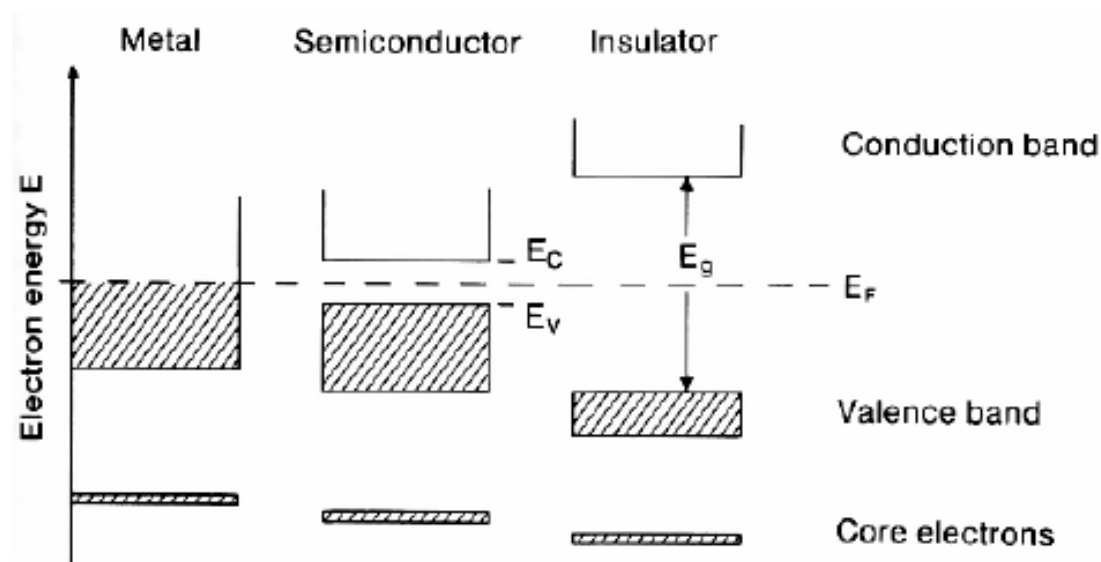
2-D stacking arrangement

Hole mobility of  $1 \text{ cm}^2/\text{Vs}$

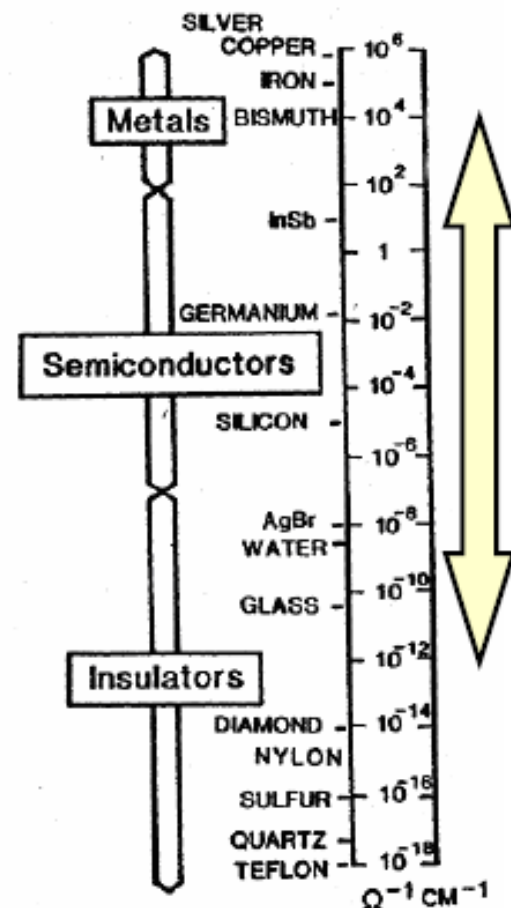
## Charge Carrier Mobility



# Electric Conductivity



Organic Conductors / Semiconductors:  
 $\sigma = 10^{-12} \dots 10^4 \text{ S/cm}$

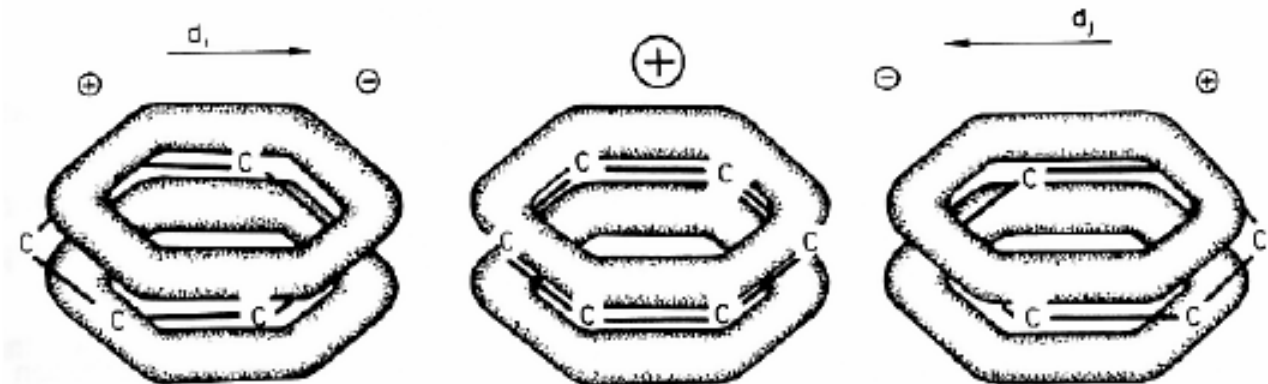


# Charge Carriers in Organic Semiconductors

## Molecular organic semiconductors

- Charge carriers in molecular organic semiconductors: molecule ions  $M^+$  resp.  $M^-$
- Gas phase:  
 $M \rightarrow M^+ + e^-$  (ionisation energy  $I_g$ )  
 $M + e^- \rightarrow M^-$  (electron affinity  $A_g$ )
- Crystal: Stabilisation of molecule ions in the solid state by polarisation energies  $P_e$  and  $P_h$  (cf. solvation energies in solutions)  
 $\rightarrow A_c > A_g$  and  $I_c < I_g$

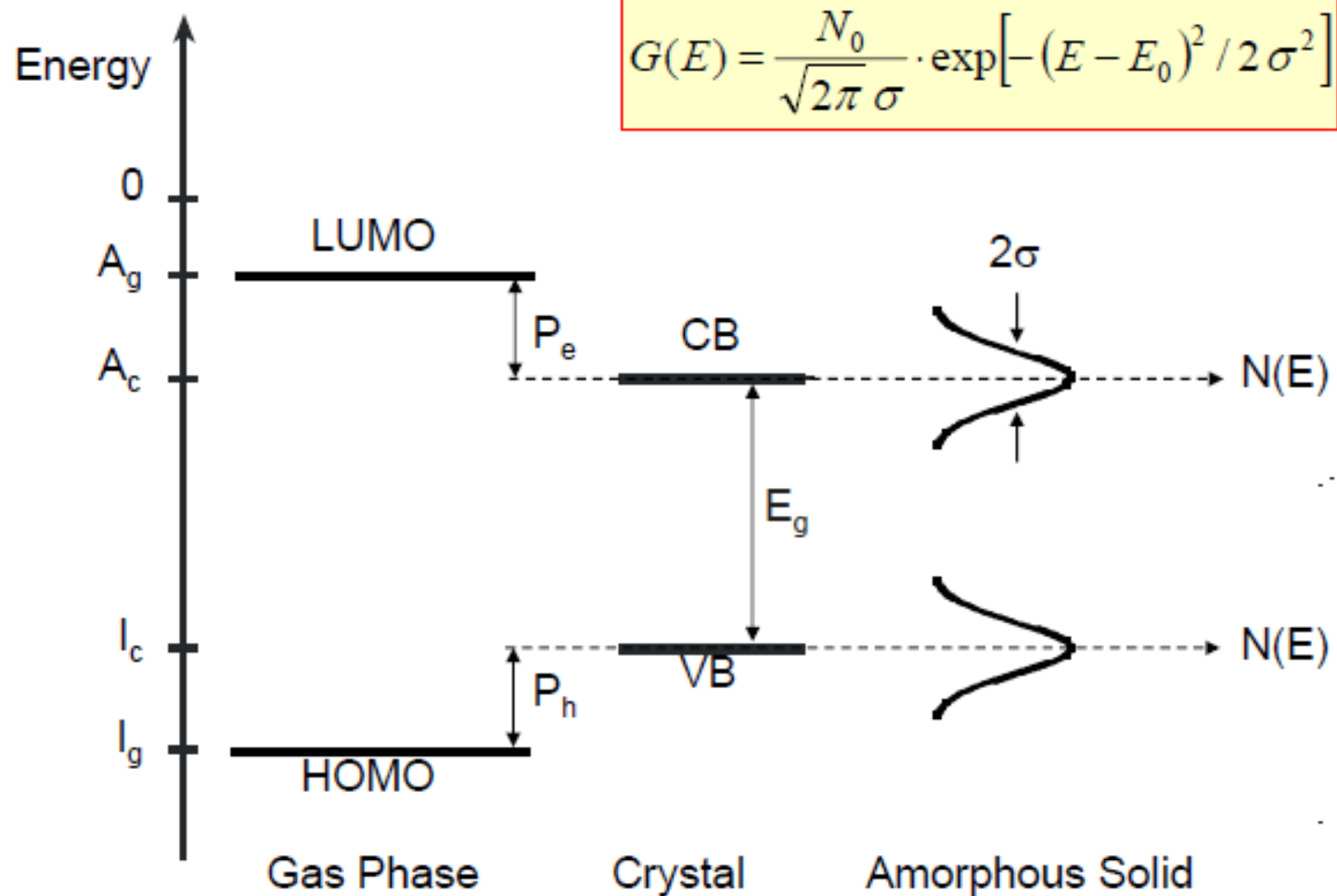
Polarisation  
in the vicinity  
of a charged  
molecule



- Charge transport via transfer of pos./neg. charge from one molecule to the next.

