

# Topological Effects

PH671 - Transport

Speakers include

Yakir Aharonov  
 Boris Altshuler  
 Yshai Avishai  
 Michael Berry  
 Markus Buttiker  
 Georgi Dvali  
 Francois Englert  
 Klaus Ensslin  
 Yuval Gefen  
 David Gross\*  
 Moty Heiblum  
 Yoseph Imry  
 Tony Klein \*  
 Charles M. Marcus  
 Yoichiro Nambu\*  
 Murray Peshkin  
 Sandu Popescu  
 Moti Segev  
 Jurgen Smet  
 David Thouless  
 Akira Tonomura  
 Chang C Tsuei  
 Sam Werner  
 Chen Ning Yang  
 \*to be confirmed



# 50 years of the Aharonov-Bohm Effect

## Concepts and Applications

Aharonov-Bohm Effect in:  
 Gauge Theories  
 Scalar form  
 Mesoscopic Systems  
 Quantum Hall Effect  
 Quantum Dots  
 Carbon Nanotubes  
 High-TC Superconductivity

Aharonov-Casher Effect  
 Quantum Geometrical Phases  
 Nonlocal Aspects of the Aharonov-Bohm effect



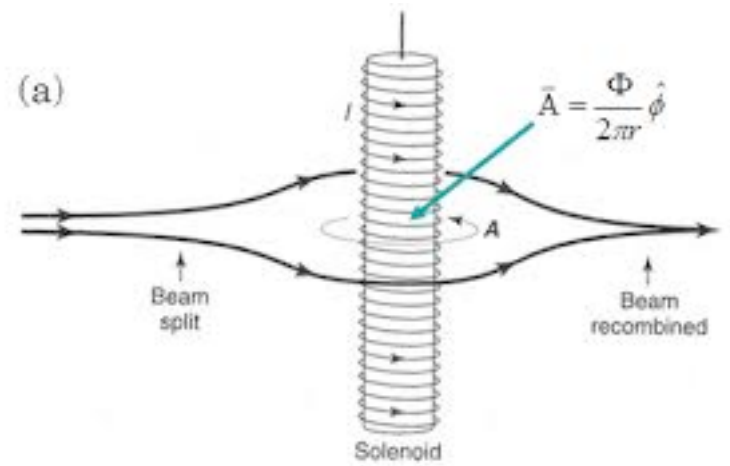
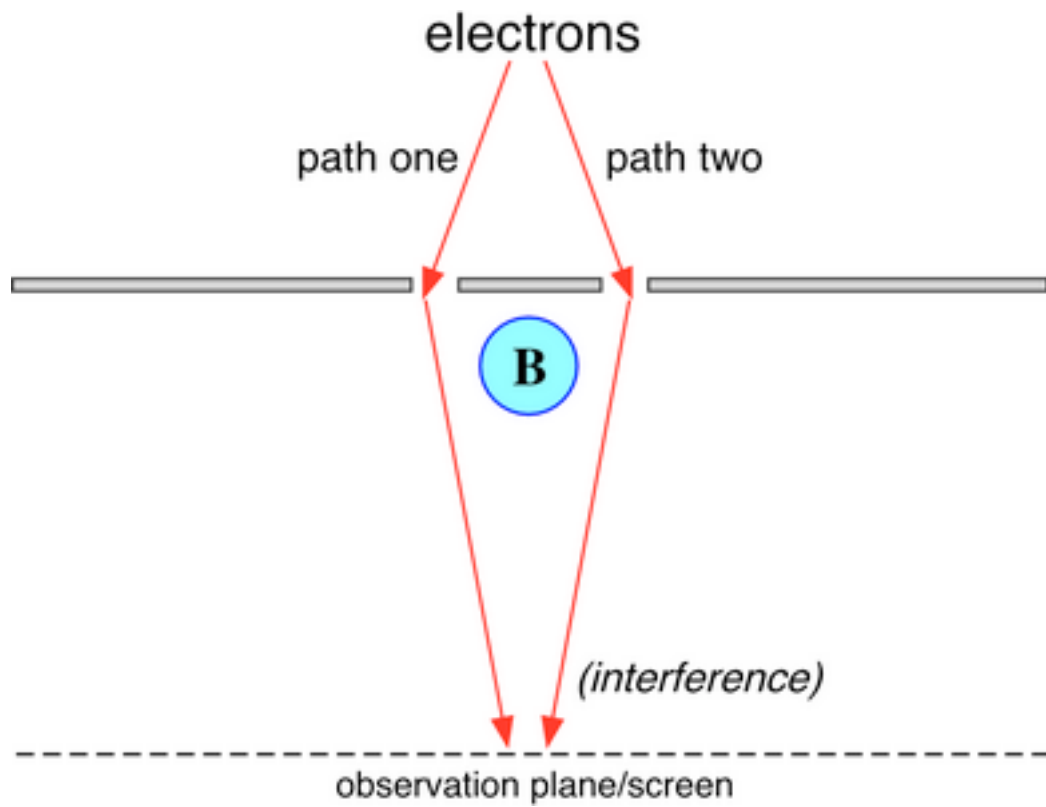
Tel Aviv University, October 11-14, 2009  
<http://www.tau.ac.il/~ab50>  
 A research workshop of the Israel Science Foundation

