

QM

Harmonic oscillator; measurements, uncertainty relationsExamples of measurements

Ex. 1 A quantum H.O. is in a state $|\psi\rangle = \sqrt{\frac{2}{3}}|0\rangle + \frac{1}{\sqrt{3}}|3\rangle$ $\leftarrow \hat{H}|n\rangle = E_n|n\rangle$

If energy is measured, what are outcomes / probabilities?

$$|\psi\rangle = \sum_n c_n |n\rangle \Rightarrow P(E_n) = |c_n|^2 \Rightarrow$$

outcomes: $E_0 = \frac{\hbar\omega}{2} \Rightarrow P_0 = \frac{2}{3}$

$$E_3 = \hbar\omega \cdot \underbrace{\left(3 + \frac{1}{2}\right)}_{7/2} \Rightarrow P_3 = \frac{1}{3}$$

Ex. 2 A quantum H.O. is in a state described by a wave function $\psi(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-\frac{m\omega x^2}{2\hbar}}$

$$\left(\frac{1}{\sqrt{2}} + \frac{2x}{\sqrt{2}} \sqrt{\frac{m\omega}{2\hbar}} \right)$$

If energy is measured, what are the outcomes & probabilities?

Need to present

$$\psi(x) = \sum_n \phi_n(x) = \frac{1}{\sqrt{2}} \phi_0(x) + \frac{1}{\sqrt{2}} \phi_1(x) \Rightarrow$$

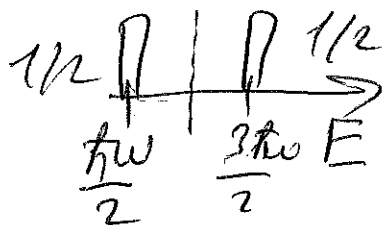
$$\hat{H}\phi_n(x) = E_n \phi_n(x)$$

$$E_0 = \frac{\hbar\omega}{2} \Rightarrow P_0 = 1/2$$

$$E_1 = 3/2 \hbar\omega \Rightarrow P_1 = 1/2$$

What is the expectation value of energy? (2)

$$\langle E \rangle = \underbrace{P_0}_{1/2} \cdot \frac{\hbar\omega}{2} + \underbrace{P_1}_{1/2} \cdot \frac{3\hbar\omega}{2} = \frac{\hbar\omega}{4} + \frac{3\hbar\omega}{4} = \hbar\omega$$



sanity check!
 $\frac{3}{2}\hbar\omega$ with equal probability
 $\hbar\omega$

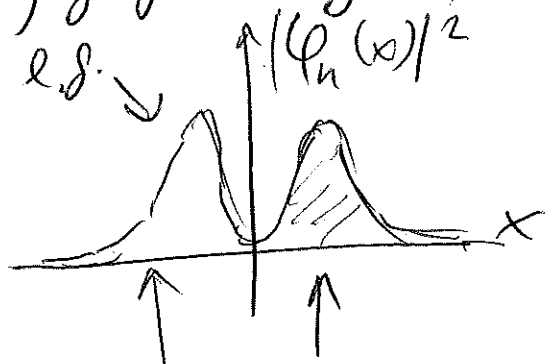
average in the middle!

Ex. 3 Find expectation

value of $\hat{X}$ if a quantum H.O. is in some state $|n\rangle$
 $\uparrow$
 position operator

$|n\rangle$
 (eigenstate of $\hat{H}$)
 so $\hat{H}|n\rangle = E_n|n\rangle$

1) graphically $\Rightarrow \psi_n(x)$ is either even or odd $\Rightarrow |\psi_n|^2$ is always even



$\uparrow$
 since $V(x) = V(-x)$
 same probability for H.O.

to find the particle at x & $-x$! $\Rightarrow \langle X \rangle = 0$

$$\begin{aligned} 2) \text{ integral } \Rightarrow \langle X \rangle &= \int_{-\infty}^{+\infty} \psi_n^*(x) X \psi_n(x) dx = \\ &= \int_{-\infty}^{+\infty} \underbrace{|\psi_n(x)|^2}_{\text{even}} \underbrace{x}_{\text{odd}} dx = 0 \end{aligned}$$

$\uparrow$ odd = even; even $|_{-\infty}^{+\infty} = 0$

3) ket (abstract) (the best!) (3)

recall: $a = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{X} + i \frac{\hat{P}}{m\omega} \right)$

$\Downarrow$ $a^\dagger = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{X} - i \frac{\hat{P}}{m\omega} \right)$

solve for $\hat{X}$ & $\hat{P}$ (in terms of $a, a^\dagger$)

$$\boxed{\hat{X} = \sqrt{\frac{\hbar}{2m\omega}} (a^\dagger + a)} \quad \boxed{\hat{P} = i \sqrt{\frac{\hbar m\omega}{2}} (a^\dagger - a)}$$