## Disorder Effects

- Variations of polarisation energies due to locally different molecular environment → distribution of energy levels (Gaussian type with typical width  $\sigma \approx 100$  meV)



## Polarisation Energy

- Origin of the polarisation energy: induced dipole moments on neighbouring molecules (plus higher order contributions)
- Simple Model:

$$P = P' + \Delta P$$

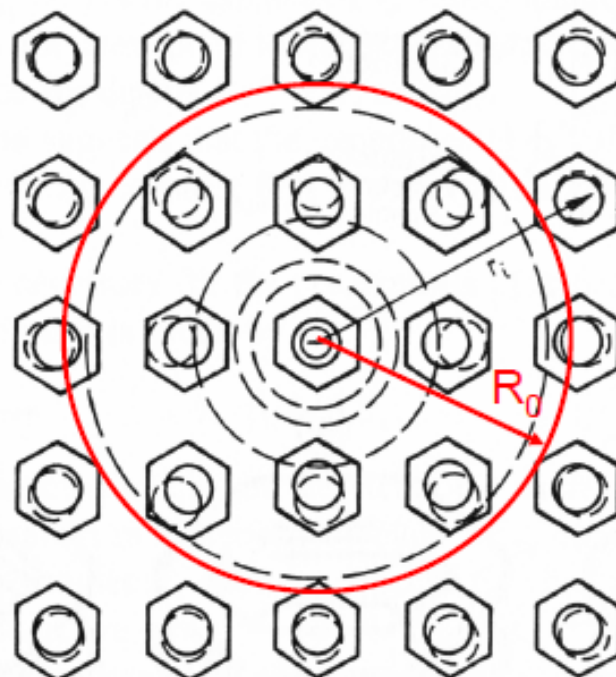
$r < R_0$  : discrete sum

$$P' = \frac{1}{2} \sum_i \vec{F}_i(\vec{r}_i) \cdot \vec{\mu}_i$$

$r > R_0$  : continuum

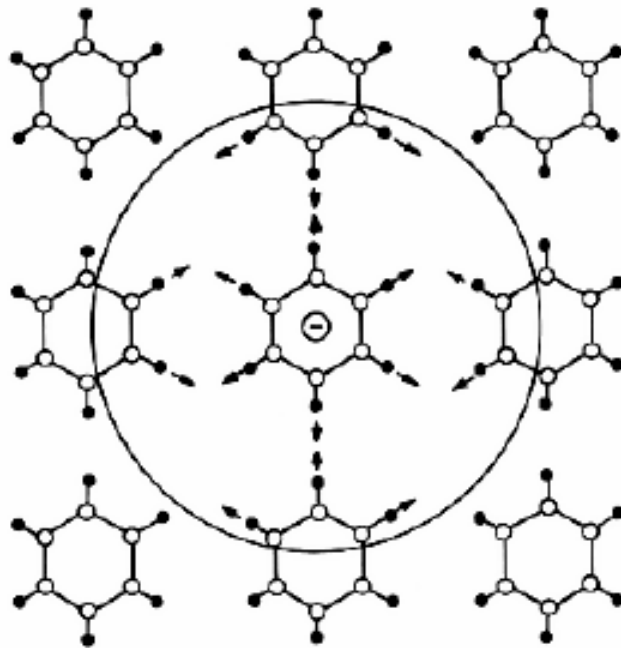
$$\Delta P = \frac{1}{4\pi\epsilon_0} \cdot \frac{e^2}{2R_0} \cdot \left(1 - \frac{1}{\epsilon}\right)$$

- Anthracene:  
 $P' = -1.26$  eV,  $\Delta P = -0.27$  eV
- In general:  $P_e \sim P_h = P = 1 \dots 2$  eV  
→ polarisation effect not negligible !



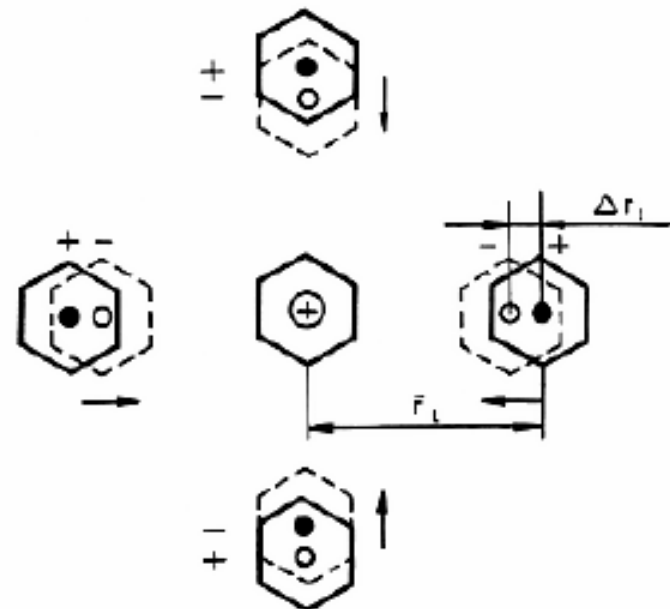
# Polarons in Molecular Crystals

- Molecular polaron: interaction with **intramolecular** vibrations



- Typical energies  $\sim 0.15\text{eV}$

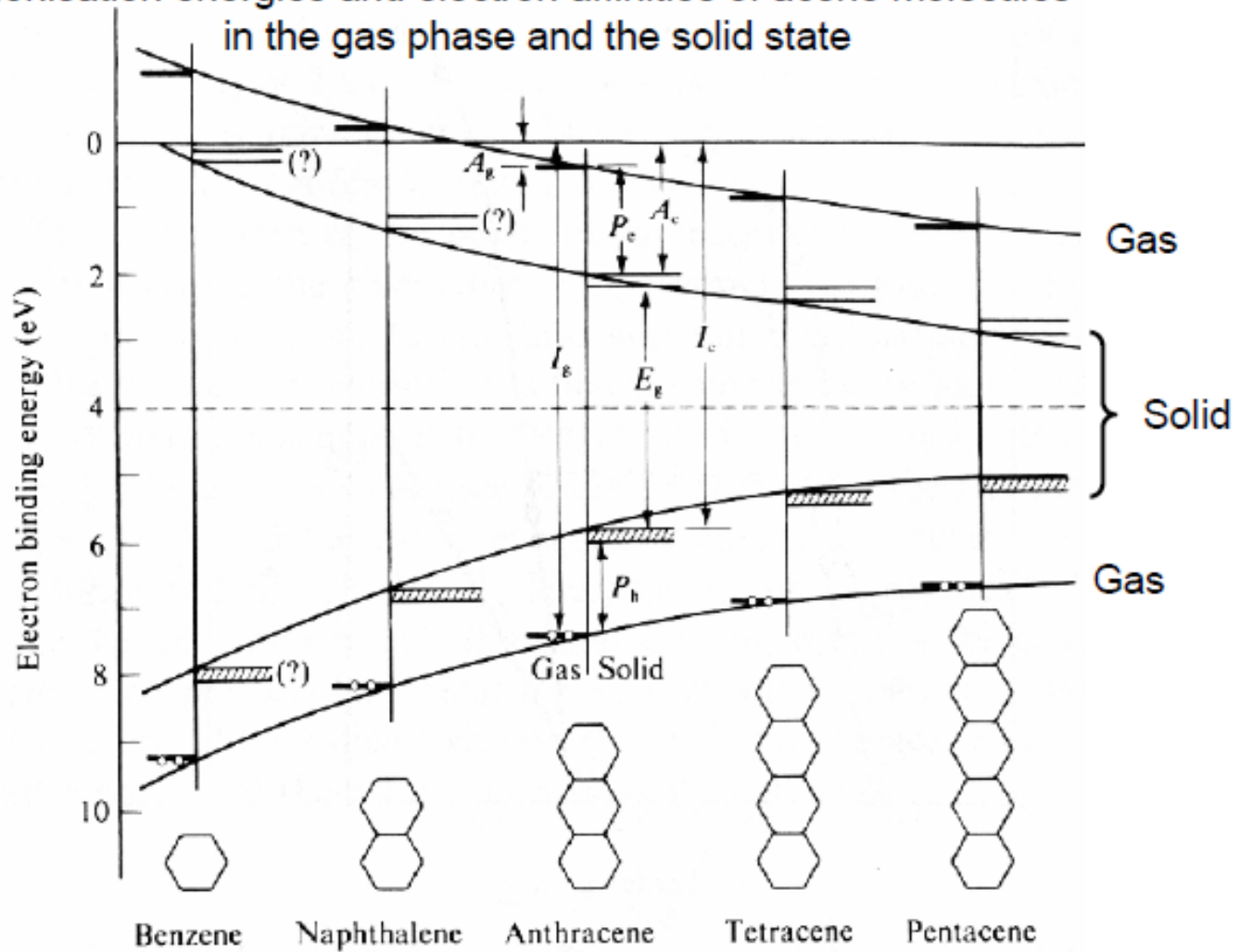
- Lattice polaron: interaction with **intermolecular** vibrations