The Organizing committee  
 Lev Vaidman (Chairman)  
 Eshel Ben Jacob  
 Guy Deutscher  
 Yaron Oz

The Program committee  
 Shmuel Nussinov (Chairman)  
 Yshai Avishai  
 Joseph Avron  
 Aharon Casher  
 Yuval Gefen  
 Benni Reznik

The Advisory committee  
 Jacob Bekenstein  
 Michael Berry  
 Moty Heiblum  
 Anthony Leggett  
 Akira Tonomura  
 Frank Wilczek  
 Joshua Yortner





# Observation of $h/e$ Aharonov-Bohm Oscillations in Normal-Metal Rings

R. A. Webb, S. Washburn, C. P. Umbach, and R. B. Laibowitz

Phys. Rev. Lett. 54, 2696 – Published 24 June 1985

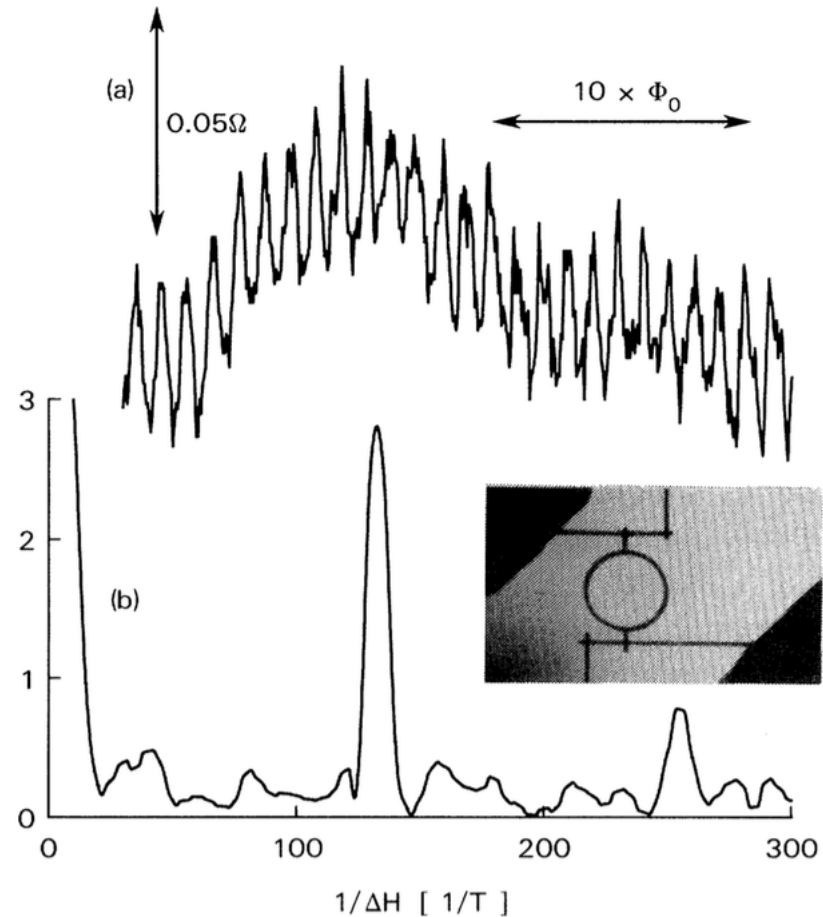


FIG. 1. (a) Magnetoresistance of the ring measured at  $T=0.01$  K. (b) Fourier power spectrum in arbitrary units containing peaks at  $h/e$  and  $h/2e$ . The inset is a photograph of the larger ring. The inside diameter of the loop is 784 nm, and the width of the wires is 41 nm.

# Electromagnetic Aharonov-Bohm effect in a two-dimensional electron gas ring

W. G. van der Wiel, Yu. V. Nazarov, S. De Franceschi, T. Fujisawa, J. M. Elzerman, E. W. G. M. Huizeling, S. Tarucha, and L. P. Kouwenhoven

