

PH 457/557  
QM

# Lecture # 14

①

## Angular momentum ladder operators

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Introduce  $J_{\pm} = J_x \pm iJ_y$

$$(J_+)^{\dagger} = J_- \leftarrow \text{not Hermitian}$$

$$[J_+, J_-] = 2\hbar J_z, \quad [\vec{J}^2, J_{\pm}] = 0$$

(WS 13)  $[J_x, J_y] = i\hbar J_z$

What do  $J_{\pm}$  do to the state  $|j m_j\rangle$ ?

Consider  $\vec{J}^2 J_{\pm} |j m_j\rangle = J_{\pm} \vec{J}^2 |j m_j\rangle =$   
recall  $\vec{J}^2$  acts on  $|j m_j\rangle$  as  $\hbar^2 j(j+1)$   
 $= \hbar^2 j(j+1) J_{\pm} |j m_j\rangle \Rightarrow J_{\pm} \text{ do not change } j$

HW #5  $\Rightarrow$  show  $[J_z, J_{\pm}] = \pm \hbar J_{\pm}$   
(here we will use the result)  $\Rightarrow$

$$J_z J_{\pm} |j m_j\rangle = J_{\pm} J_z |j m_j\rangle \pm \hbar m_j |j m_j\rangle \quad (2)$$

$$\pm \hbar J_{\pm} |j m_j\rangle = \hbar (m \pm 1) J_{\pm} |j m_j\rangle$$

$$J_z |? \rangle = \hbar (m \pm 1) |? \rangle$$

$$\Rightarrow c |j m \pm 1\rangle$$

$$\text{HW \#5} \Rightarrow J_- J_+ = J^2 - J_z^2 - \hbar J_z$$

$$\langle j m_j | J_- J_+ | j m_j \rangle = \langle j m_j | J_+ J_- | j m_j \rangle$$

$$= \| J_+ | j m_j \rangle \|^2 = \langle j m_j | J^2 - J_z^2 - \hbar J_z | j m_j \rangle$$

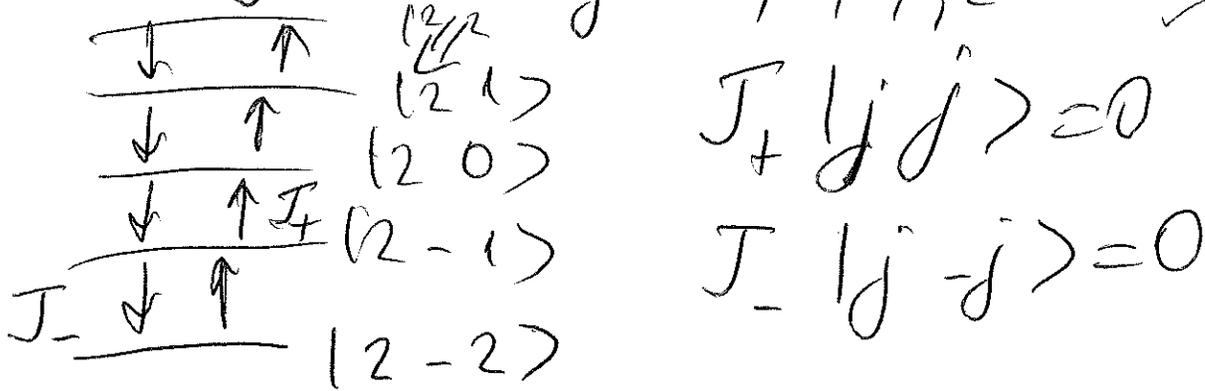
$$= \hbar^2 j(j+1) - \hbar^2 m^2 - \hbar^2 m$$

$$= \hbar^2 (j(j+1) - m(m+1)) = c^2 \Rightarrow$$

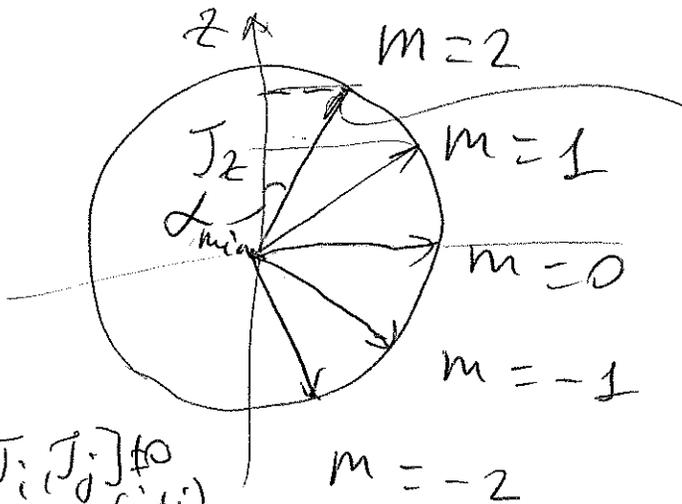
$$c = \hbar \sqrt{j(j+1) - m(m+1)}$$

So,  $\hat{J}_\pm |j, m_j\rangle = \hbar \sqrt{j(j+1) - m_j(m_j \pm 1)}$  ③

E.g.  $j=2 \Rightarrow m_j = -2, -1, 0, 1, 2$   $|j, m_j \pm 1\rangle$



Another visualisation:



$|\vec{J}| = \hbar \sqrt{j(j+1)} = \sqrt{6} \hbar$

$J_\pm \Rightarrow$  don't rescale  $\vec{J}$ , but step through available projections on  $z$

$[J_i, J_j] = i\hbar \epsilon_{ijk} J_k$   
can't know all comp.