- Typical energies  $< 0.1\text{eV}$

# Polarisation Energies in Acenes

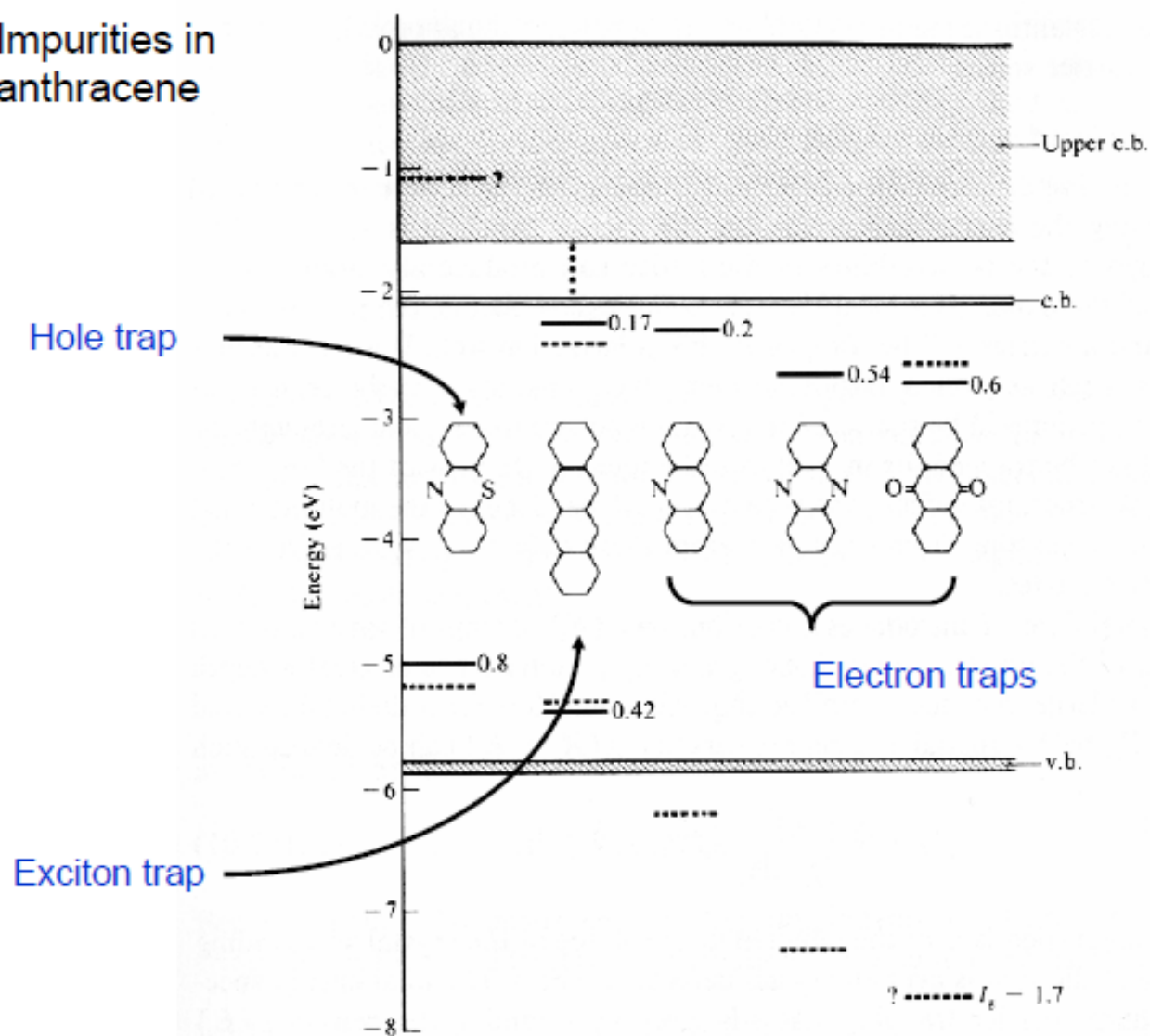
Ionisation energies and electron affinities of acene molecules  
in the gas phase and the solid state





# Charge Carrier Traps in Molecular Crystals

Impurities in anthracene



Additionally:

- structural defects
  - disorder
- trap distribution (exponential / Gaussian)

# Charge Carrier Mobility

---

**Band Transport:** energy bands

$$\mu \propto T^{-n} \quad (n \approx 1.5)$$

T = 300 K:  $\mu \sim 1000 \text{ cm}^2/\text{Vs}$  (crystalline Si)

→ **Phonon scattering**

(at low T: impurity scattering  $\mu \propto T^{+n}$ )

**Hopping-Transport:** localised carriers

T = 300 K: →  $\mu \ll 1 \text{ cm}^2/\text{Vs}$  (amorphous materials)

$$\mu \propto \exp(-E_A/k_B T)$$

→ **Phonon-assisted hopping**

## Charge Carrier Density

---

Intrinsic carrier density:  $n_i = N_0 \cdot \exp(-E_g / 2k_B T)$

T = 300 K

• Si:  $E_g = 1.12 \text{ eV}$ ;  $N_0 = 10^{19} \text{ cm}^{-3}$   $\rightarrow n_i = 10^{10} \text{ cm}^{-3}$   
(Doping:  $n = 10^{15} \dots 10^{19} \text{ cm}^{-3}$ )

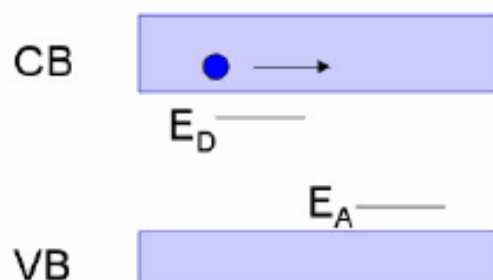
• Organic semiconductor  
with  $E_g = 2.5 \text{ eV}$ ;  $N_0 = 10^{21} \text{ cm}^{-3}$   $\rightarrow n_i = 1 \text{ cm}^{-3}$

→ Peculiarities of organic semiconductors:

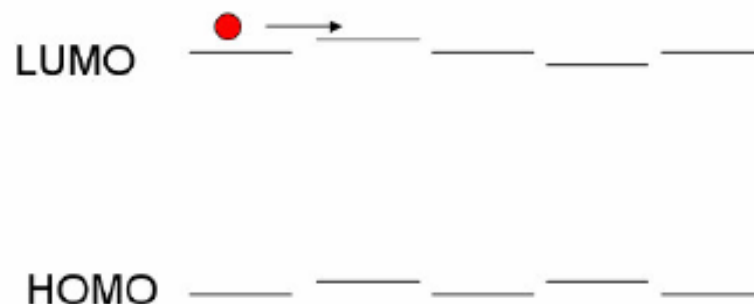
- Current carried by excess carriers  $\rightarrow$  non-linear I-V characteristics
- Strongly field and temperature dependent mobility
- Anisotropic conductivity and mobility  $\rightarrow$  tensors

## Inorganic vs. Organic semiconductors

- broad bands
- small correlation energies ( $e-h \approx 4\text{meV}$ )
- hydrogen model works

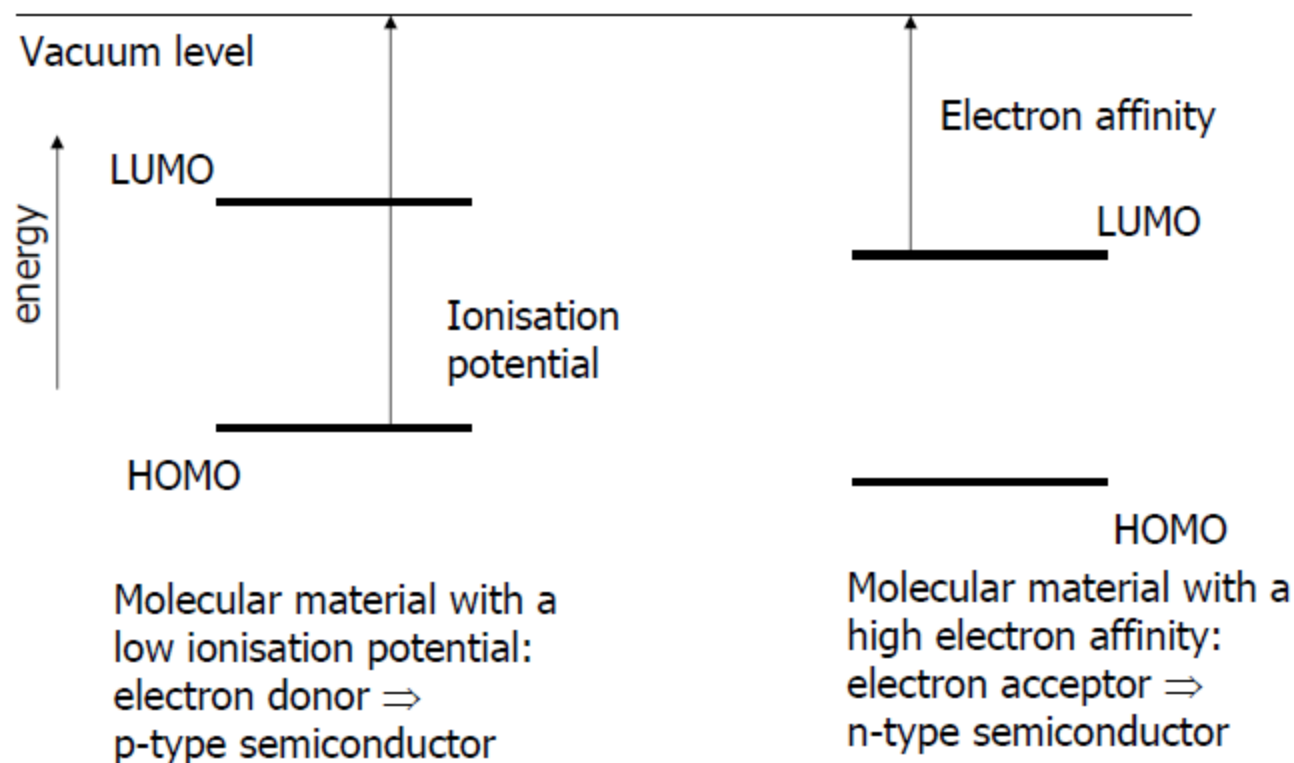


- hopping transport
- large correlation ( $e-h \approx 1\text{ eV}$ )
- polaron effects important



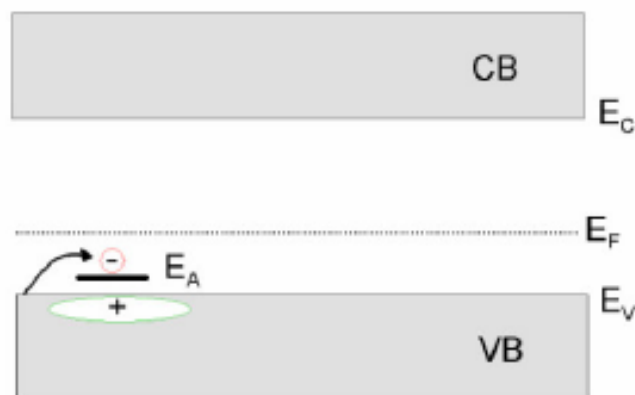
Inorganic Semiconductor		Organic Semiconductor
Band	Transport Mechanism	Hopping
$10^2\text{-}10^3$	RT mobility ( $\text{cm}^2/\text{Vs}$ )	$10^{-6}\text{-}1$ (typ. $10^{-3}$ )
$10^{15}\text{-}10^{18}$ (doping controlled)	Charge carrier concentration	$10^{10}\text{-}10^{16}$ (injection controlled)
$<<10^{15}$	Electr. Active impurit. ( $\text{cm}^{-3}$ )	$\approx 10^{17}$