Phys. Rev. B 67, 033307 –  
Published 29 January 2003

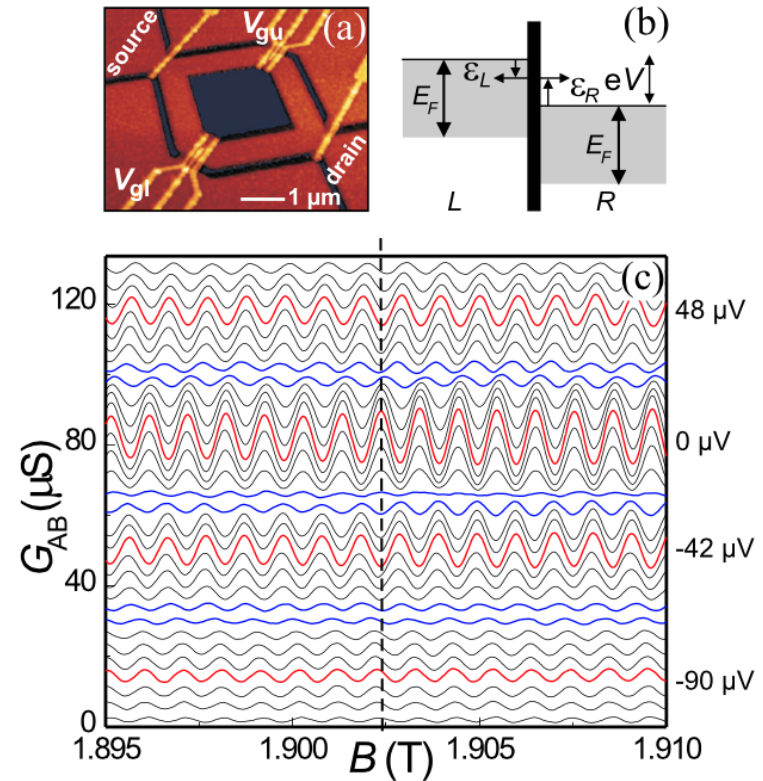


FIG. 1. (a) Atomic force microscope image of the device. The AB ring is defined in a 2DEG by dry etching (dark regions, depth  $\sim 75$  nm). The 2DEG with electron density  $3.4 \times 10^{15} m^{-2}$  is situated 100 nm below the surface of an Al-GaAs/GaAs heterostructure. In both arms of the ring (lithographic width  $0.5 \mu m$ ; perimeter  $6.6 \mu m$ ) a barrier can be defined by applying negative voltages to the gate electrodes  $V_{gu}$  and  $V_{gl}$ . The other gates are not used. (b) Schematic energy picture at one of the barriers (see text). (c) Differential AB conductance,  $G_{AB}$ , as function of  $B$  at 15 mK for different source drain voltages,  $V$ , separated by  $5.3 \mu V$ . The traces have a vertical offset. Three values for  $V$  are indicated on the right. The dashed line highlights a phase change by a change in the amplitude and sign of the AB oscillations, indicating an electrostatic AB effect.



# Aharonov-Bohm effect in charge-density wave loops with inherent temporal current switching

M. Tsubota, K. Inagaki, T. Matsuura and S. Tanda

EPL, 97 (2012) 57011

doi: 10.1209/0295-5075/97/57011

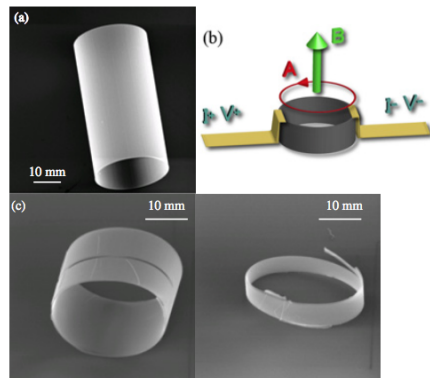


Fig. 1: (Colour on-line) (a) Scanning electron microscope (SEM) image of a TaS<sub>3</sub> tube-shaped crystal. (b) Contact configuration for transport measurement. Two gold electrodes (100 nm thick, 5 μm wide) were evaporated. (c) SEM image of a TaS<sub>3</sub> tube/ring crystal before cutting, (left) after cutting (right).