$\Downarrow$
 $\langle n | \hat{X} | n \rangle = \sqrt{\frac{\hbar}{2m\omega}} \langle n | a^\dagger + a | n \rangle = \sqrt{\frac{\hbar}{2m\omega}} \cdot$

$\cdot \left(\underbrace{\sqrt{n+1}}_0 \langle n | a^\dagger | n \rangle + \underbrace{\sqrt{n}}_0 \langle n | a | n \rangle \right) = 0$ $a | n \rangle = \sqrt{n} | n-1 \rangle$
 $a^\dagger | n \rangle = \sqrt{n+1} | n+1 \rangle$

(Note: similarly $\langle P \rangle \sim \langle n | a^\dagger - a | n \rangle = 0$)

What about $\langle X^2 \rangle$? $\Rightarrow$ use abstract! $\Rightarrow$

$$\begin{aligned} \langle n | \hat{X}^2 | n \rangle &= \frac{\hbar}{2m\omega} \langle n | (a^\dagger + a)^2 | n \rangle = \\ &= \frac{\hbar}{2m\omega} \langle n | \underbrace{a a^\dagger}_{n+1} + \underbrace{a^\dagger a}_n | n \rangle = \frac{\hbar}{2m\omega} (2n+1) \end{aligned}$$

Uncertainty relations

(4)

Recall: $\Delta X = \sqrt{\langle X^2 \rangle - \langle X \rangle^2}$
 $\Delta P = \sqrt{\langle P^2 \rangle - \langle P \rangle^2} \Leftrightarrow$ standard deviation

In general, $\Delta A \cdot \Delta B \geq \frac{1}{2} |\langle [\hat{A}, \hat{B}] \rangle|$

if $[\hat{A}, \hat{B}] \neq 0 \Rightarrow \hat{A}$ & $\hat{B}$ are incompatible observable

if $\Delta A \rightarrow 0$

$\Delta B \rightarrow \infty$

(and the other way around)

if $[\hat{A}, \hat{B}] \neq 0$

can't be

$\Leftarrow$ simultaneously measured with high degree of accuracy

can achieve $\downarrow \Delta A \rightarrow 0$ for $[\hat{A}, \hat{B}] = 0$
 $\Delta B \rightarrow 0$ but not for $[\hat{A}, \hat{B}] \neq 0$

If $\hat{A} = \hat{X}$, $\hat{B} = \hat{P} \Rightarrow \Delta X \Delta P \geq \frac{\hbar}{2}$

$[\hat{X}, \hat{P}] = i\hbar$

So how does this look like for quantum H.O.?

$\Delta X \Rightarrow \langle X \rangle = 0, \langle X^2 \rangle = \frac{\hbar}{2m\omega} (2n+1)$

$$\text{So } \Delta x = \sqrt{\frac{\hbar}{2m\omega}} \sqrt{2n+1}$$

(5)

$$\Delta p \Rightarrow \langle p \rangle = 0, \quad \langle p^2 \rangle = \frac{m\hbar\omega}{2} (2n+1)$$

↑
HW # 2

$$\text{So } \Delta p = \sqrt{\frac{m\hbar\omega}{2}} \sqrt{2n+1}$$

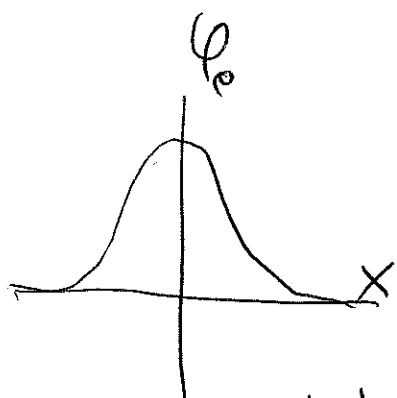
$$\begin{aligned} \text{Then, } \Delta x \Delta p &= \sqrt{\frac{\hbar}{2m\omega}} \sqrt{\frac{m\hbar\omega}{2}} (2n+1) = \\ &= \frac{\hbar}{2} (2n+1) \end{aligned}$$

So, for H.O, the uncertainty relations

$$\text{are: } \frac{\hbar}{2} (2n+1) \geq \frac{\hbar}{2}$$

• Observation 1: ground-state $\Rightarrow n=0$

$$\text{get } \frac{\hbar}{2} = \frac{\hbar}{2} \leftarrow \text{minimal uncertainty!}$$



$$\leftarrow \text{recall } \psi_0(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-m\omega x^2/2\hbar}$$

Ch. 6

minimal uncertainty wavepacket $\leftarrow$ Gaussian !!

↑ is this expected?

• Observation 2 : as $n \uparrow \Rightarrow$

(6)

the uncertainty $\uparrow$

harder $\swarrow$ to make precise measurements
for excited-states