## n- or p-type molecular semiconductors

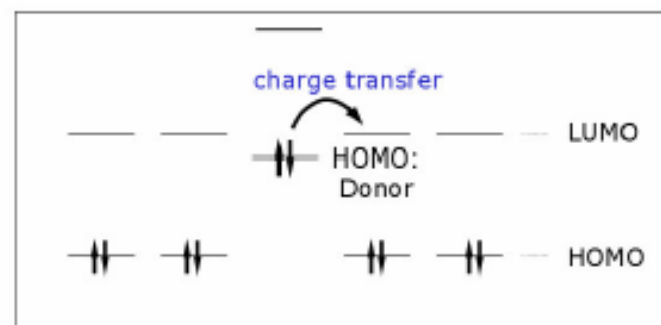


# Doping organic semiconductors

## Inorganic



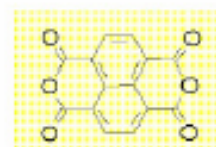
## Organic



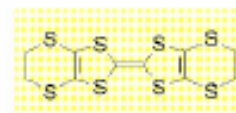
Organic dopants: electron transfer from high lying HOMO to matrix LUMO

Both n and p doping are possible in organic semiconductors by introducing charge donor or acceptor dopants into the molecular host matrix.

Charge donation of metal atoms (for example, alkaline elements) is very effective in modifying drastically (several orders of magnitude) the conductivity of organic materials.



Matrix:  
Naphthalin-tetracarboxylic-  
dianhydride  
(NTCDA)



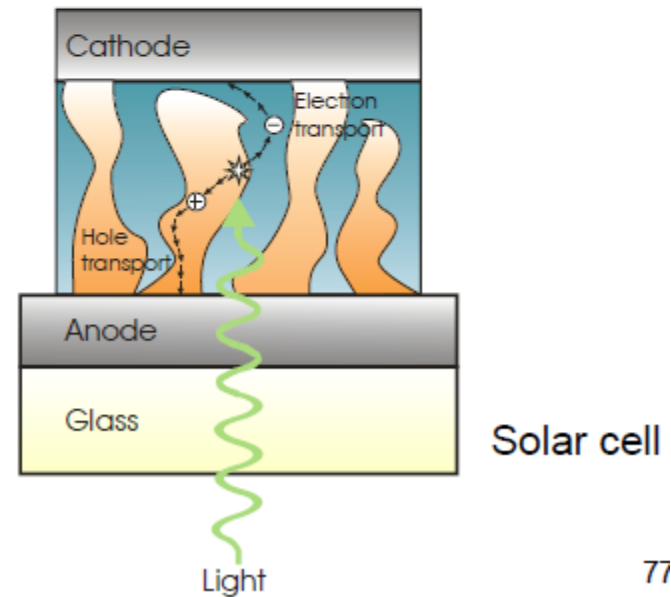
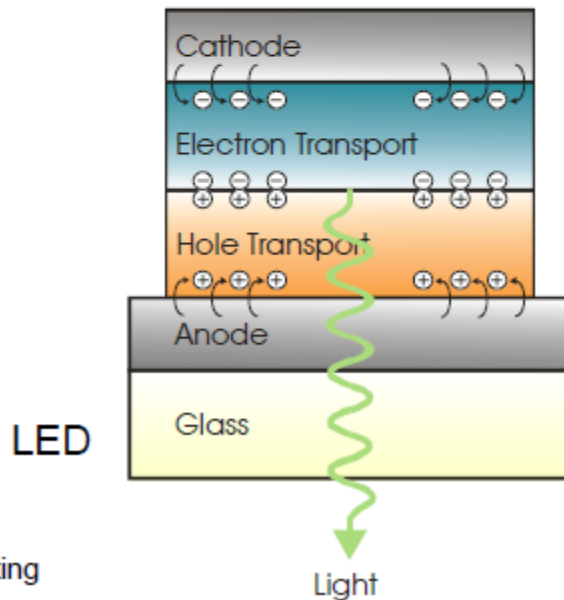
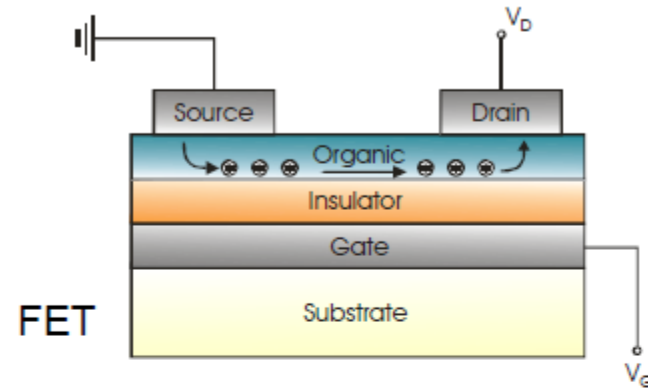
Donor:  
Bisethylenedithio-  
tetrathiafulvalene  
(BEDT-TTF)

# Charge (photo)generation

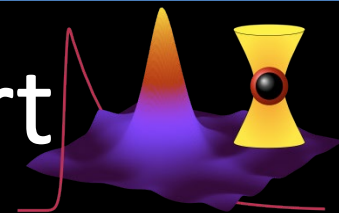
November 12, 2025

# Charge Carrier Generation

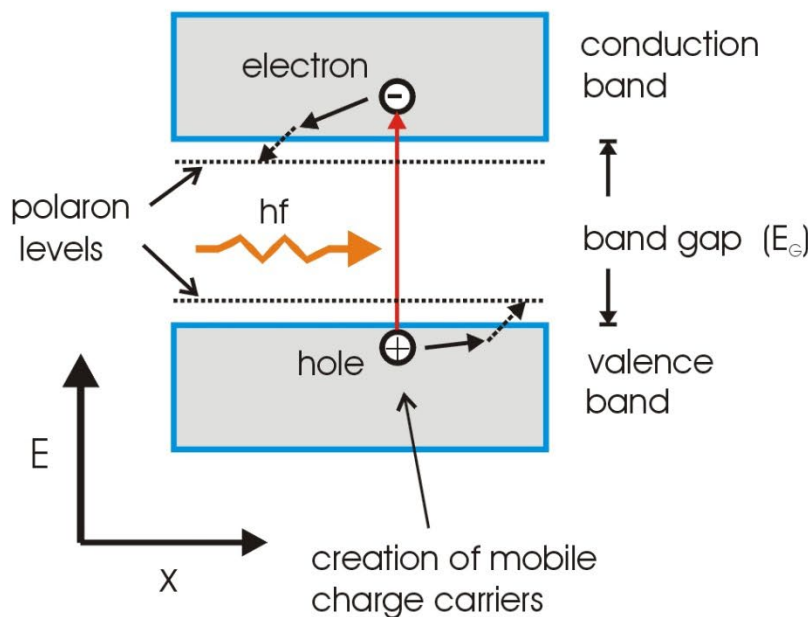
- Injection from contacts:  
→ light-emitting diodes
- Field-effect doping:  
→ field-effect transistors
- Photogeneration:  
→ photodiodes and solar cells
- (Electro-)chemical doping



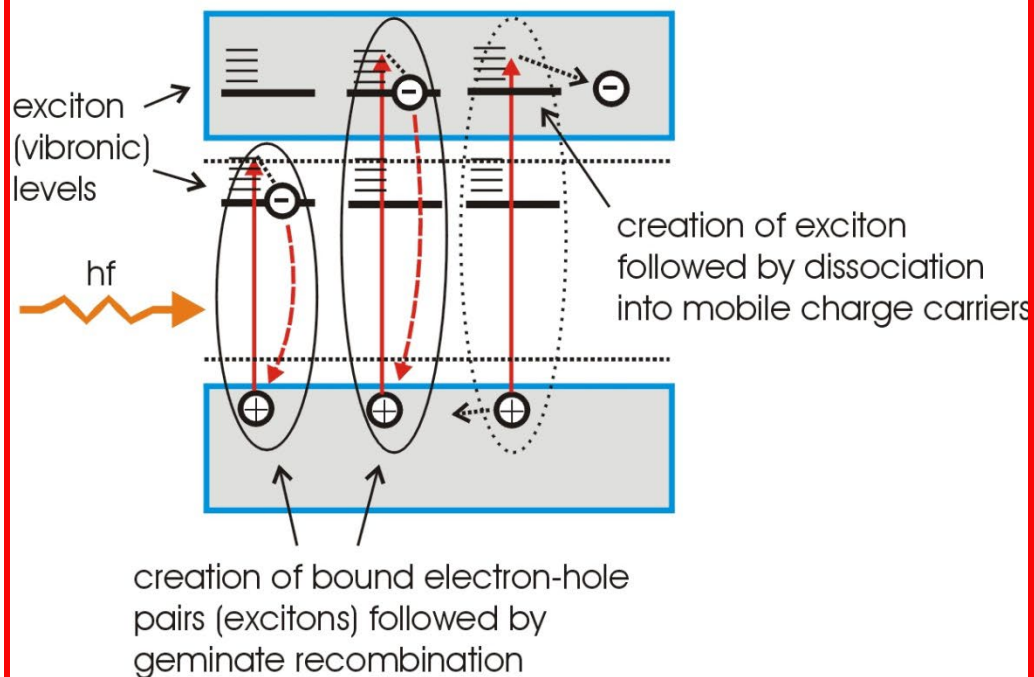




## SEMICONDUCTOR BAND MODEL



## MOLECULAR EXCITON MODEL



- Exciton binding energies  $\sim 0.1-0.5$  eV: difficult to dissociate
- Use D-A systems to improve charge photogeneration

