

PH421: Oscillations – lecture 1

*Reading: Taylor 4.6
(Thornton and Marion 2.6)
Knight 10.7*

Goals for the pendulum module:

- (1) CALCULATE the period of oscillation if we know the potential energy; specific example is the pendulum*
- (2) MEASURE the period of oscillation as a function of oscillation amplitude*
- (3) COMPARE the measured period to models that make different assumptions about the potential*
- (4) PRESENT the data and a discussion of the models in a coherent form consistent with the norms in physics writing*
- (5) CALCULATE the (approximate) motion of a pendulum by solving Newton's $F=ma$ equation*

How do you calculate how long it takes to get from one point to another?

$$\Delta t = \frac{\Delta x}{v}$$

But what if v is not constant?

$$\frac{dx}{dt} = v(x)$$

$$dt = \frac{dx}{v(x)}$$

Separation of x and t variables!

$$\int_{t_{start}}^{t_{finish}} dt = \int_{x(t=t_{start})}^{x(t=t_{finish})} \frac{dx}{v(x)}$$

The case of a conservative force

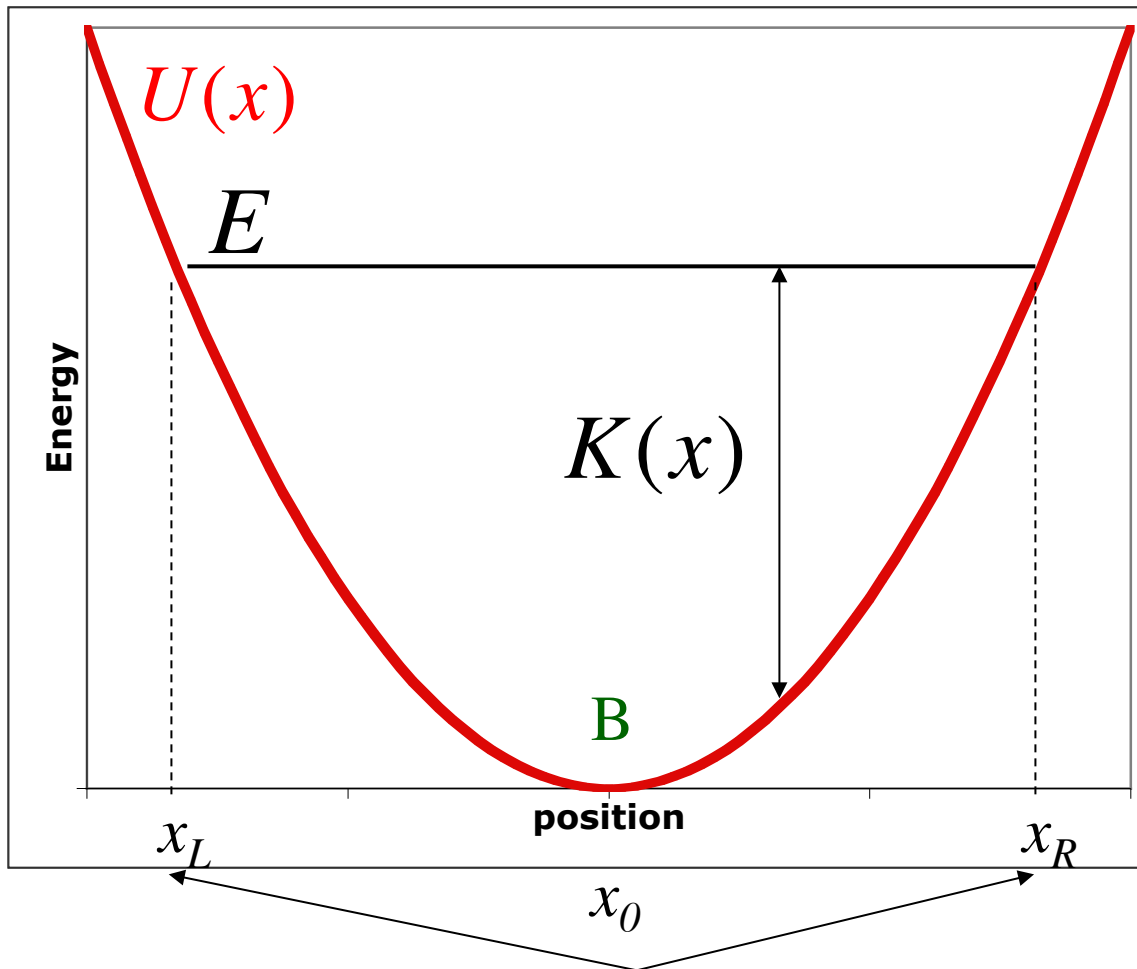
$$E = K(x) + U(x)$$

Suppose total energy is **CONSTANT**
(we have to know it, or be able to
find out what it is)

