

WHY DFT? Besides $E(k)$ and DoS?

Equilibrium crystal structures

Elastic constants

Charge densities

Vibrational properties

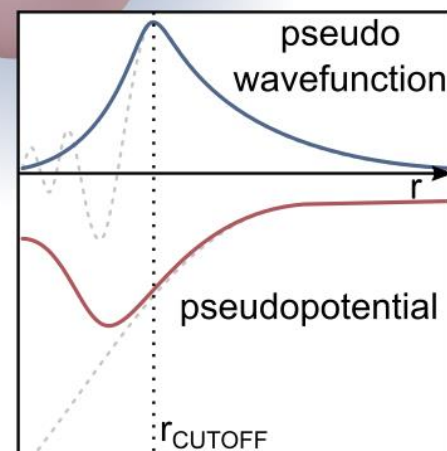
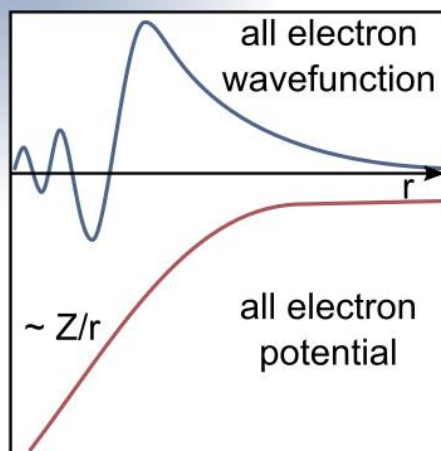
Qualitative prediction of band topology

How to Do DFT with Quantum Espresso?

- 1) Understand your structure
- 2) Obtain (generate or steal) pseudopotentials
- 3) Test your pseudopotentials on known systems (if possible)
- 4) Test your system for convergence in PW energy cutoff
- 5) Test your system for convergence in k-point sampling
- 6) Do project-specific calculations and associated tests



replacing inner electrons
with pseudopotential



What will almost always get wrong?

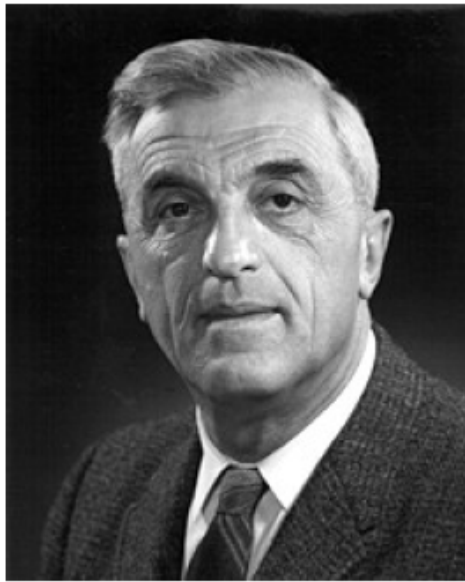
Excited state properties:

- Band gaps [can do GW instead]
- Optical properties (electron–hole interactions) [can do BSE]

Other missing bits and trouble areas:

- Overbinding in LDA, underbinding in GGA
- Strongly correlated systems (narrow d– and f– bands)
- Lack of van der Waals interactions [can be added through the use of various schemes in addition to DFT]
- Presence of self–interaction error (electron interacts with its own charge density) [reduced by using hybrid functionals with some fraction of exact–exchange, but expensive]

- Wavefunctions
- Collective motions of the atoms



Bloch

$$\psi_{\mathbf{k}}(\mathbf{r} + \mathbf{T}) = e^{i\mathbf{k} \cdot \mathbf{T}} \psi_{\mathbf{k}}(\mathbf{r})$$

equivalently:

$$\psi_{\mathbf{k}}(\mathbf{r}) = e^{i\mathbf{k} \cdot \mathbf{r}} u_{\mathbf{k}}(\mathbf{r})$$

$$u_{\mathbf{k}}(\mathbf{r} + \mathbf{T}) = u_{\mathbf{k}}(\mathbf{r})$$

Allowed form of the Bloch vectors:

$$\mathbf{k} = \sum_{i=1}^3 \frac{m_i}{N_i} \mathbf{b}_i$$

where m_i are integers

We say the solutions are “Plane Waves”
→ requires periodic boundary conditions

Bloch form of solution:

$$\psi_{\mathbf{k}}(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}} u_{\mathbf{k}}(\mathbf{r})$$

can be satisfied by expansion of the form:

$$\phi_{\mathbf{k}}(\mathbf{r}) = \frac{e^{i\mathbf{k}\cdot\mathbf{r}}}{\sqrt{V}} \sum_{\mathbf{G}=0}^{\infty} c_{\mathbf{k}}(\mathbf{G}) e^{i\mathbf{G}\cdot\mathbf{r}}$$

AB INITIO SOFTWARE

Plane Wave Software (for periodic system)

Quantum Espresso

VASP

Requires PERIODIC boundary conditions to work (not great for molecules)

→ Solution for molecules make the unit cell so large (typically 1.5 nm is enough) that the atoms in the cell can't possibly interact (how to do this)
<https://www.youtube.com/watch?v=Z8TEpaW7RNc>)