Note:  $|J_z| < |\vec{J}| \Rightarrow m^2 < j(j+1)$

there is a minimal angle  $\alpha_{min} \Rightarrow \cos \alpha_{min} = \frac{J_z}{|\vec{J}|} = \frac{m_{max}}{\sqrt{j(j+1)}} = \frac{j}{\sqrt{j(j+1)}}$

if  $\alpha_{min} = 0 \Rightarrow J_x = J_y = 0, J_z = \pm \hbar j$   $\leftarrow$  Principle of uncertainty  $\leftarrow$  know all comp.

# Hyperfine interaction

(4)

$$\vec{\mu}_e = -g_e \frac{e}{2m_e} \vec{S} = -g_e \underbrace{\mu_B}_{\substack{\text{Bohr magneton} \\ \mu_B = \frac{e\hbar}{2m_e}}} \frac{\vec{S}}{\hbar}$$

↑  
magnetic moment of  $e^-$

↑  
gyromag. ratio ( $\approx 2$ )

$$\vec{\mu}_p = g_p \frac{e}{2m_p} \vec{I} = g_p \underbrace{\mu_N}_{\substack{\text{nuclear magneton} \\ \mu_N = \frac{e\hbar}{2m_p}}} \frac{\vec{I}}{\hbar}$$

↑  
magnetic moment of proton

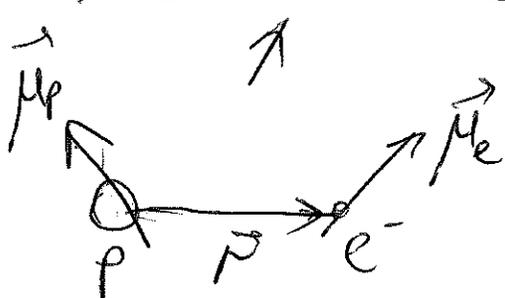
↑  
nuclear magneton

Note:  $\mu_N \ll \mu_B$

$$H'_{hf} = \vec{\mu}_p \cdot \frac{\mu_0}{4\pi} \frac{e}{m_p^3} \vec{L} + \frac{\mu_0}{4\pi} \frac{1}{r^3} \left[ \vec{\mu}_e \cdot \vec{\mu}_p - \frac{3(\vec{\mu}_e \cdot \vec{r})(\vec{\mu}_p \cdot \vec{r})}{r^2} \right] - \frac{\mu_0}{4\pi} \frac{8\pi}{3} \vec{\mu}_e \cdot \vec{\mu}_p \delta(\vec{r})$$

0 for s-states

← orb. ang. mom. of  $e^-$



dipole-dipole interaction  
(like in ESM)

↑ less obvious, but also 0 for s-states

Fermi's contact term (when  $e^-$  is very close to p)

↓  
s-states  
( $l=0$ )

We will focus on  $H'_{hf}$  effect on (5)  
 the 1s state of (H)-atom (ground state)  
 $\uparrow$   
 $|100\rangle$

$$H'_{hf} = - \frac{\mu_0}{4\pi} \frac{8\pi}{3} \underbrace{\vec{\mu}_e \cdot \vec{\mu}_p}_{\text{}} \delta(\vec{r}) = \frac{\mu_0}{4\pi} \frac{g_e g_p \mu_B \mu_N}{\hbar^2} \delta(\vec{r}) \vec{S} \cdot \vec{I}$$

Next, use  $H'_{hf}$  as a perturbation to  $H_0 \Rightarrow$

$\frac{\hbar \pm 1}{|100\rangle} \Rightarrow ?$  Note:  $H'_{hf} \sim \vec{S} \cdot \vec{I}$   
 $H_0 + H'_{hf}$   
 $\downarrow$   
 acts on spin variables for  $e^-$  &  $p^+$ !

For a complete description of the state, need

$$|100\rangle \otimes |S m_S\rangle \otimes |I m_I\rangle \quad \text{or, simply,}$$

$\uparrow \quad \uparrow \quad \uparrow \quad \uparrow$   
 $1/2 \quad \pm 1/2 \quad 1/2 \quad \pm 1/2$   
 for electron      for proton

$|100; m_S m_I\rangle \leftarrow$  4-fold degenerate