$$K(x) = \frac{1}{2} m \left[\frac{dx}{dt} \right]^2 \Rightarrow \frac{dx}{dt} = \pm \sqrt{\frac{2}{m} [E - U(x)]}$$

$$dt = \frac{dx}{\pm \sqrt{\frac{2}{m} [E - U(x)]}}$$

Example: $U(x) = \frac{1}{2} kx^2$, the harmonic oscillator



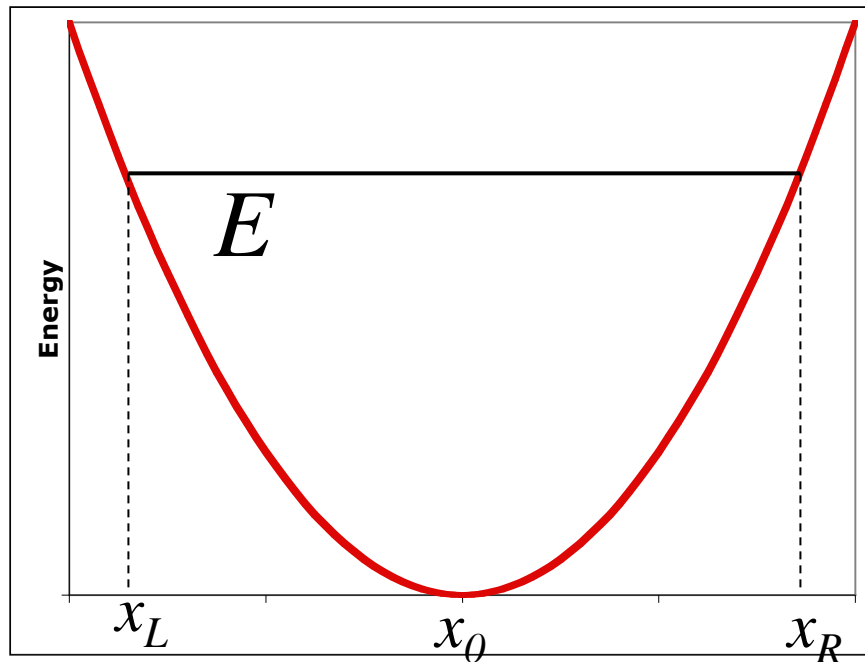
$$U(x) = \frac{1}{2} kx^2 \quad \text{PE}$$

$$F(x) = -\frac{dU}{dx} = -kx$$

$$K(x) = \frac{1}{2} m\dot{x}^2 \quad \text{KE}$$

Classical turning points

$$\dot{x} \equiv \frac{dx}{dt}$$



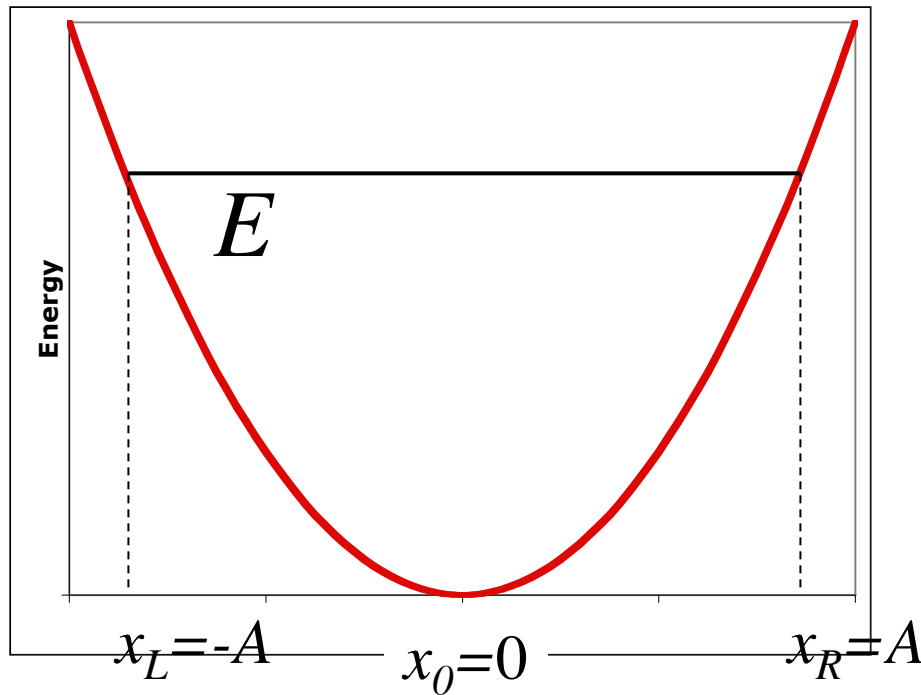
Symmetry - time to go there is the same as time to go back (no damping)