Atomic based (site-based)

Gaussian

QChem


The 1st calculation iteration is completed when all states are filled below a cutoff energy

$$\phi_{n\mathbf{k}}(\mathbf{r}) = \sum_{\mathbf{G}=0}^{\infty} c_{n\mathbf{k}}(\mathbf{G}) \phi_{\mathbf{G}}^{\mathbf{k}}(\mathbf{r})$$

$$\phi_{\mathbf{G}}^{\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{V}} e^{i(\mathbf{k}+\mathbf{G})\cdot\mathbf{r}}$$

Coefficients decay with increasing wavevector magnitude, and thus can be truncated without loss of accuracy:

$$\frac{\hbar^2}{2m} |\mathbf{k} + \mathbf{G}|^2 < E_{\text{cut}}$$



Plane wave energy cutoff, typically expressed in Ry (e.g., QE) or eV (e.g. VASP)

- Many-body wavefunction (keep it simple: only 3 electrons)

$$\psi(\mathbf{r}) \rightarrow \Psi = \Psi(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3)$$

- Probability of finding electron #1 at the point \mathbf{r}

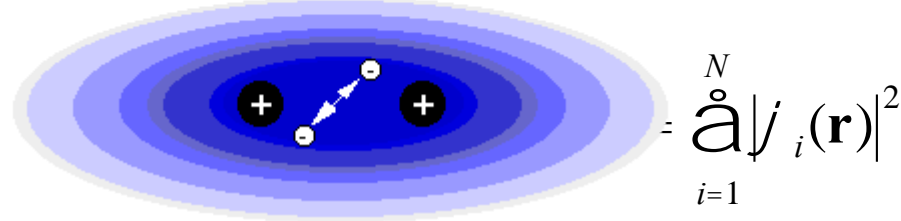
$$\text{prob}(\mathbf{r}_1 = \mathbf{r}) = \int |\Psi(\mathbf{r}, \mathbf{r}_2, \mathbf{r}_3)|^2 d\mathbf{r}_2 d\mathbf{r}_3$$

- Electron density at the point \mathbf{r}

$$n(\mathbf{r}) = \text{prob}(\mathbf{r}_1 = \mathbf{r}) + \text{prob}(\mathbf{r}_2 = \mathbf{r}) + \text{prob}(\mathbf{r}_3 = \mathbf{r})$$

- Electrons are indistinguishable

$$n(\mathbf{r}) = 3 \int |\Psi(\mathbf{r}, \mathbf{r}_2, \mathbf{r}_3)|^2 d\mathbf{r}_2 d\mathbf{r}_3$$



Kohn and Sham said the (Helmholtz) Energy of the system is:

$$F_{HK}[\rho] = F_{KS}[\rho] = T_s[\rho] + \frac{1}{2} \int \int \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r} d^3\mathbf{r}' + E_{xc}[\rho]$$

$$= T_s[\rho] + E_{Hartree}[\rho] + E_{xc}[\rho]$$

1) Kinetic energy of the system of **non-interacting** electrons at the same density.

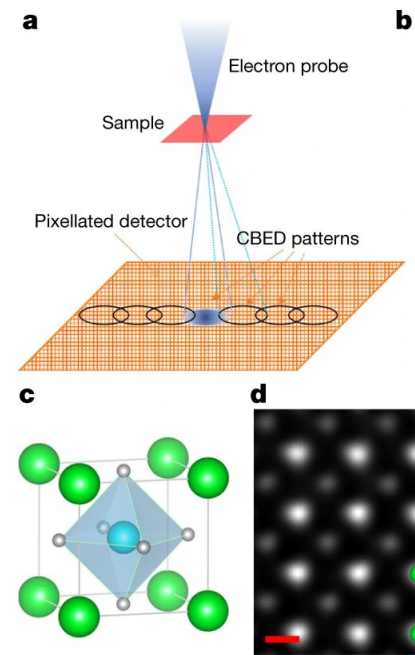
2) Coulomb is the electrostatic term (Hartree)

3) Exchange-correlation is everything else $V_{xc}[\mathbf{r}] = \frac{\partial E_{xc}[\rho(\mathbf{r})]}{\partial \rho(\mathbf{r})}$

How would you compute the total charge density of a crystal of N electrons?

$$\begin{aligned}
 n(\mathbf{r}) &= \int \Psi^*(\mathbf{r}, \mathbf{r}_2, \dots, \mathbf{r}_N) \Psi(\mathbf{r}, \mathbf{r}_2, \dots, \mathbf{r}_N) d^3r_2 d^3r_3 \dots d^3r_N + \\
 &+ \int \Psi^*(\mathbf{r}_1, \mathbf{r}, \dots, \mathbf{r}_N) \Psi(\mathbf{r}_1, \mathbf{r}, \dots, \mathbf{r}_N) d^3r_1 d^3r_3 \dots d^3r_N + \dots \\
 &+ \int \Psi^*(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}) \Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}) d^3r_1 d^3r_2 d^3r_3 \dots d^3r_{N-1} = \\
 &= \int (\delta(\mathbf{r} - \mathbf{r}_1) + \delta(\mathbf{r} - \mathbf{r}_2) + \dots + \delta(\mathbf{r} - \mathbf{r}_N)) \\
 &\quad \Psi^*(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N) \Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N) d^3r_1 d^3r_2 d^3r_3 \dots d^3r_N = \\
 &= \sum_{i=1}^N \int \langle \Psi | \mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N \rangle \delta(\mathbf{r} - \mathbf{r}_i) \langle \mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N | \Psi \rangle d^3r_1 d^3r_2 d^3r_3 \dots d^3r_N = \\
 &= N \int \langle \Psi | \mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N \rangle \delta(\mathbf{r} - \mathbf{r}_1) \langle \mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N | \Psi \rangle d^3r_1 d^3r_2 d^3r_3 \dots d^3r_N =
 \end{aligned}$$