## Wannier-Mott Excitons

Coulombic interaction between the hole and the electron is given by

$$E_{EX} = -e^2/\epsilon r$$

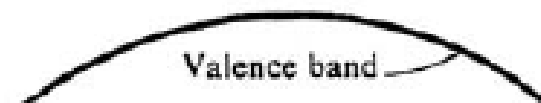
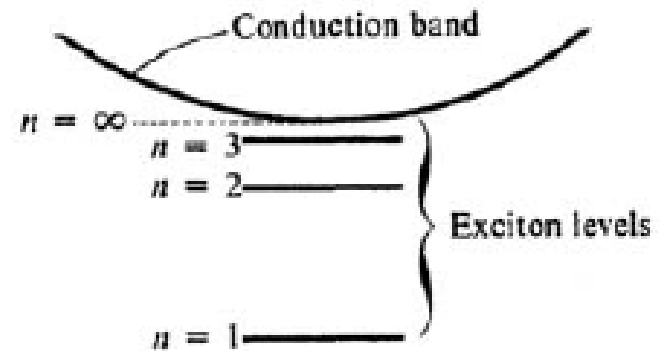
The exciton energy is then

$$E = E_{ION} - E_{EX}/n^2, \quad n = 1, 2, \dots$$

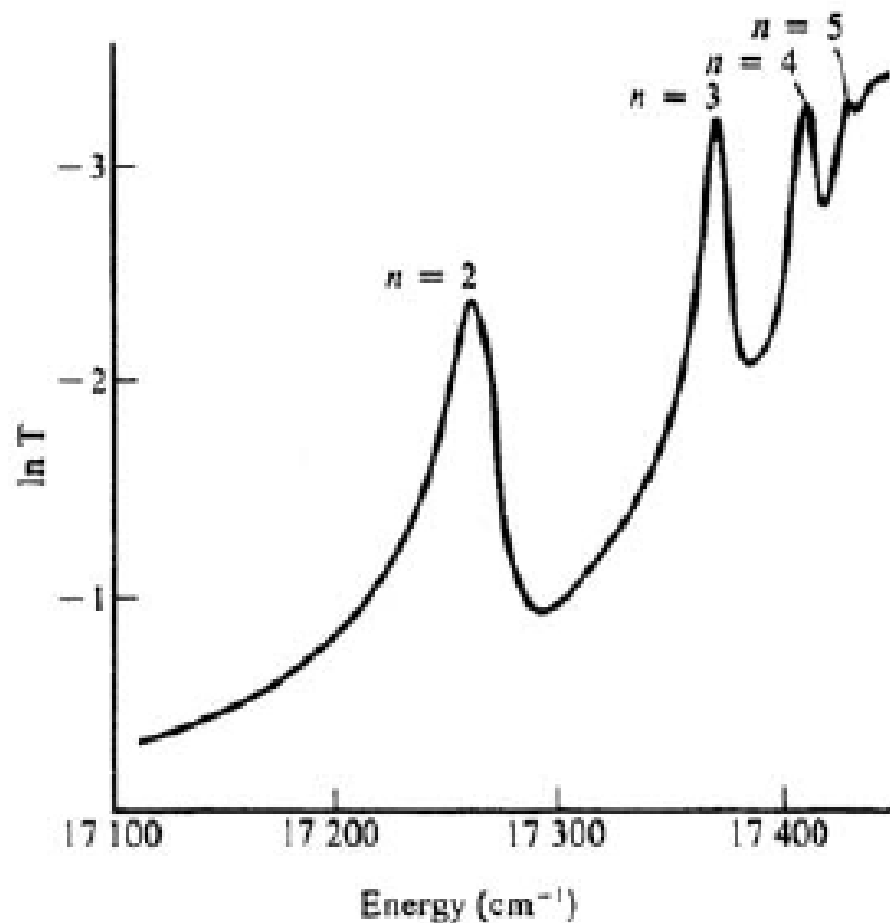
$E_{ION}$  - energy required to ionize the molecule  
 $n$  - exciton energy level

$$E_{EX} = 13.6 \text{ eV } \mu / m_e$$

$$\mu\text{- reduced mass} = m_e m_h / (m_e + m_h)$$



## An Example of Wannier-Mott Excitons



exciton progression  
fits the expression

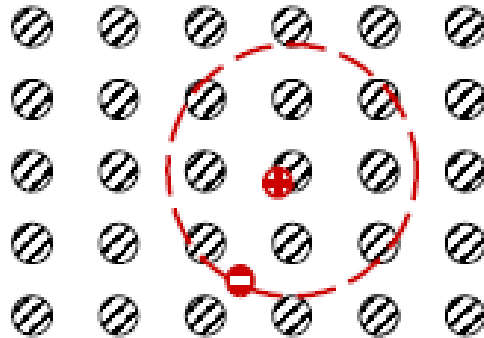
$$\nu[\text{cm}^{-1}] = 17,508 - 800/n^2$$

corresponding to  
 $\mu = 0.7$  and  $\epsilon = 10$

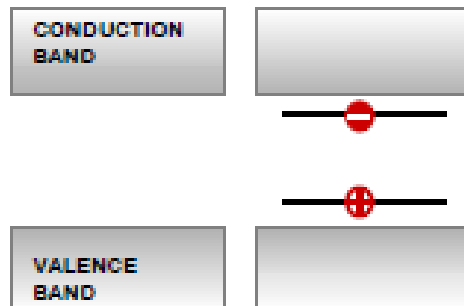
**Fig. I.D.28.** The absorption spectrum of  $\text{Cu}_2\text{O}$  at 77 K, showing the exciton lines corresponding to several values of the quantum number  $n$ . (From Baumeister 1961)

### Wannier exciton

(typical of inorganic semiconductors)



#### SEMICONDUCTOR PICTURE



GROUND STATE      WANNIER EXCITON

binding energy  $\sim 10\text{meV}$   
radius  $\sim 100\text{\AA}$

## Excitons (bound electron-hole pairs)

treat excitons  
as **chargeless**  
**particles**  
capable of  
diffusion,

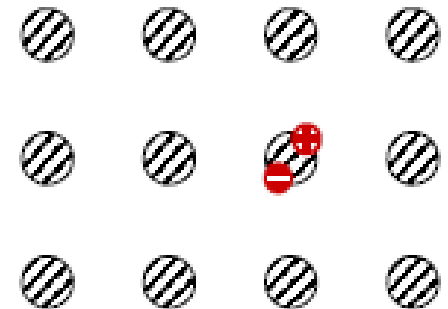
also view  
them as  
excited states  
of the  
molecule

### Charge Transfer (CT)

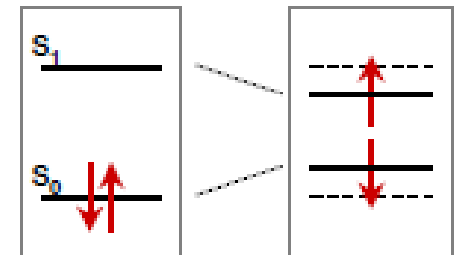
Exciton  
(typical of organic  
materials)

### Frenkel exciton

(typical of organic  
materials)



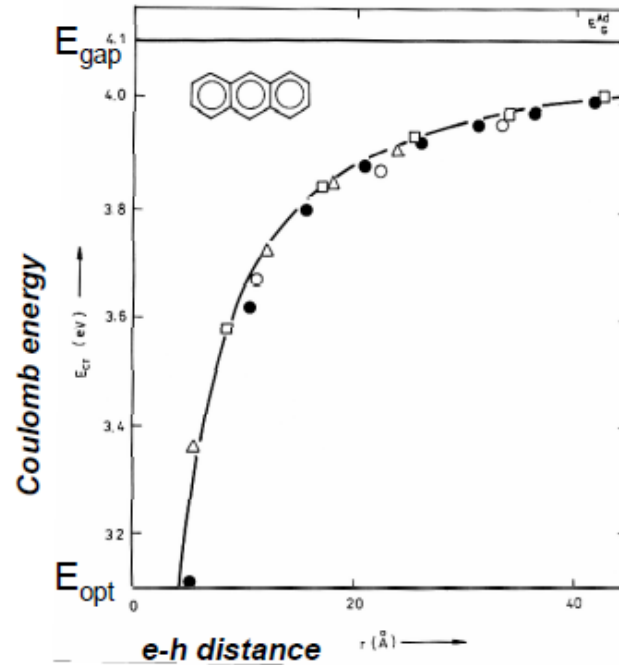
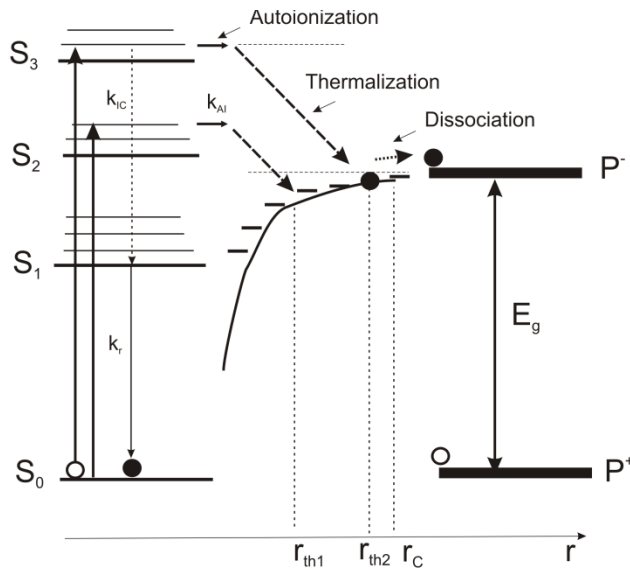
#### MOLECULAR PICTURE