$$T = 2 \int_{x_L}^{x_R} \frac{dx}{\sqrt{\frac{2}{m} [E - U(x)]}}$$

SHO - symmetry about x_0

$x_L \rightarrow x_0$ same time as for $x_0 \rightarrow x_R$

$$T = 4 \int_{x_0}^{x_R} \frac{dx}{\sqrt{\frac{2}{m} [E - \frac{1}{2} kx^2]}}$$



Another way to specify E is via the amplitude A

SHO - do we get what we expect?

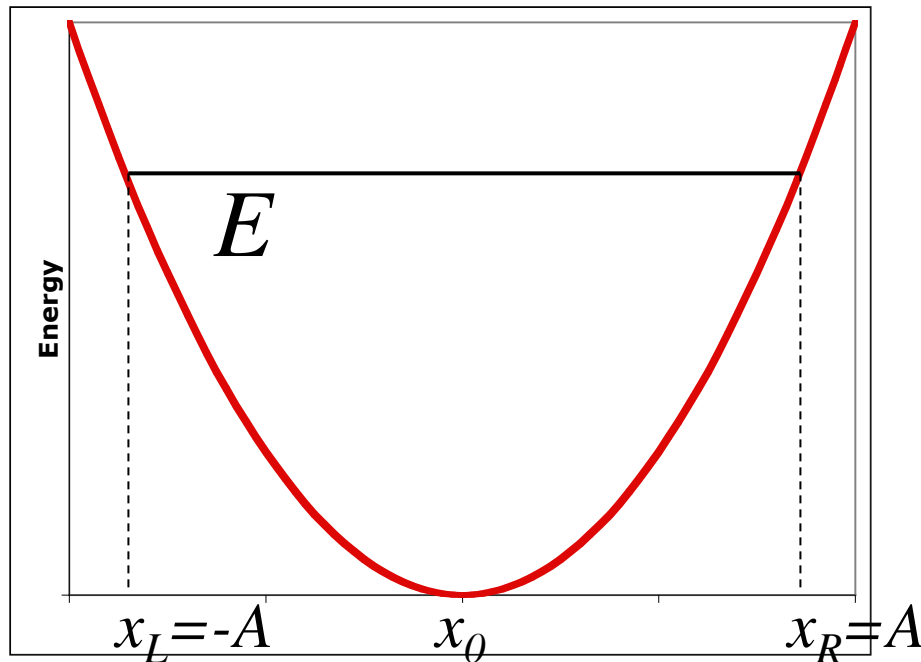
$$T = 4 \int_{x=0}^{x=A} \frac{dx}{\sqrt{\frac{2}{m} \left[E - \frac{1}{2} kx^2 \right]}}$$

$$u = \frac{x}{A} \quad \boxed{E = \frac{1}{2} kA^2}$$

$$\boxed{T = 4 \int_{u=0}^{u=1} \frac{du}{\sqrt{\frac{k}{m} [1 - u^2]}}}$$

Independent of A !