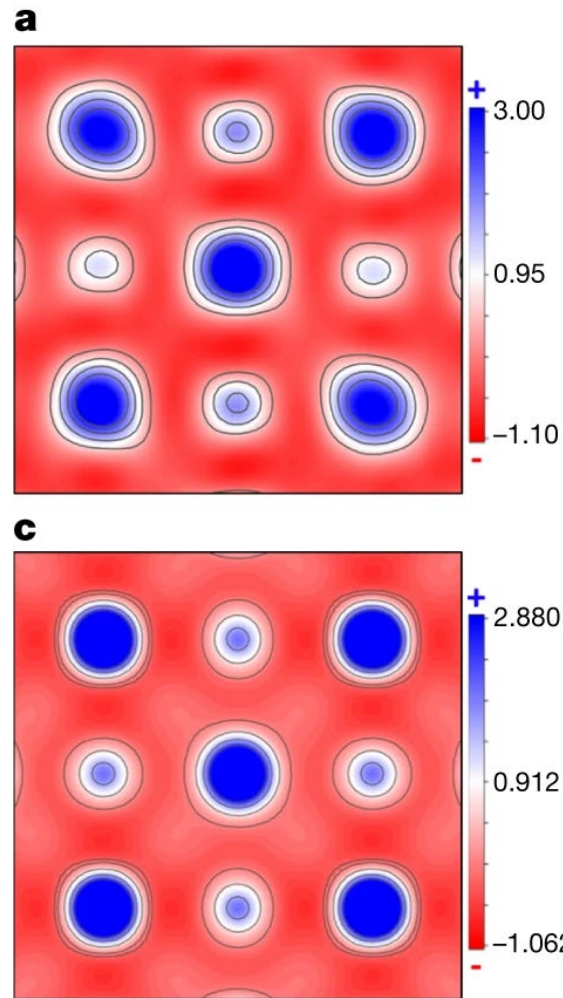
ELECTRON DENSITY MAPS, $\rho(r)$



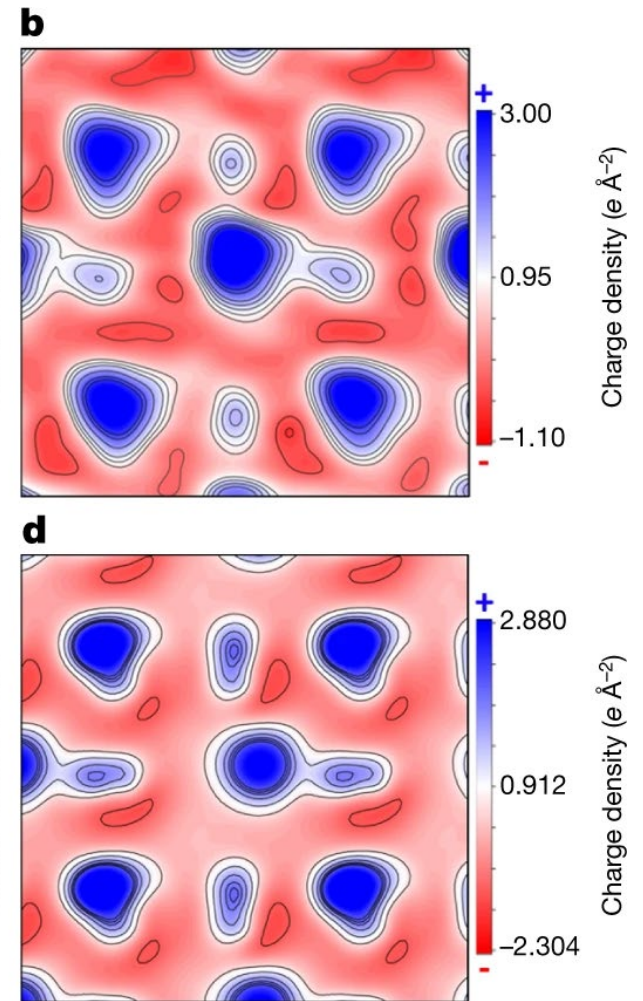
EXPERIMENT

THEORY (DFT)