GROUND STATE      FRENKEL EXCITON

binding energy  $\sim 1\text{eV}$   
radius  $\sim 10\text{\AA}$

# Charge Separation and Recombination



Polaron gap ~4.1eV

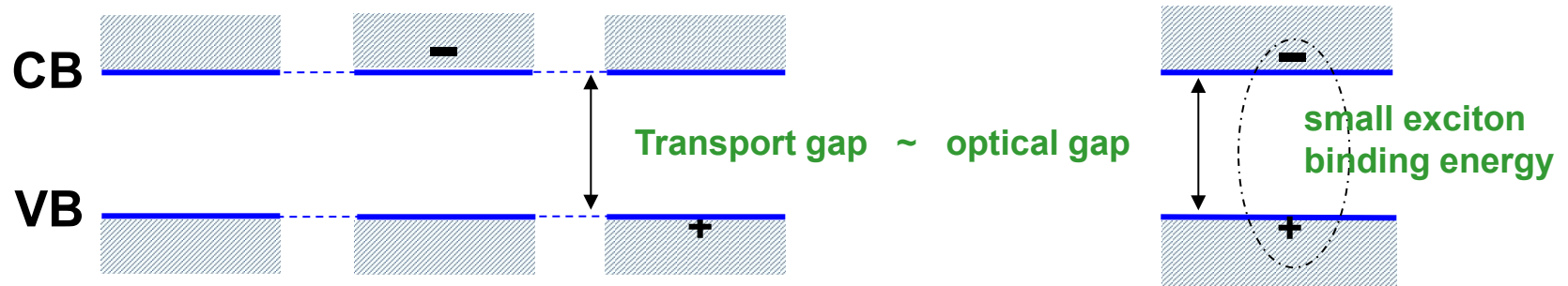
Optical gap ~3.1eV

**Critical distance:**  
**→ Coulomb radius**

$$\frac{1}{4\pi\epsilon\epsilon_0} \cdot \frac{e^2}{r_C} = kT \quad (\epsilon = 3, T = 300K)$$

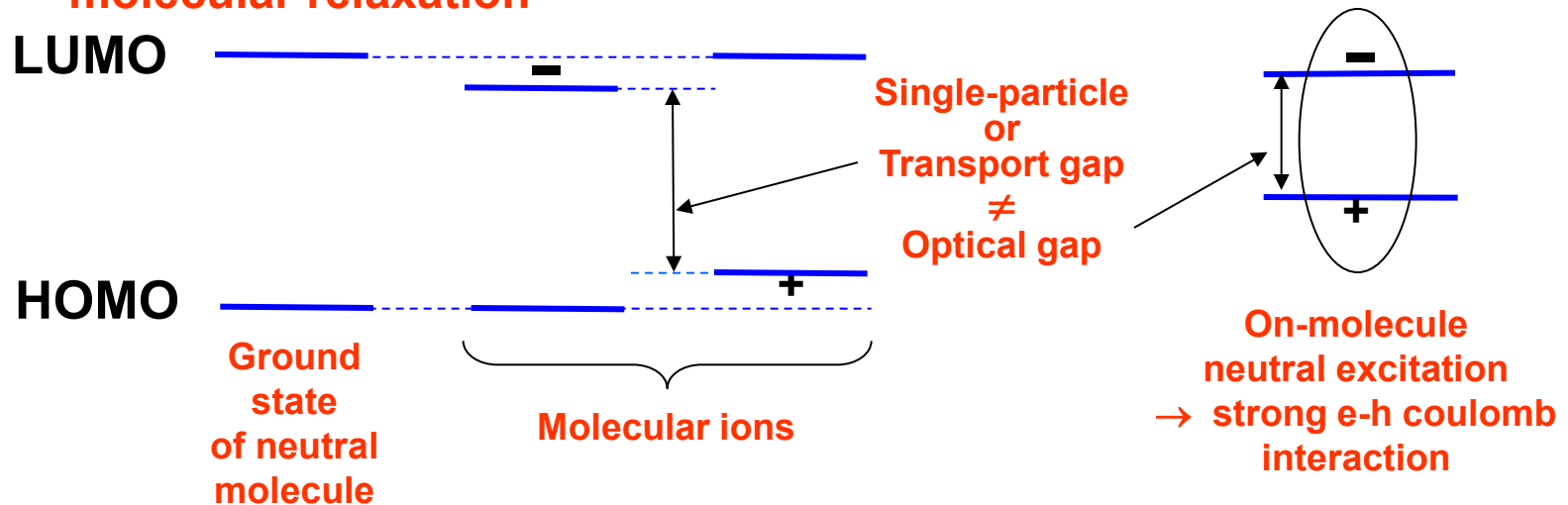
$$\Rightarrow r_C \approx 190 \text{ Å} \quad \text{"Coulomb radius"}$$

- Inorganic semiconductor:** wide bands and delocalized states



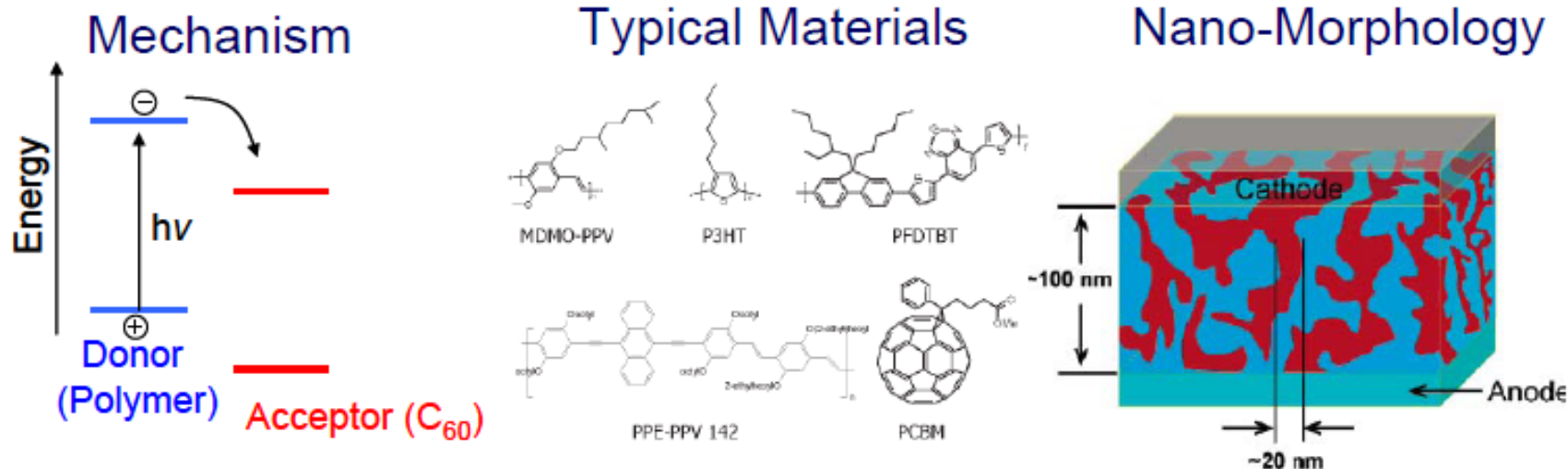
Energy levels and transport: Bloch states and single-electron approximation

- Organic molecular solid:** small transfer integral between molecules; charge carrier = molecular ions; electronic polarization + molecular relaxation

