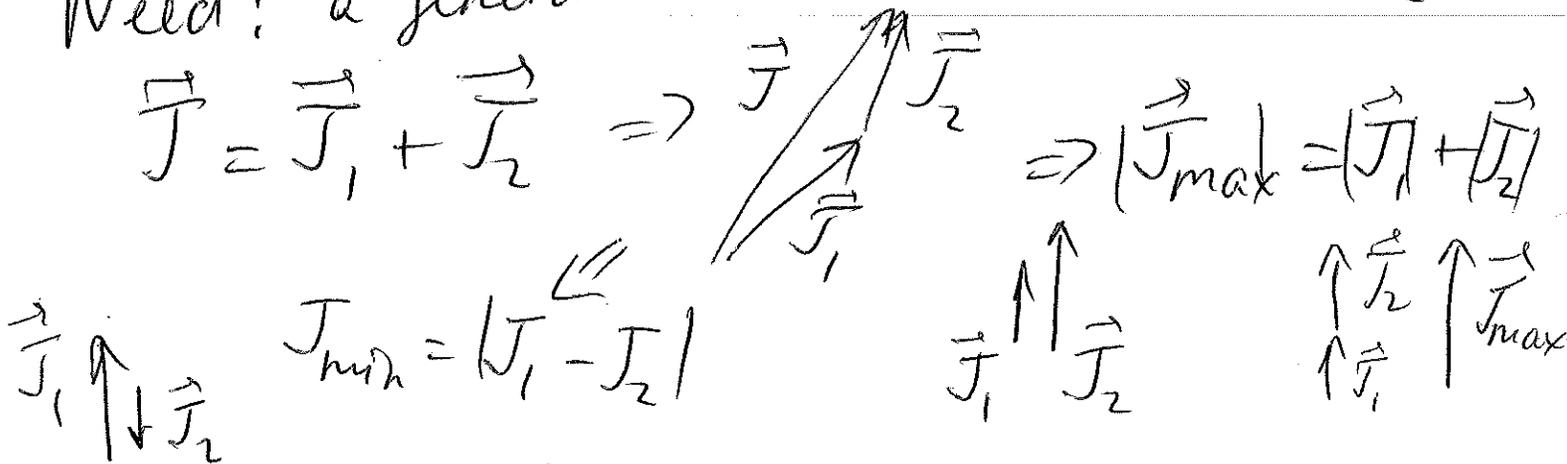


Addition of generalised angular momenta.

Need: a general rule of adding \vec{J}_1, \vec{J}_2

$\vec{J} = \vec{J}_1 + \vec{J}_2 \Rightarrow \vec{J} \Rightarrow |\vec{J}_{max}| = |\vec{J}_1| + |\vec{J}_2|$



Classically: $|\vec{J}|$ varies from $||\vec{J}_1| - |\vec{J}_2||$ to $|\vec{J}_1| + |\vec{J}_2|$

QM: in quantized steps: $J = j_1 + j_2, j_1 + j_2 - 1, \dots$

j_1, j_2 : quantum number correspondingly to \vec{J}_1, \vec{J}_2
 quantum number correspond to J

$-j_1 \leq m_1 \leq j_1$
 $-j_2 \leq m_2 \leq j_2$
 $M = m_1 + m_2$
 $M: -J \leq M \leq J$
 $-J, -J+1, -J+2, \dots, J$

Ex. $j_1 = 2; j_2 = 3 \rightarrow J = ?$

$J = \underbrace{2+3}_5, 4, 3, 2, 1$

Note that $J \geq 0!$

$|2-3| = 1$

$$\{ \vec{J}_1^2, \vec{J}_2^2, J_{1z}, J_{2z} \} \Rightarrow |j_1 j_2 m_1 m_2\rangle \quad (2)$$

↑
set of commuting
↓
observables

↕ uncoupled
basis

$$\{ \vec{J}_1^2, \vec{J}_2^2, \vec{J}^2, J_z \} \Rightarrow |j_1 j_2 JM\rangle \equiv |JM\rangle$$

coupled basis

Recall that we established the connection between coupled & uncoupled bases

for a specific case of $j_1 = j_2 = 1/2 \Rightarrow$

$$J = 1, 0 \Rightarrow |11\rangle = |++\rangle \equiv \left| \frac{1}{2} \frac{1}{2} \frac{1}{2} \frac{1}{2} \right\rangle$$

$$|10\rangle = \frac{1}{\sqrt{2}} (|+-\rangle + |-+\rangle) \equiv$$

$$\equiv \frac{1}{\sqrt{2}} \left(\left| \frac{1}{2} \frac{1}{2} \frac{1}{2} -\frac{1}{2} \right\rangle + \left| \frac{1}{2} \frac{1}{2} -\frac{1}{2} \frac{1}{2} \right\rangle \right)$$

$$|1-1\rangle = |--\rangle \equiv \left| \frac{1}{2} \frac{1}{2} -\frac{1}{2} -\frac{1}{2} \right\rangle$$

$$|00\rangle = \frac{1}{\sqrt{2}} (|+-\rangle - |-+\rangle) \equiv$$

$$\equiv \frac{1}{\sqrt{2}} \left(\left| \frac{1}{2} \frac{1}{2} \frac{1}{2} -\frac{1}{2} \right\rangle - \left| \frac{1}{2} \frac{1}{2} -\frac{1}{2} \frac{1}{2} \right\rangle \right)$$

Is there a generalized approach?

↓

yes! \Rightarrow

What is the dimensionality of the 3 uncoupled & coupled bases? (must match!)