SrTiO₃

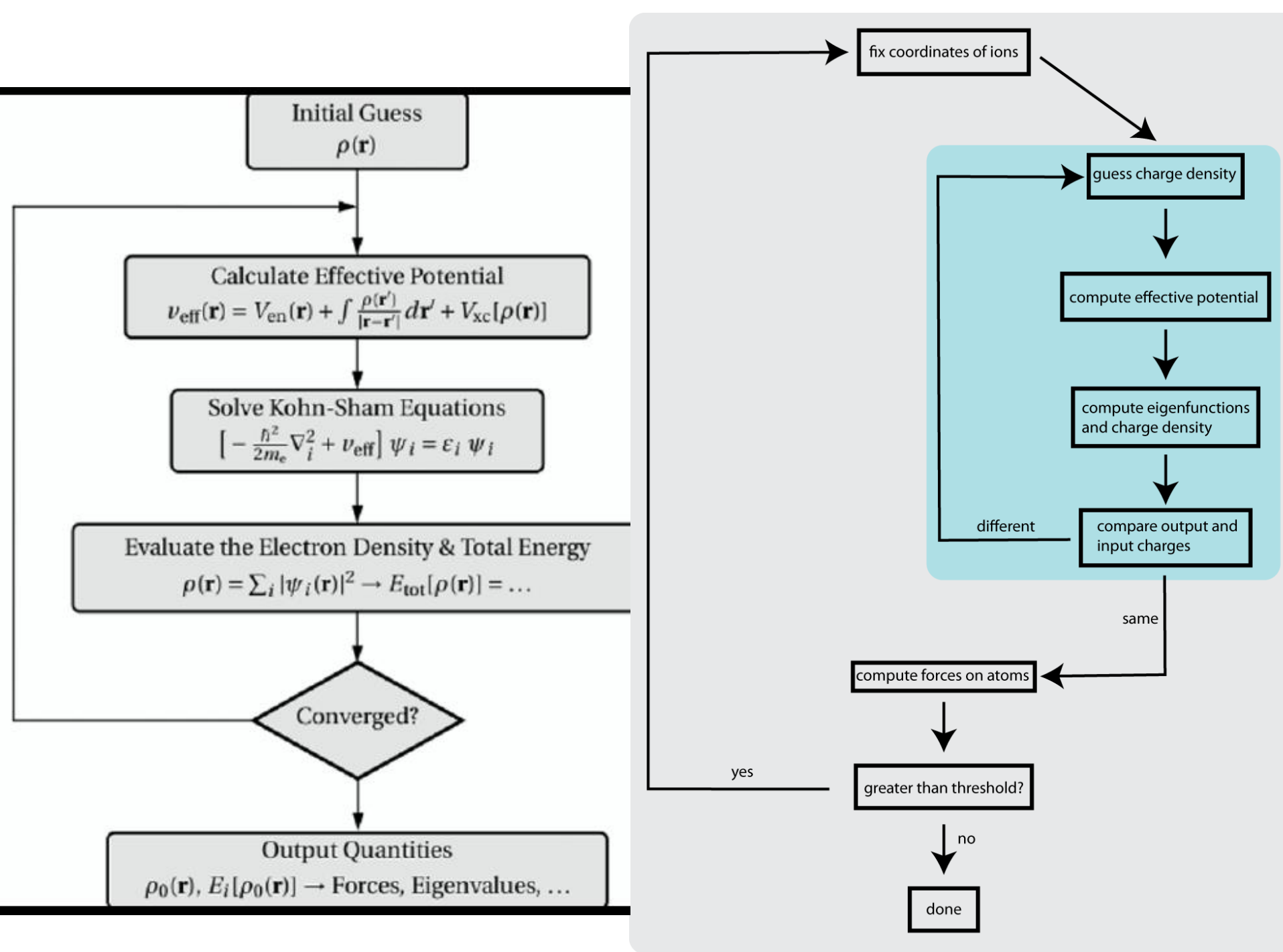


BiFeO₃



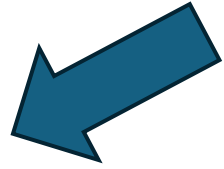
Real-space charge-density imaging with sub-ångström resolution by four-dimensional electron microscopy. *Nature* **575**, 480–484 (2019).

<https://doi.org/10.1038/s41586-019-1649-6>



$$F_{HK}[\rho] = F_{KS}[\rho] = T_s[\rho] + \frac{1}{2} \int \int \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r} d^3\mathbf{r}' + E_{xc}[\rho]$$

$$= T_s[\rho] + E_{Hartree}[\rho] + E_{xc}[\rho]$$

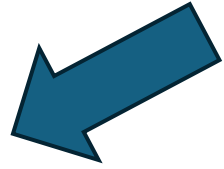


Take the functional derivative wrt ρ ,
and evaluating at LCAO orbitals gives
the Kohn-Sham equations

$$\left[-\frac{\hbar^2}{2m} + V_{ks}[\rho](\mathbf{r}) \right] \phi_i(\mathbf{r}) = \epsilon_i \phi_i(\mathbf{r})$$

$$F_{HK}[\rho] = F_{KS}[\rho] = T_s[\rho] + \frac{1}{2} \int \int \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r} d^3\mathbf{r}' + E_{xc}[\rho]$$

$$= T_s[\rho] + E_{Hartree}[\rho] + E_{xc}[\rho]$$



Take the functional derivative wrt ρ ,
and evaluating at LCAO orbitals gives
the Kohn-Sham equations

$$\left[-\frac{\hbar^2}{2m} + V_{ks}[\rho](\mathbf{r}) \right] \phi_i(\mathbf{r}) = \varepsilon_i \phi_i(\mathbf{r})$$

where the Kohn-Sham potential is given as:

$$V_{ks}[\rho](\mathbf{r}) = V_{ext}(\mathbf{r}) + V_{Hartree}[\rho](\mathbf{r}) + V_{xc}[\rho](\mathbf{r})$$

$$V_{Hartree}[\rho](\mathbf{r}) = \int \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r}'$$

$$V_{xc}[\rho](\mathbf{r}) = \frac{\delta E_{xc}[\rho]}{\delta \rho(\mathbf{r})}$$

WHY?