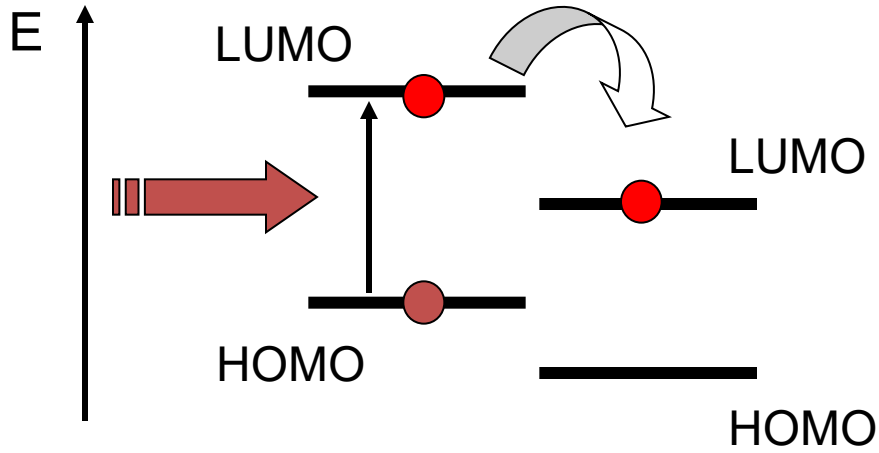


Transport gap – optical gap = exciton binding energy

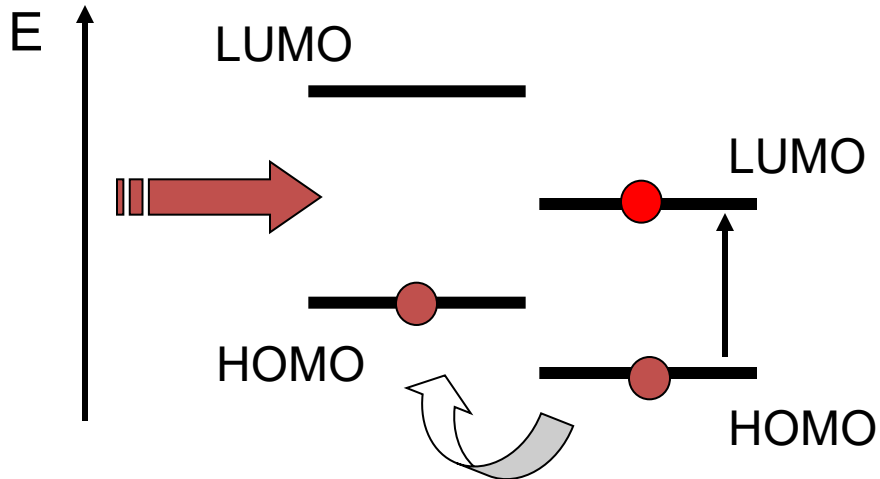
# Solar Cell Characteristics



# Photoinduced Charge Transfer



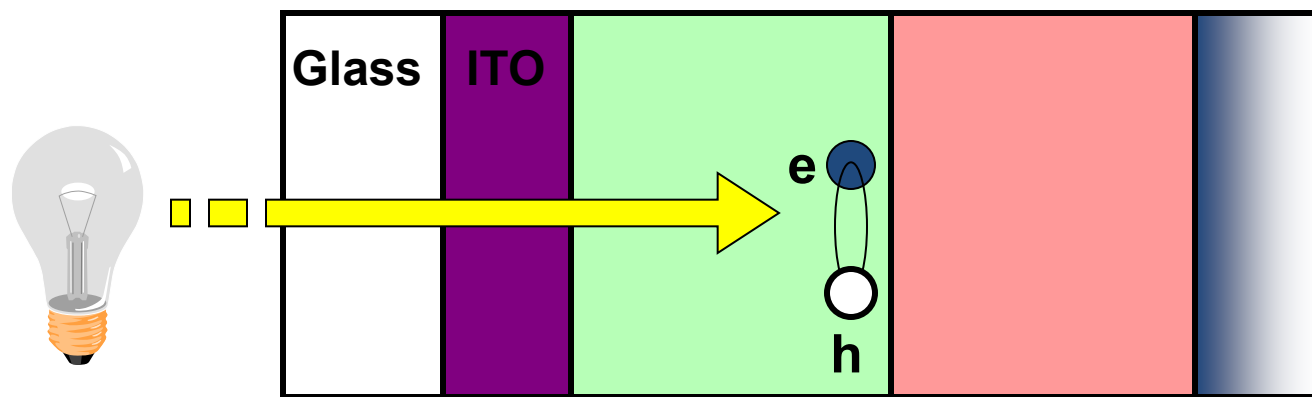
Photoinduced ELECTRON  
transfer



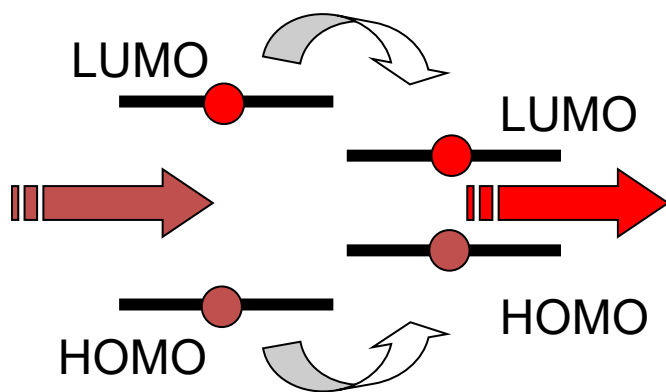
Photoinduced HOLE  
transfer



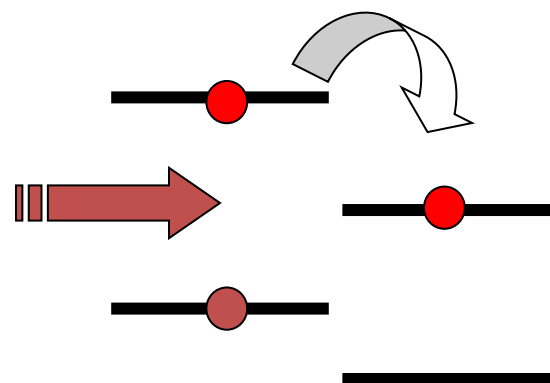
# Charge and energy transfer in conjugated polymers



**Organic Solar Cells**



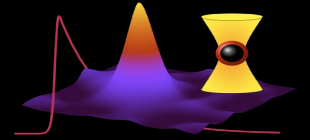
**Energy transfer**



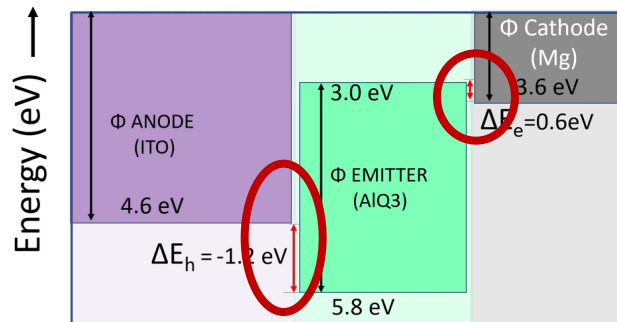
**Charge transfer**

**TABLE 1.1.** *Comparison of Characteristic Features of Organic Molecular and Covalent (Atomic) Crystals*

Molecular crystal (e.g., anthracene-type crystal)	Covalent (atomic) crystal (e.g., silicon-type crystal)
Weak Van der Waals type of interaction (characteristic interaction energies $E_{\text{vdW}} = 10^{-3} - 10^{-2}$ eV)	Strong covalent-type interaction (characteristic interaction energies $E_{\text{cov}} = 2 - 4$ eV)
Marked tendency of charge carrier and exciton localization	Pronounced charge carrier delocalization
Self-energy of charge carriers and excitons determined by many-electron interaction (polarization) effects	Single-electron approximation valid
Charge carriers and excitons as polaron-type quasiparticles	Charge carriers as free holes and electrons
Low charge carrier mobilities ( $\mu \approx 1$ cm <sup>2</sup> /Vs) and small mean free path ( $l \approx a_0 =$ lattice constant) at room temperatures	High charge carrier mobilities and large mean free path [ $l = (100 - 1000) a_0$ ]
Large effective mass of charge carriers $m_{\text{eff}} = (10^2 - 10^3) m_e$	Small effective mass of charge carriers $m_{\text{eff}} \leq m_e$
Hopping-type charge carrier transport dominant	Band-type charge carrier transport dominant
Excitons as molecular Frenkel-type quasiparticles	Excitons as Wannier-type quasiparticles
Low melting and sublimation temperatures, low mechanic strength, high compressibility	High melting and sublimation temperatures, high mechanical strength, low compressibility

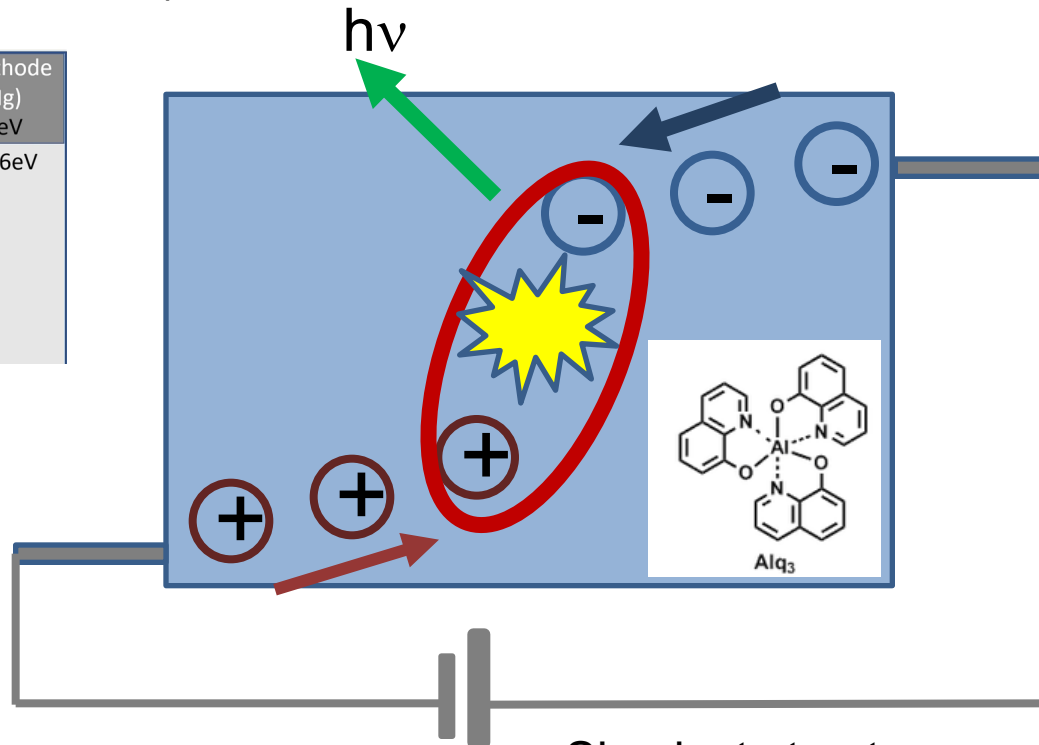


Step 1: inject holes and electrons  
(choose appropriate electrodes)



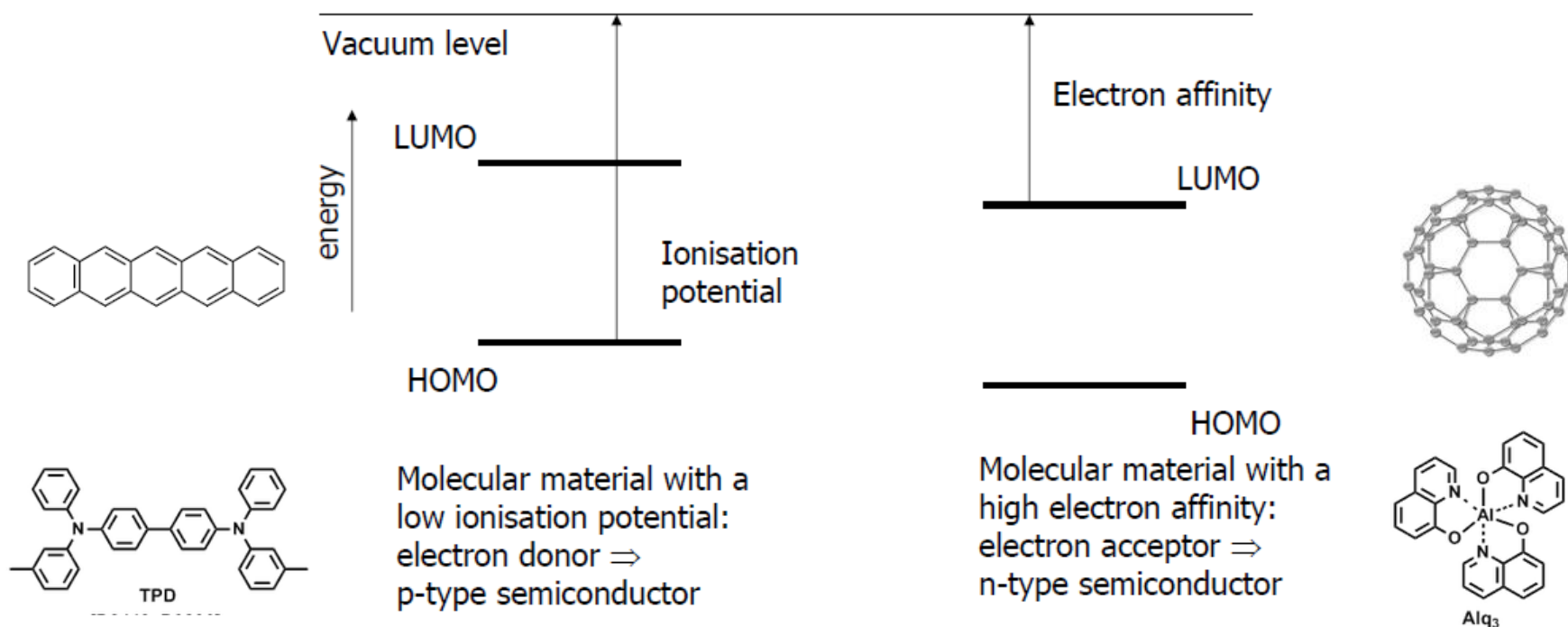
- Work function: how much energy is needed to release the charge
- Minimize energy barriers
- One of the electrodes (ITO) transparent