\Downarrow

$|j_1 j_2 m_1 m_2\rangle \Rightarrow N_{j_1 j_2} = \underbrace{(2j_1+1)(2j_2+1)}_{\substack{\text{dimensionality} \\ \downarrow \text{of basis}}} \xrightarrow{\text{change basis for convenience}} \text{should not affect physics}$

\Downarrow

for $j_{1,2}$ there are $2j_{1,2}+1$ states ($|j_1 m_1\rangle$ or $|j_2 m_2\rangle$)

On the other hand, for coupled basis \Rightarrow

$$|JM\rangle \Rightarrow N_J = 2J+1$$

$N_{j_1 j_2} \stackrel{\text{show!}}{=} N_J = \sum_{J=|j_1-j_2|}^{j_1+j_2} (2J+1)$

same dimensionality!

$(2j_1+1)(2j_2+1)$

Any vector $|JM\rangle$ can be presented as a superposition of $|j_1 j_2 m_1 m_2\rangle$ vectors

$\underbrace{\hspace{10em}}$ form complete & orthonormal basis

How does completeness look like here? (4)

$$I = \sum_{m_1=-j_1}^{j_1} \sum_{m_2=-j_2}^{j_2} |j_1 j_2 m_1 m_2\rangle \langle j_1 j_2 m_1 m_2|$$

↑
identity

$$\begin{aligned} \text{So: } |JM\rangle &= I |JM\rangle = \\ &= \sum_{m_1} \sum_{m_2} |j_1 j_2 m_1 m_2\rangle \langle j_1 j_2 m_1 m_2 | JM \rangle \end{aligned}$$

relate coupled & uncoupled bases

← Clebsch-Gordan coefficients

↑ $j_1 j_2 J$
 $m_1 m_2 M$

↑ Tabulated, also can be derived

Properties of C.-G. coefficients:

- 1) $\langle j_1 j_2 m_1 m_2 | JM \rangle \neq 0$ only if $M = m_1 + m_2$ & $|j_1 - j_2| \leq J \leq j_1 + j_2$
- 2) coeff. are real (convention)
- 3) $\langle j_1 j_2 \underbrace{j_1}_{\substack{\text{max} \\ \text{value}}} \underbrace{J-j_1}_{m_2} | \underbrace{J}_{M=J} \rangle > 0$ (convention)

Example

$$\begin{matrix} \leftarrow F & \leftarrow M_F \\ J & M \end{matrix}$$

$$|00\rangle = \frac{1}{\sqrt{2}} \left(|+\!-\rangle - |-\!+\rangle \right) \equiv \frac{1}{\sqrt{2}} \left(\left| \frac{1}{2} \frac{1}{2} \frac{1}{2} -\frac{1}{2} \right\rangle - \left| \frac{1}{2} \frac{1}{2} -\frac{1}{2} \frac{1}{2} \right\rangle \right)$$

Recall: $j_1 = j_2 = 1/2$ case (5)
 (from hyperfine discussion $S = I = 1/2$)

C.-G. coeff. Based on this:

$$\left\langle \frac{1}{2} \frac{1}{2} \frac{1}{2} -\frac{1}{2} \mid 00 \right\rangle = \frac{1}{\sqrt{2}}$$

$$\left\langle \frac{1}{2} \frac{1}{2} -\frac{1}{2} \frac{1}{2} \mid 00 \right\rangle = -\frac{1}{\sqrt{2}}$$

These must be > 0 by convention

(also; see the Table)

Why is it $(|+\!-\rangle - |-\!+\rangle)$ and not

$(|-\!+\rangle - |+\!-\rangle)$ (given that the global phase doesn't matter)

convention \Rightarrow
 (C.-G. property #3)

$$\left\langle \frac{1}{2} \frac{1}{2} \frac{1}{2} -\frac{1}{2} \mid 00 \right\rangle > 0$$

$$\begin{matrix} \uparrow & \uparrow & \uparrow & \uparrow & \uparrow \\ m_1 = j_1 & & m_2 = J - j_1 & & \\ & & \uparrow & \uparrow & \\ & & 0 & 1/2 & \end{matrix}$$

P. 4

Spectroscopic notation

$$\vec{J} = \vec{L} + \vec{S}$$

Specifies atom state

$$2S+1$$

J

total orb.
L = 0 1 2 3

S P D F ...

$$2S_{1/2} \rightarrow L=0, S=1/2, J=1/2$$

For single electron:
l=0 \Rightarrow s
1 p
2 d

Example

A (H) -atom is in a $2p_{1/2}$ state

$$j = l + s, \dots, |l - s|$$

with total ang. momentum up along z-axis. With what probability will the e^- be found with spin down?

So $l + \frac{1}{2}, l - \frac{1}{2}$
if $l \geq 1$ or $\frac{1}{2}$
for $l=0$

Solution! $2P_{1/2} \Rightarrow \left. \begin{matrix} l=1 & (=j_1) \\ s=1/2 & (=j_2) \\ J=1/2, M=1/2 \end{matrix} \right\} \text{given!}$

Need: Prob. that e^- is in the $M_S = -\frac{1}{2}$ state!

$$| \frac{1}{2} \frac{1}{2} \rangle = \frac{1}{\sqrt{3}} | \frac{1}{2} \frac{1}{2} 1 -\frac{1}{2} \rangle - \frac{1}{\sqrt{3}} | \frac{1}{2} \frac{1}{2} 0 \frac{1}{2} \rangle \Rightarrow P = \frac{2}{3}$$