<http://www.wellesley.edu/Physics/Yhu/Animations/sho.html>



$$T = 4 \int_{u=0}^{u=1} \frac{du}{\sqrt{\frac{k}{m} [1 - u^2]}}$$

$$u = \sin \phi$$

$$du = \cos \phi d\phi$$

$$\sqrt{1 - u^2} = \cos \phi$$

↓

$$T = 4 \int_{\phi=0}^{\phi=\pi/2} \frac{\cos \phi d\phi}{\sqrt{\frac{k}{m} \cos \phi}}$$

$$T_{SHO} = 2\pi \sqrt{\frac{m}{k}} \equiv \frac{2\pi}{\omega_0}$$

You have seen this before in intro PH, but you didn't derive it this way.

$$T_{SHO} = \frac{2\pi}{\omega_0}$$

Period of SHO is INDEPENDENT OF AMPLITUDE
Why is this surprising or interesting?

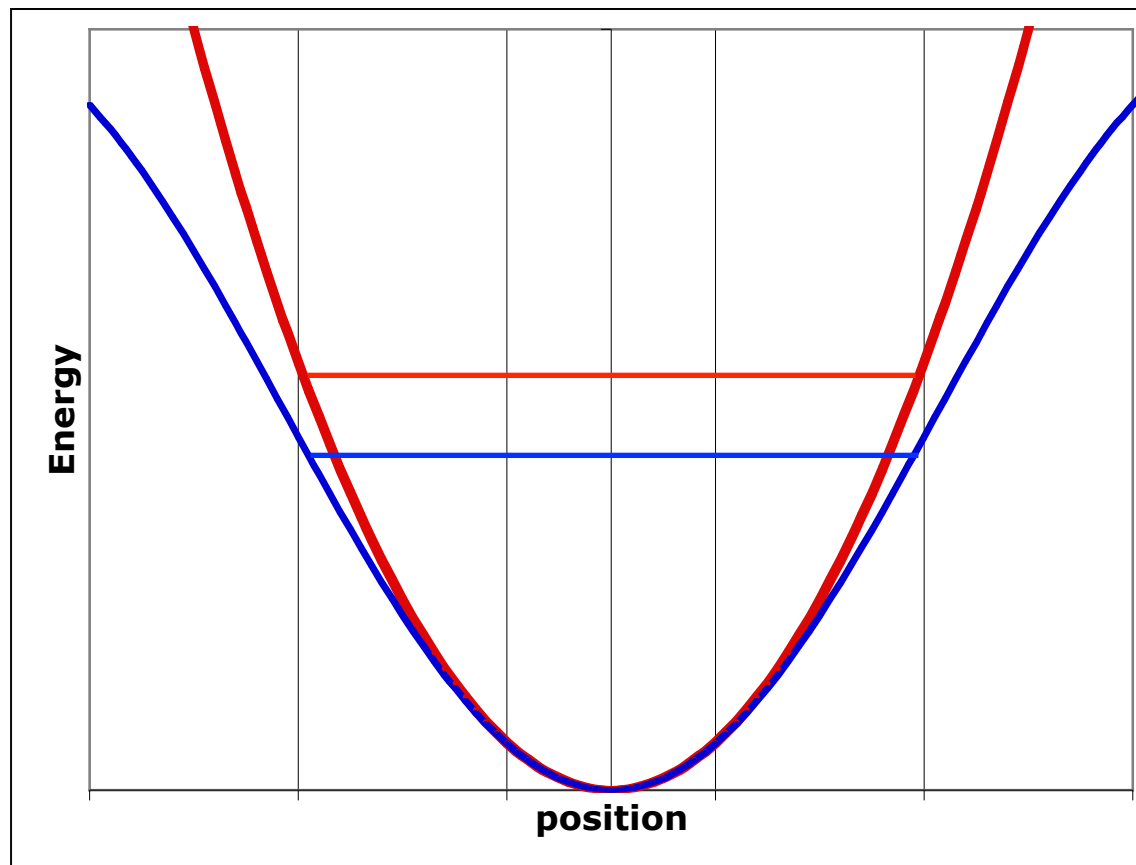
As A increases, the distance and velocity change. How does this affect the period for ANY potential?