In the Kohn-Sham
equation the atoms are
not interacting (i.e
Schrodinger Equation),
**Instead ρ keeps track of
all interactions!**

V_{xc} : LDA Local Density Approximation

Approximate as that of a homogeneous electron gas where each nuclei replaced by a uniform positively charge background electron density ρ_{gs}

$$V_{xc}(\mathbf{r}; n) = \frac{\delta E_{xc}[n]}{\delta n(\mathbf{r})} = \epsilon_{xc}(\mathbf{r}; n) + \int n(\mathbf{r}') \frac{\delta \epsilon_{xc}(\mathbf{r}'; n)}{\delta n(\mathbf{r})} d^3 r'$$

LDA potential:

$$V_{xc}(\mathbf{r}; n) = \epsilon_{xc}^{LD}(n) + n \frac{d\epsilon_{xc}^{LD}(n)}{dn} = \frac{d}{dn} (n\epsilon_{xc}^{LD}(n)) = V_{xc}^{LD}(n)$$

the electron gas exchange term

$$\epsilon_x^{LD}(n) = -\frac{3}{4\pi} (3\pi^2 n)^{\frac{1}{3}} \longrightarrow V_x^{LD} = -\frac{1}{\pi} (3\pi^2 n)^{\frac{1}{3}} = \frac{4}{3} \epsilon_x^{LD}$$

$$V_c^{LD} = \frac{A}{2} \left\{ \ln \left(\frac{y^2}{Y(y)} \right) + \frac{2b}{Q} \arctan \left(\frac{Q}{2y+b} \right) + \right.$$

$$\left. -\frac{by_0}{Y(y_0)} \left[\ln \left(\frac{(y-y_0)^2}{Y(y)} \right) + \frac{2(b+2y_0)}{Q} \arctan \left(\frac{Q}{2y+b} \right) \right] \right\} +$$

$$-\frac{A}{6} \frac{c(y-y_0) - by_0 y}{(y-y_0)Y(y)}$$

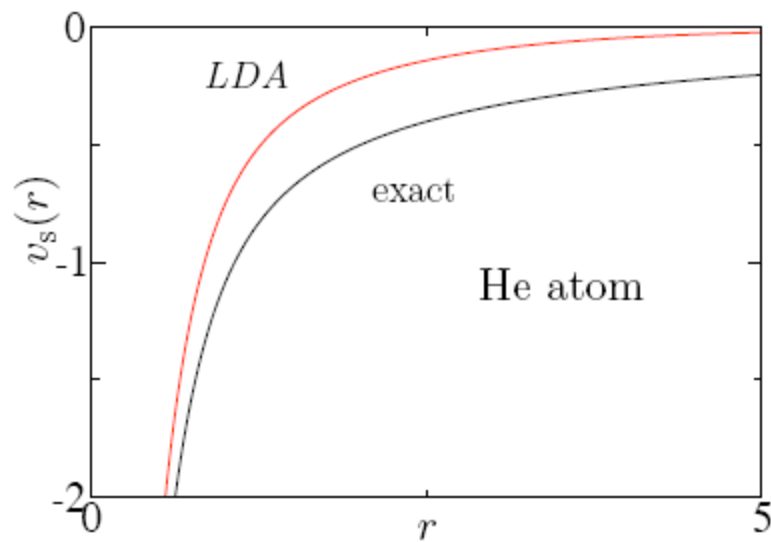


FIG. 7: Exact and LDA KS potentials for the He atom.

LDA potential for the He atom. It decays far too rapidly, and so its orbitals are far too shallow. The HOMO is at -0.5704 H, while the LUMO is not bound at all.

V_{xc} : GGA Generalized Gradient Approximation

**Generalized Gradient Approximation,
the derivative of the charge computed numerically**



$$V_{xc}^{GGA}(\vec{r}) = \frac{\delta E^{GGA}[\rho(\vec{r}'), |\nabla \rho(\vec{r}')|]}{\delta \rho(\vec{r})}$$

$$V_{xc}^{GGA}(\rho(\vec{r}), |\nabla \rho(\vec{r})|, \nabla^2 \rho(\vec{r}), \nabla \rho(\vec{r}) \cdot \nabla |\nabla \rho(\vec{r})|)$$

**Density gradient need not be provided, since they are
calculated numerically using the density at the grid points**

$$\frac{\partial \rho}{\partial x} = \frac{\rho_{i+1} - \rho_{i-1}}{x_{i+1} - x_{i-1}} \Rightarrow E_{xc}^{GGA}(\rho_1, \rho_2, \dots)$$