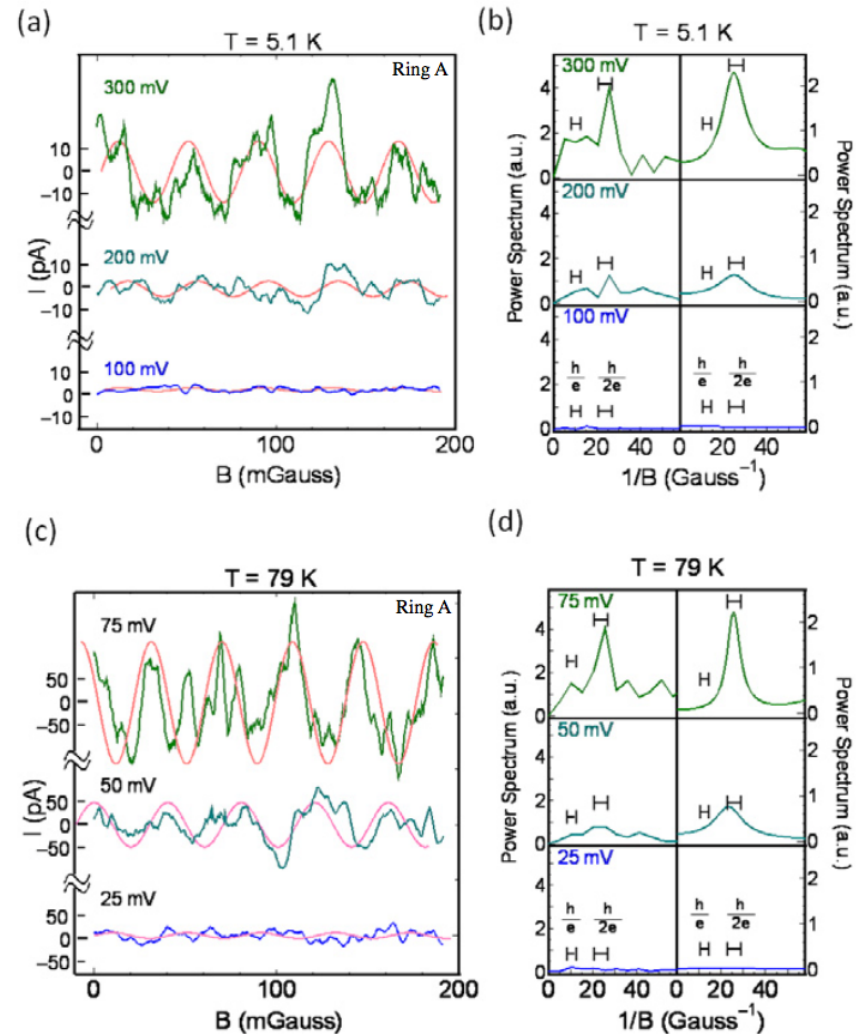


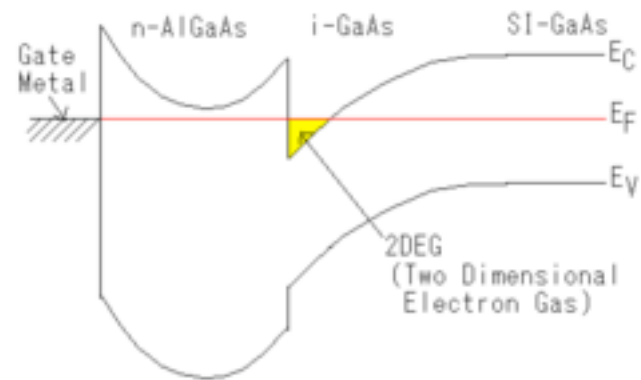
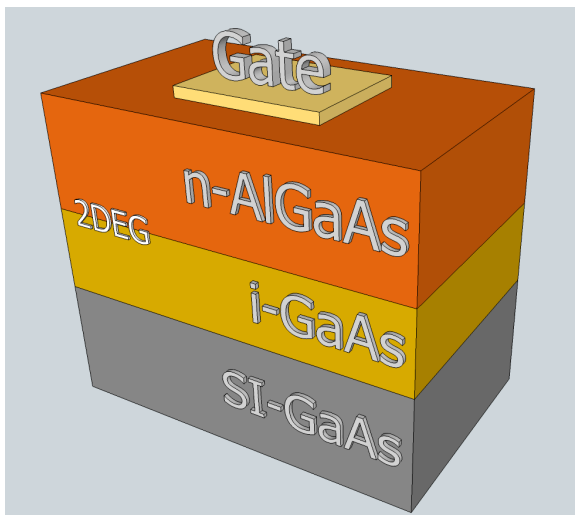
Fig. 2: (Colour on-line) (a) Change in current value as magnetic field at 5.1 K, applied voltage 100, 200 and 300 mV. (b) The power spectra of the observed oscillation using DFT (left) and MEM (right). The red lines show the sinusoidal oscillation estimated by IDFT and MEM. (c) Current depending on the magnetic field and (d) the evolution of the power spectra DFT (left) and MEM (right), 25, 50 and 75 mV, from bottom to top at 79 K.

Aharonov-Bohm effect in charge-density  
wave loops with inherent temporal  
current switching

M. Tsubota, K. Inagaki, T. Matsuura and  
S. Tanda

EPL, 97 (2012) 57011

doi: 10.1209/0295-5075/97/57011



Interferometer-Based Studies of Quantum Hall Phenomena, Douglas T. McClure III (PhD thesis, Harvard, 2012)

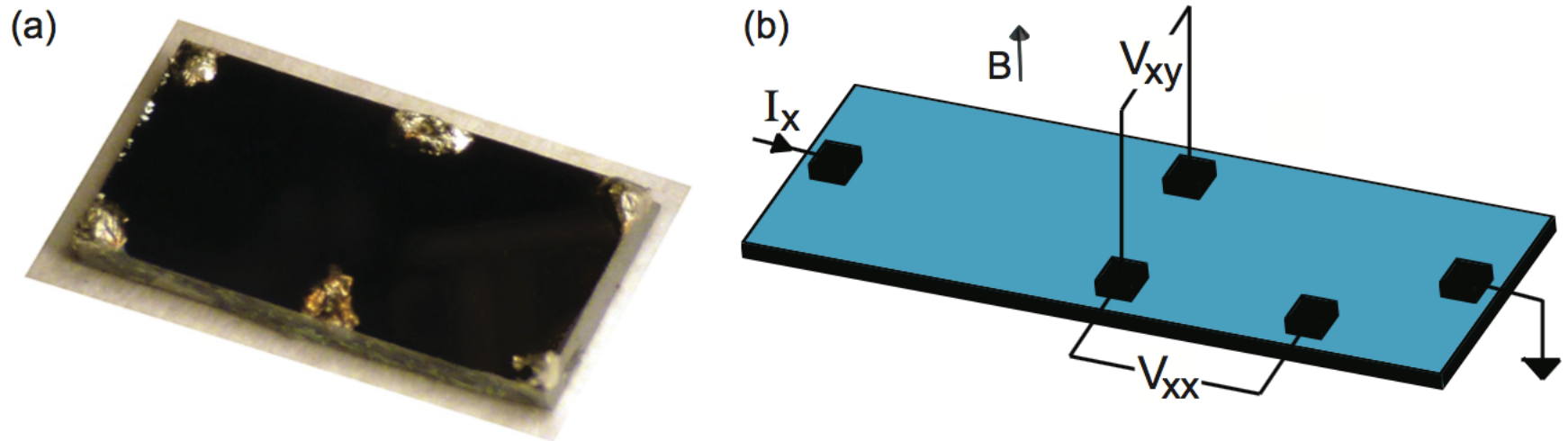


Figure 1.1: (a) GaAs Hall bar with six indium contacts. (b) Schematic Hall bar indicating a typical measurement configuration. A current bias  $I_x$  applied at one end of the sample gives rise to voltages  $V_{xy}$  and  $V_{xx}$ , from which the Hall and longitudinal resistances are calculated as  $R_{xy} = V_{xy}/I_x$  and  $R_{xx} = V_{xx}/I_x$ .