A increases \rightarrow further to travel \rightarrow distance increases
 \rightarrow period increases

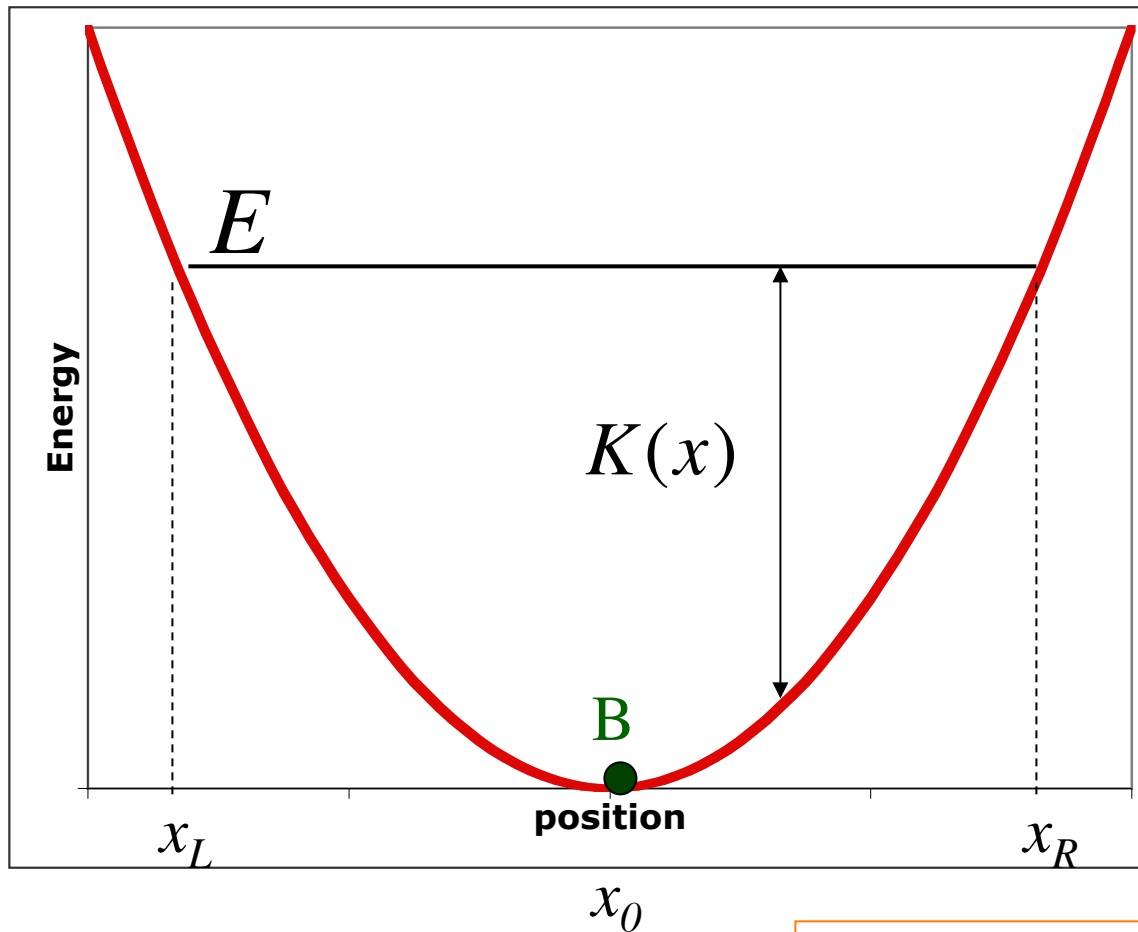
A increases \rightarrow more energy \rightarrow velocity increases
 \rightarrow period decreases

Which one wins, or is it a tie?

Compare $1 - \cos \theta$ with $\frac{1}{2!} \theta^2$



Period increases because $v(x)$ is smaller at every x (*why?*) in the trajectory. Effect is magnified for larger amplitudes.



$$\text{Eqm: } \left. \frac{dU}{dx} \right|_{x_0} = 0$$

$$\text{If } \left. \frac{d^2U}{dx^2} \right|_{x_0} = k$$

“Everything is a SHO!”

$$U(x) = U(x_0) + \frac{1}{2} k(x - x_0)^2$$

$$U(x) = U(x_0) + (x - x_0) \left. \frac{dU}{dx} \right|_{x_0} + \frac{(x - x_0)^2}{2!} \left. \frac{d^2U}{dx^2} \right|_{x_0} + \frac{(x - x_0)^3}{3!} \left. \frac{d^3U}{dx^3} \right|_{x_0}$$