A finer grid is required for GGA

$$V_{xc}^{GGA}(\vec{r}_i) \equiv \frac{\partial E_{xc}^{GGA}}{\partial \rho_i}$$

III. Exchange-correlation density functional

Performances: GGA vs LDA

Atoms:

Exchange-correlation energies in Ha

| Atom | LSD | GGA | Exact |
|------|--------|--------|--------|
| H | -0.29 | -0.31 | -0.31 |
| He | -1.00 | -1.06 | -1.09 |
| Li | -1.69 | -1.81 | -1.83 |
| Be | -2.54 | -2.72 | -2.76 |
| N | -6.32 | -6.73 | -6.78 |
| Ne | -11.78 | -12.42 | -12.50 |

Molecules:

Atomization energies in Ha

| Molecule | LSD | GGA | Exact |
|------------------|------|------|-------|
| H ₂ | 4.9 | 4.6 | 4.7 |
| CH ₄ | 20.0 | 18.2 | 18.2 |
| NH ₃ | 14.6 | 13.1 | 12.9 |
| H ₂ O | 11.6 | 10.1 | 10.1 |
| CO | 13.0 | 11.7 | 11.2 |
| O ₂ | 7.6 | 6.2 | 5.2 |

Pop science readers



Copenhagen

Many Worlds

QBism

de Broglie-Bohm

Consistent histories