Interferometer-Based Studies of Quantum Hall Phenomena, Douglas T. McClure  
III (PhD thesis, Harvard, 2012)

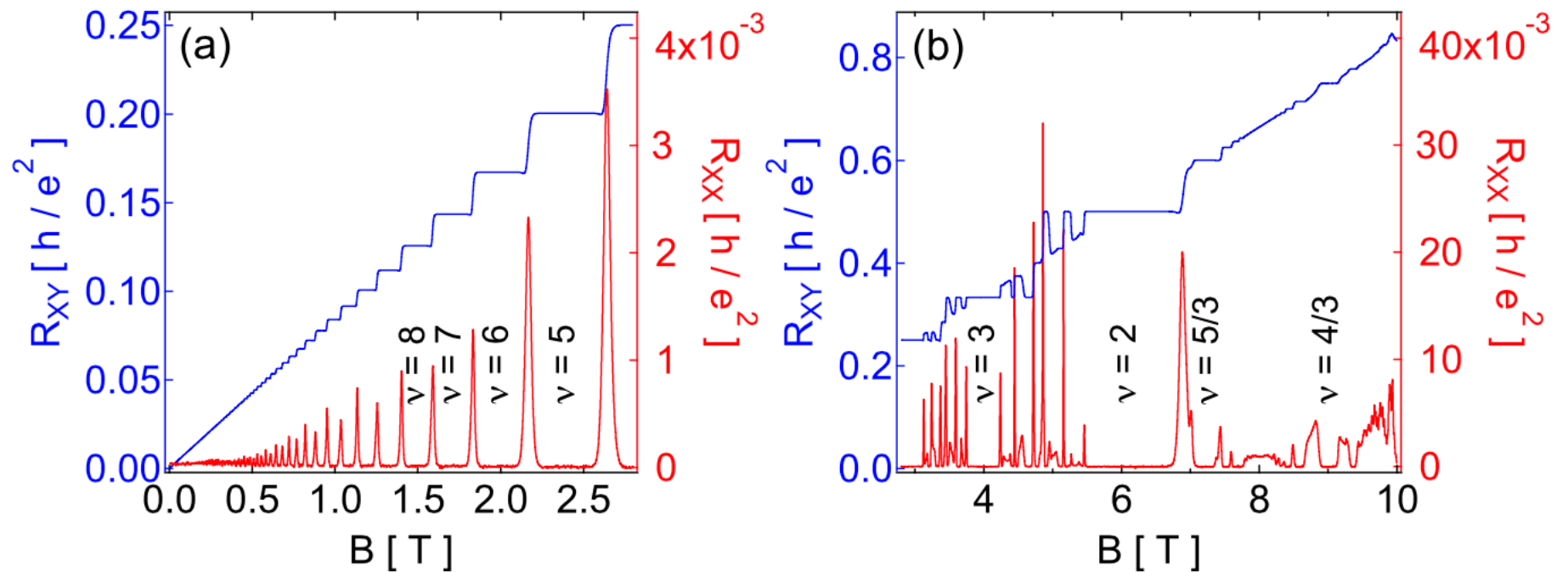


Figure 1.2: Hall (blue) and longitudinal (red) resistances as a function of magnetic field at a temperature of about 0.01 K. (a) As magnetic field is increased, a transition from classical to quantum Hall behavior is observed. (b) At even higher magnetic fields, the Hall resistance exhibits plateaus at fractional filling factors and even non-monotonic features.