E.g. use FTO as a transparent Hole injector and Ytterbium ( $\Phi = 2.6$  eV) as an electron injector

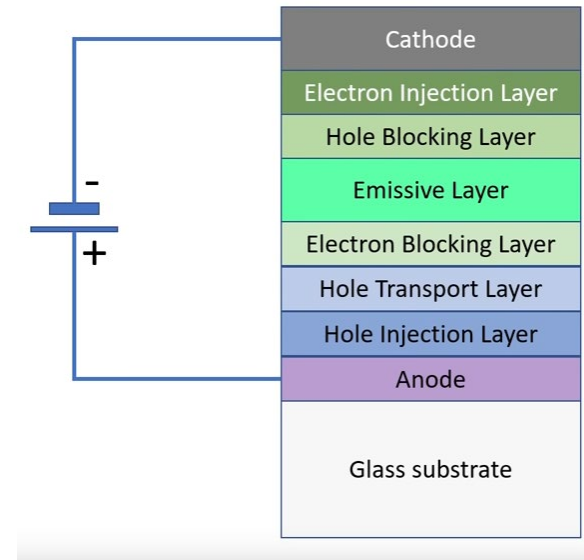
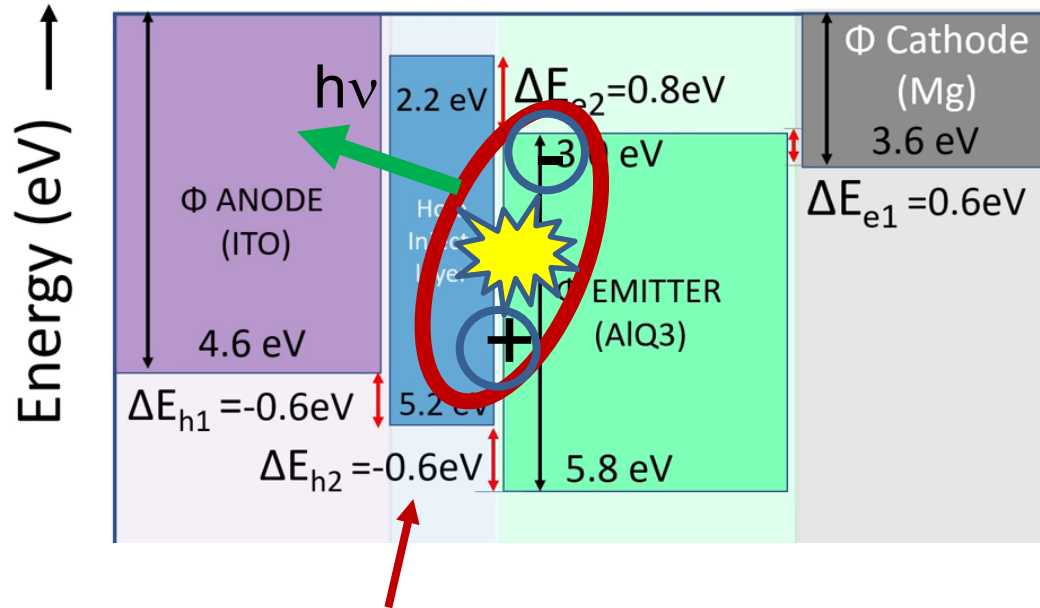
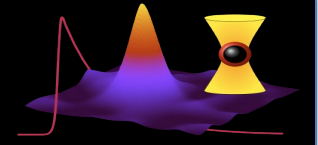


Simplest structure:  
electrode/Alq3/electrode  
Problem: Alq3 is better for electron than hole transport, and the hole injection barrier too high

# n- or p-type molecular semiconductors



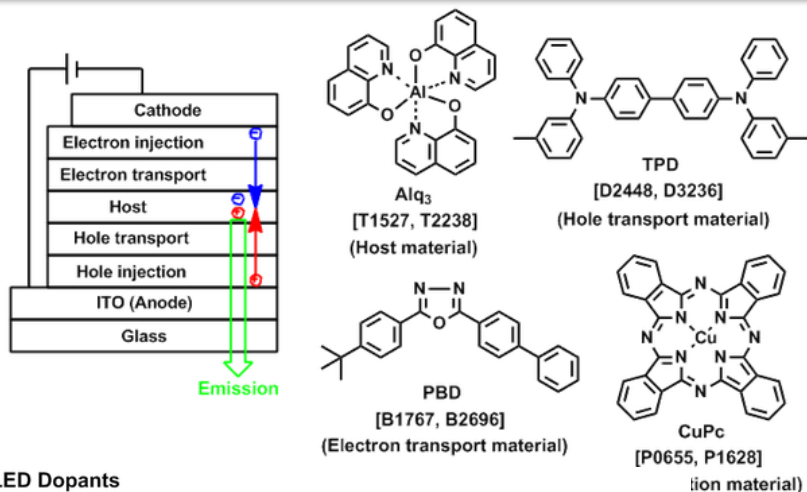
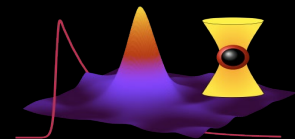
Depending on the molecular energy levels, some molecules transport holes and other transport electrons more efficiently



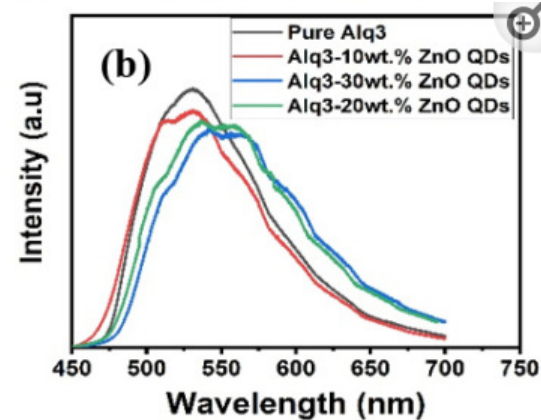
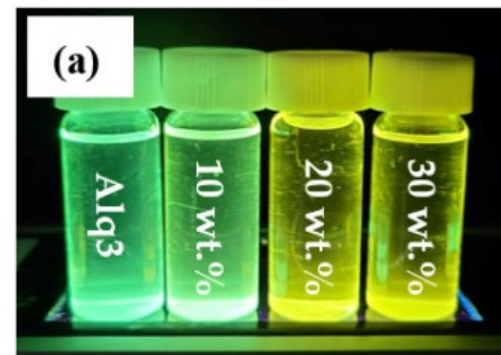
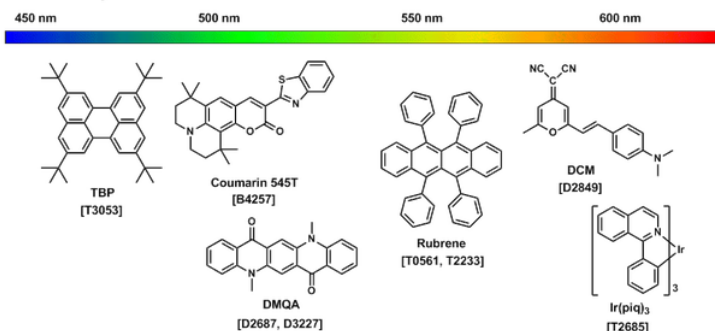
Insert a hole injection layer to help inject and transport holes (e.g. TPD)

More complicated structures

Example structure: glass/FTO/TPD (hole injection/transport)/Alq3 (electron injection/emitter)/Ytterbium



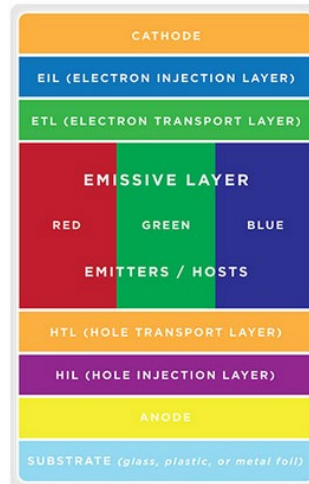
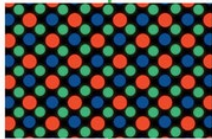
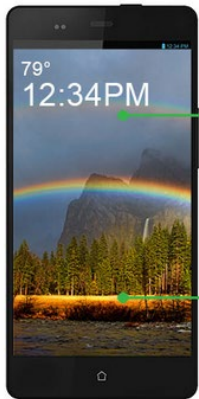
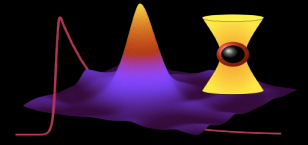
OLED Dopants



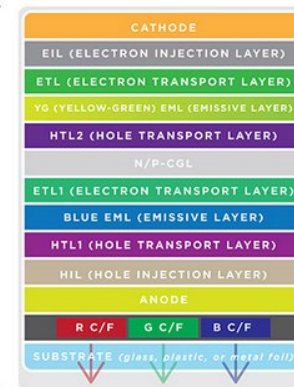
<https://www.tcichemicals.com/US/en/c/127>

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# Multicolor OLEDs



Incorporate multiple emitters



<https://oled.com/oleds/>