Ensemble

Transactional

Quantum darwinism

Objective-collapse

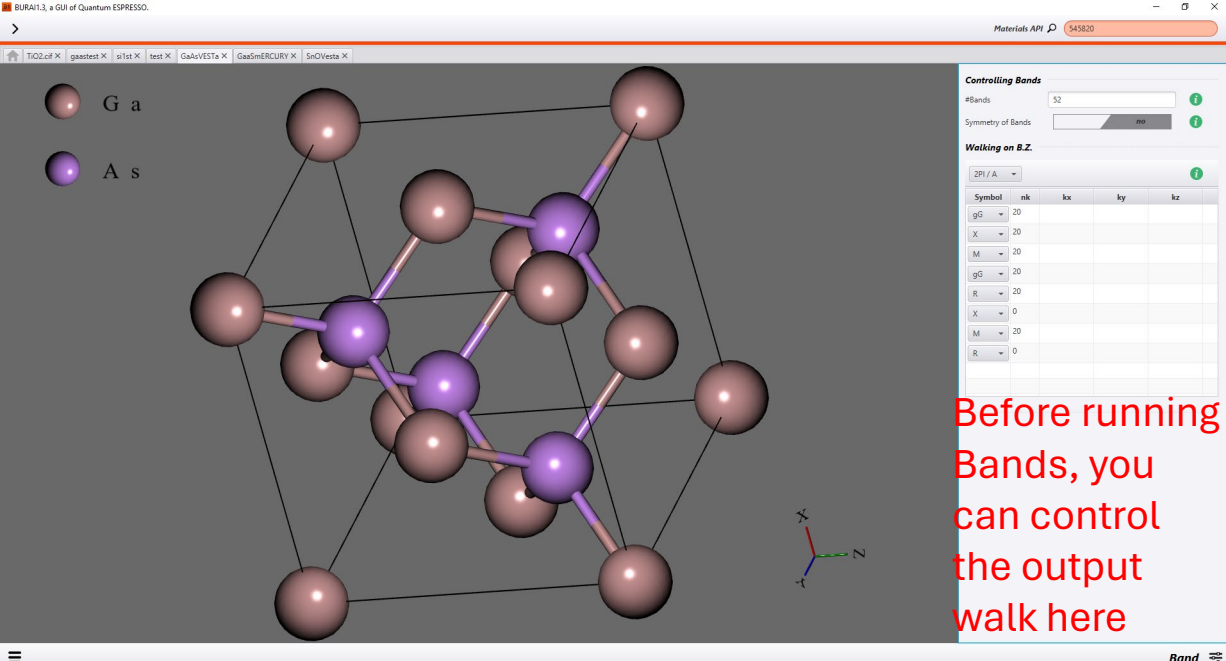
Wigner

Physicists



Shut up and calculate

No one cares



Controlling Bands

#Bands: 52

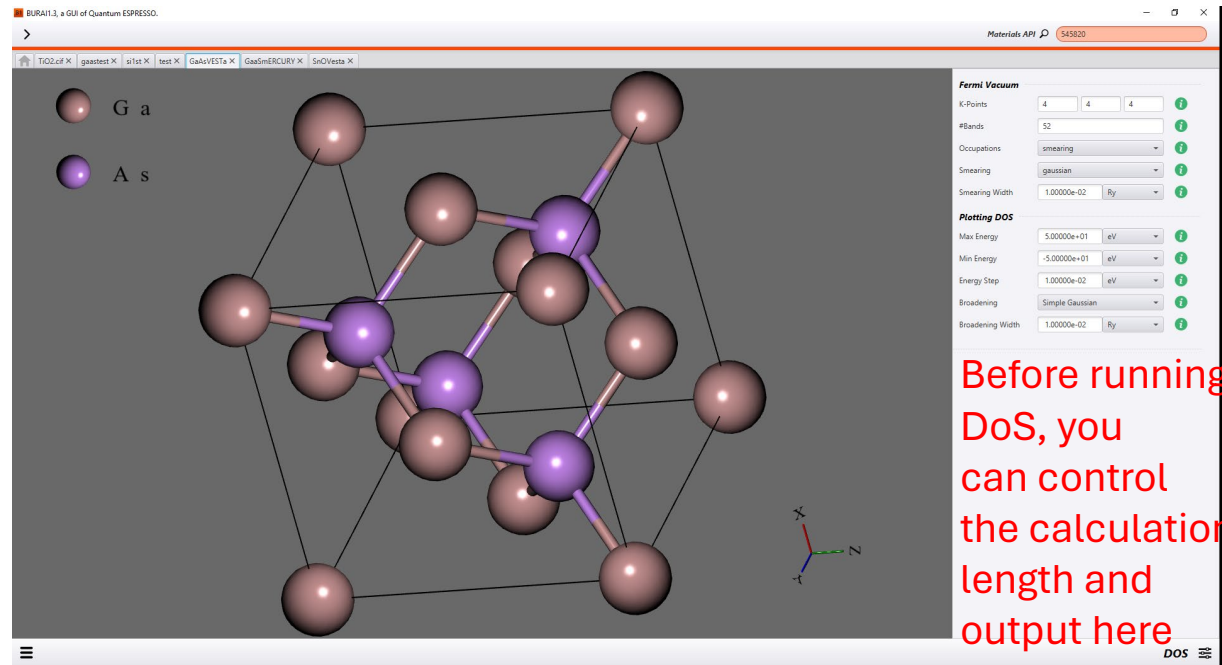
Symmetry of Bands: no

Walking on B.Z.

ZPI / A:

| Symbol | nk | kx | ky | kz |
|--------|----|----|----|----|
| gG | 20 | | | |
| X | 20 | | | |
| M | 20 | | | |
| gG | 20 | | | |
| R | 20 | | | |
| X | 0 | | | |
| M | 20 | | | |
| R | 0 | | | |

Before running Bands, you can control the output walk here



Fermi Vacuum

K-Points: 4 4 4

#Bands: 52

Occupations: smearing

Smearing: gaussian

Smearing Width: 1.00000e-02 Ry

Plotting DOS

Max Energy: 5.00000e+01 eV

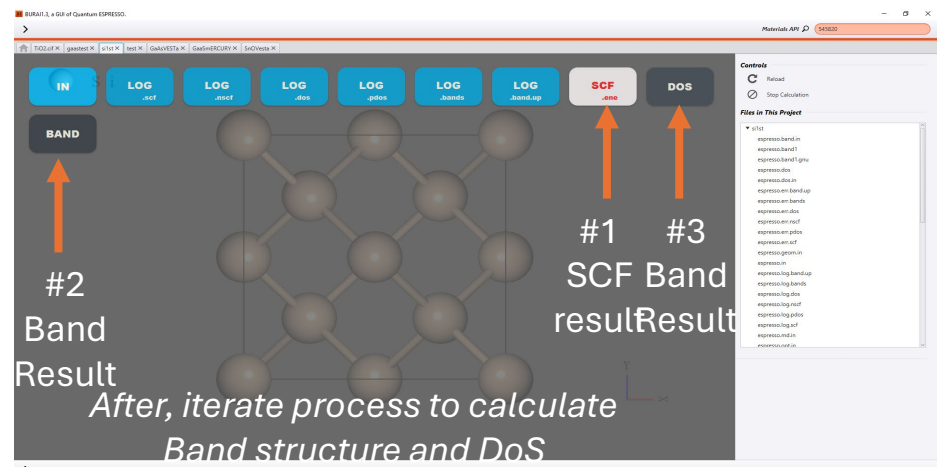
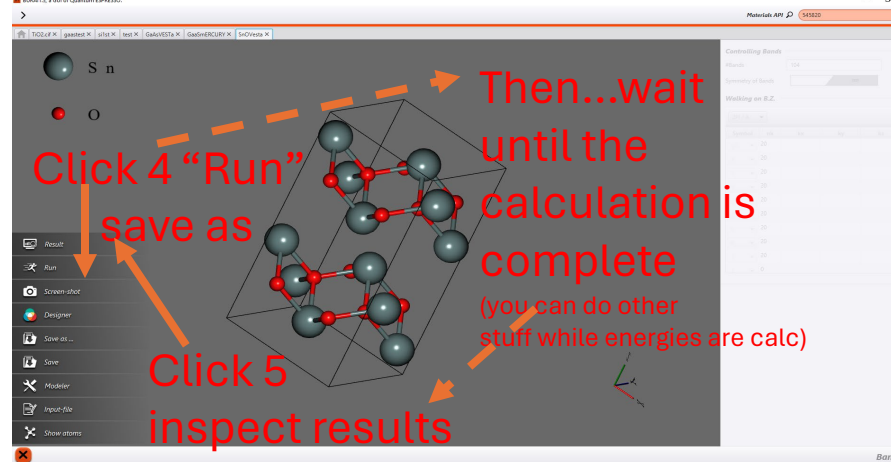
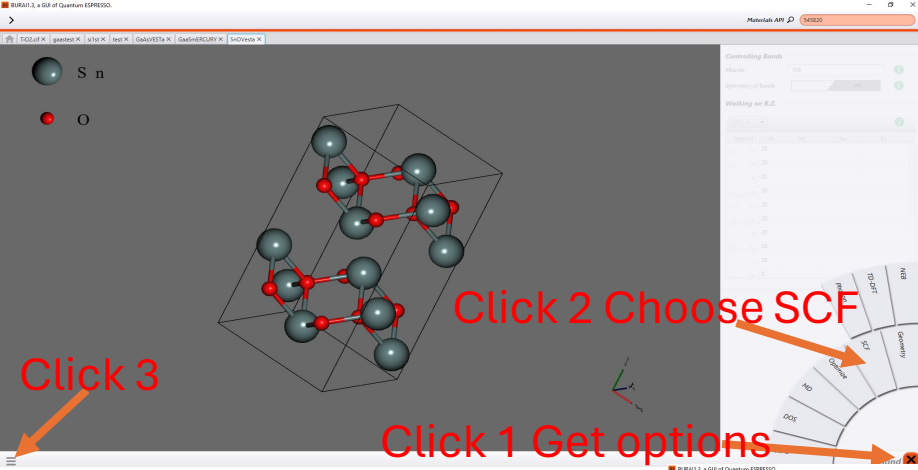
Min Energy: -5.00000e+01 eV