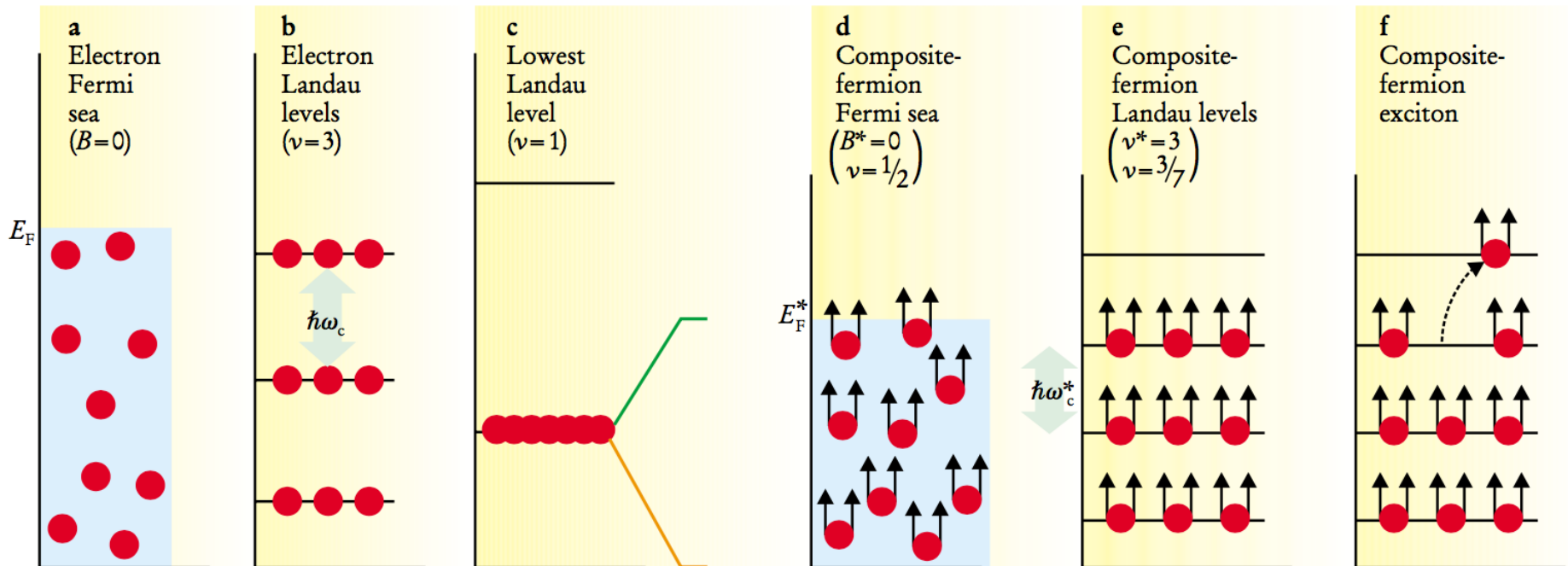
# The Composite Fermion: A Quantum Particle and Its Quantum Fluids

Jainendra K. Jain

Citation: Physics Today **53**(4), 39 (2000); doi: 10.1063/1.883035

## Integer QHE

## Fractional QHE



DENSITY OF STATES

FIGURE 1. EVOLUTION OF A TWO-DIMENSIONAL ELECTRON SYSTEM as the transverse magnetic field  $B$  is increased. For independent electrons, the Fermi sea (a) (filled to Fermi energy  $E_F$ ) at  $B = 0$ , splits into Landau levels (b) separated by the cyclotron energy. The lowest Landau level (c) is split by interactions into energy levels of composite fermions with attached flux quanta, which fill a composite-fermion Fermi sea (d) at  $\nu = 1/2$  and occupy composite-fermion Landau levels (e) at other filling factors. A jump out of such a level (f) creates an exciton, a neutral particle-hole excitation. At still higher fields, this scenario (d–f) repeats itself, but now with composite fermions carrying four or more flux quanta.

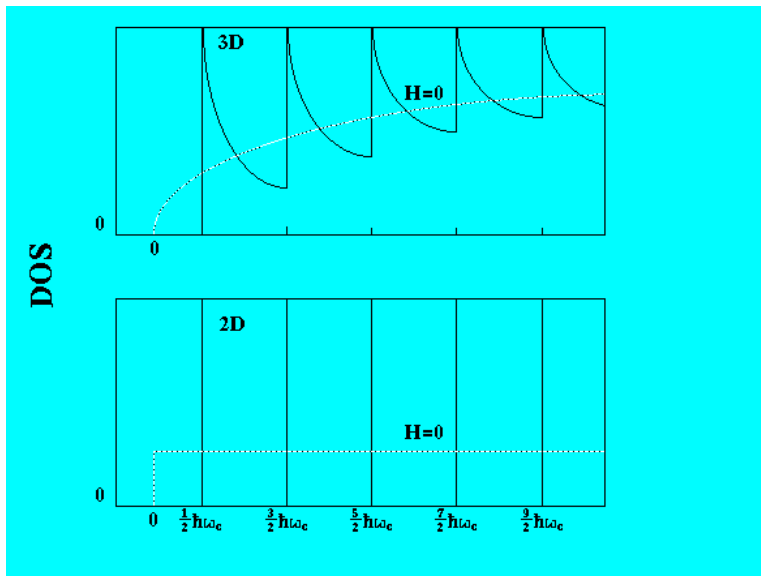
# The Composite Fermion: A Quantum Particle and Its Quantum Fluids

Jainendra K. Jain

Citation: Physics Today **53**(4), 39 (2000); doi: 10.1063/1.883035

Integer QHE

Fractional QHE



Interferometer-Based Studies of Quantum Hall Phenomena, Douglas T. McClure III (PhD thesis, Harvard, 2012)