Energy Step: 1.00000e-02 eV

Broadening: Simple Gaussian

Broadening Width: 1.00000e-02 Ry

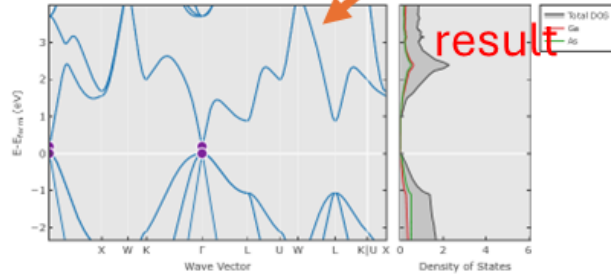
Before running DoS, you can control the calculation length and output here



Electronic Structure

Data Methods API

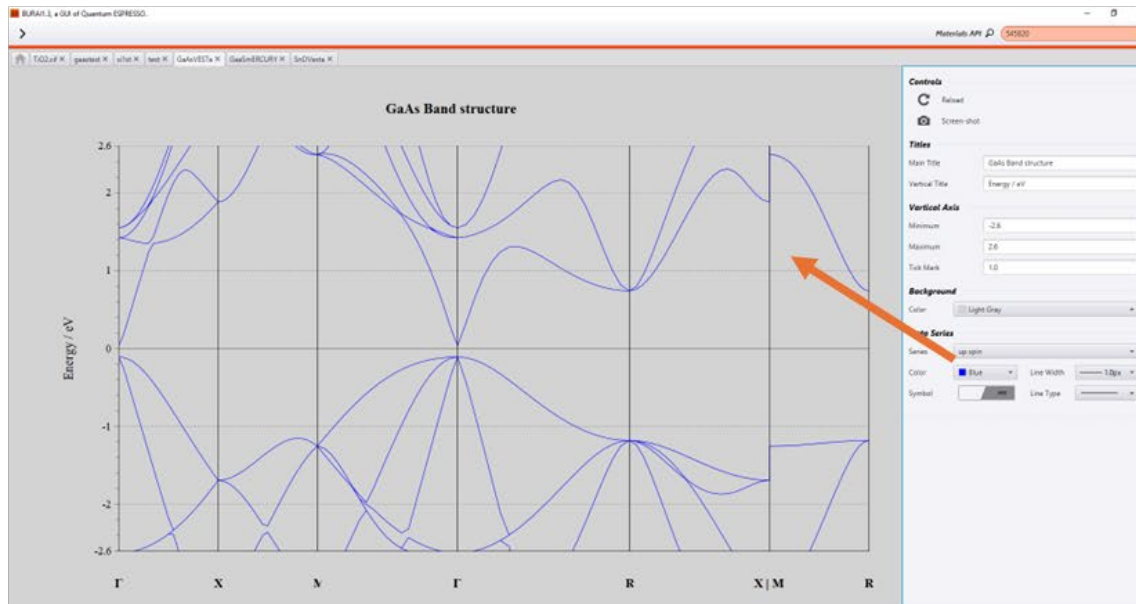
| | |
|------------|---------|
| Band Gap | 0.19 eV |
| Direct Gap | Yes |
| Metallic | No |
| CBM | (0 0 0) |
| VBM | (0 0 0) |



Click electron structure
at the “CIF file site”

Zoom in literature
“interactive graph”
and compare to you
result

*Does your result
make sense with
literature?*



Mostly good! But you took
slight different path through
k-space!
Change or address in your report!

Lastly, load the output file in Excel,
adjust and re-plot in your favorite
Sci-graphing program!