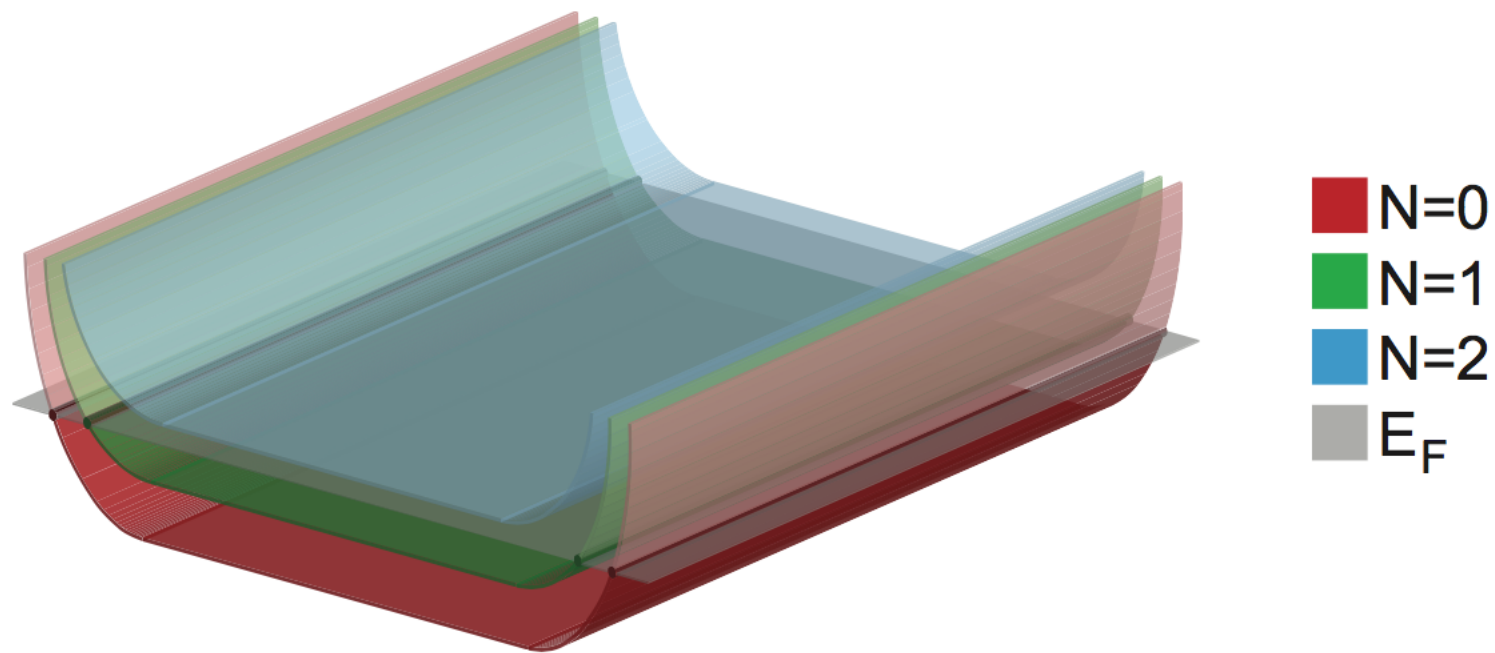


Figure 1.4: Schematic diagram of Landau levels in a 2DEG. In the bulk of the 2DEG, LL's below the Fermi energy are occupied. At the 2DEG edges, LL energies are bent upward by the confining potential, leading to the formation of chiral, current-carrying edge states.

Interferometer-Based Studies of Quantum Hall Phenomena,  
 Douglas T. McClure III (PhD thesis , Harvard, 2012)  
 Chapter 2: Distinct signatures for Coulomb blockade and  
 Aharonov-Bohm interference in electronic Fabry-Perot  
 interferometers

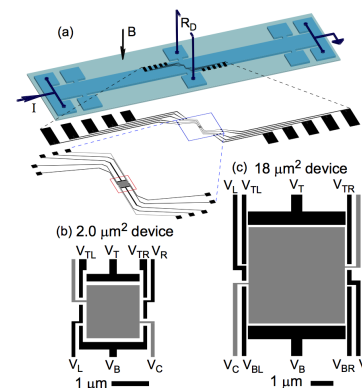


Figure 2.1: Measurement setup and devices. (a) Diagram of the wet-etched Hall bar, surface gates, and measurement configuration. Diagonal resistance,  $R_D$ , is measured directly across the Hall bar, with current bias,  $I$ . Subsequent zoom-ins of the surface gates are also shown; the red box encloses the detailed gate layouts for the device shown in (c). (b,c) Gate layouts for the  $2.0 \mu\text{m}^2$  and  $18 \mu\text{m}^2$  devices, respectively. The areas quoted refer to those under  $V_C$ .

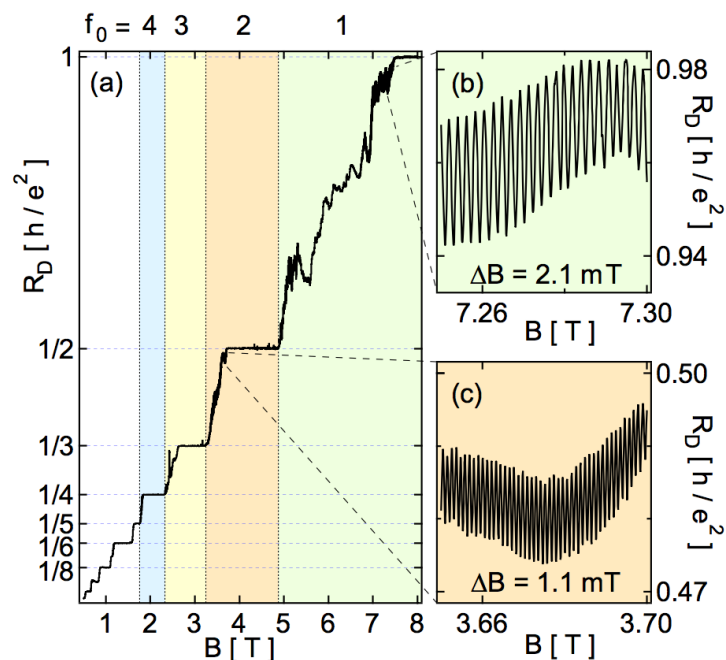


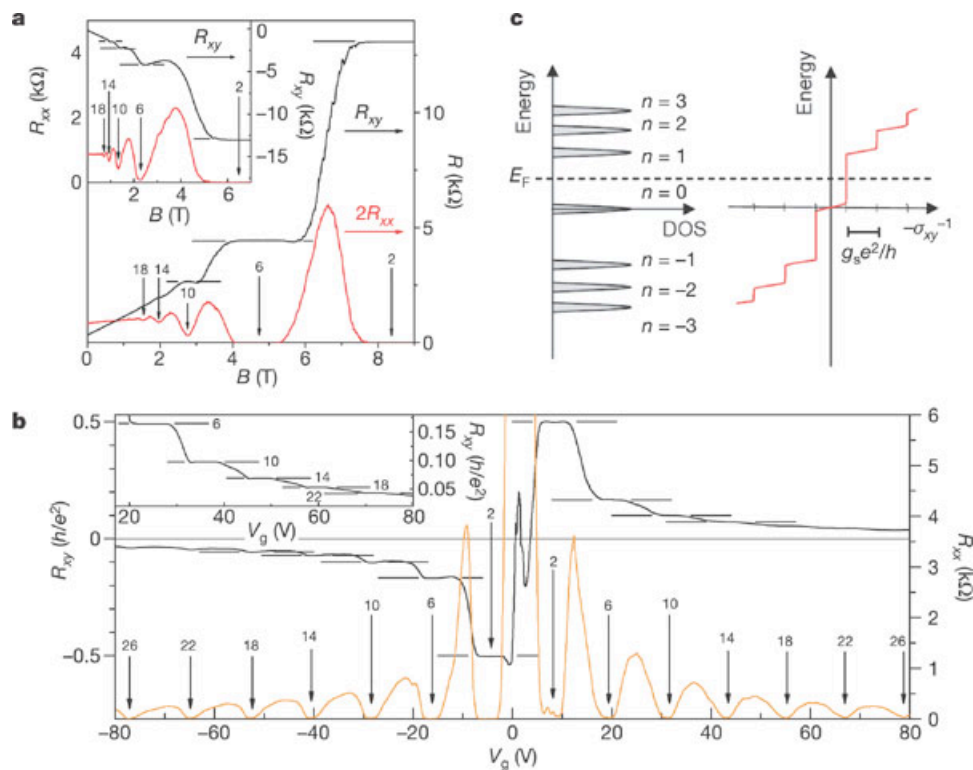
Figure 2.2: Oscillations in  $R_D$  as a function of magnetic field,  $B$ , for the  $2.0 \mu\text{m}^2$  device. (a)  $R_D$  as a function of  $B$ , showing well-quantized integer plateaus. Different colored backgrounds indicate different numbers of fully-occupied LL's,  $f_0$ , through the device. (b, c) Zoom-ins of the data in (a), at  $f_0 = 1$  and  $2$ , respectively, showing oscillations in  $R_D$ , and their  $B$  periods,  $\Delta B$ .

# Experimental observation of the quantum Hall effect and Berry's phase in graphene

Yuanbo Zhang, Yan-Wen Tan, Horst L. Stormer and Philip Kim

Nature 438, 201-204 (10 November 2005)

doi:10.1038/nature04235



a, Hall resistance (black) and magnetoresistance (red) measured in the device in Fig. 1 at  $T = 30$  mK and  $V_g = 15$  V. The vertical arrows and the numbers on them indicate the values of  $B$  and the corresponding filling factor  $\nu$  of the quantum Hall states. The horizontal lines correspond to  $h/e2\nu$  values. The QHE in the electron gas is shown by at least two quantized plateaux in  $R_{xy}$ , with vanishing  $R_{xx}$  in the corresponding magnetic field regime. The inset shows the QHE for a hole gas at  $V_g = -4$  V, measured at 1.6 K. The quantized plateau for filling factor  $\nu = 2$  is well defined, and the second and third plateaux with  $\nu = 6$  and  $\nu = 10$  are also resolved. b, Hall resistance (black) and magnetoresistance (orange) as a function of gate voltage at fixed magnetic field  $B = 9$  T, measured at 1.6 K. The same convention as in a is used here. The upper inset shows a detailed view of high-filling-factor plateaux measured at 30 mK. c, A schematic diagram of the Landau level density of states (DOS) and corresponding quantum Hall conductance ( $\sigma_{xy}$ ) as a function of energy. Note that, in the quantum Hall states, Unfortunately we are unable to provide accessible alternative text for this. If you require assistance to access this image, or to obtain a text description, please contact [npg@nature.com](mailto:npg@nature.com). The LL index  $n$  is shown next to the DOS peak. In our experiment the Fermi energy  $E_F$  can be adjusted by the gate voltage, and  $R_{xy}$  changes by an amount  $g_s e^2/h$  as  $E_F$  crosses a LL.