## 1. A. PRACTICE

Read Main and Taylor on the topic of damped and driven oscillators. You should check out this website (which I'll also show in class). Be warned that at first sight, it is very busy and confusing! But after you study it for a while, you'll realize it has a great deal of good information about the circuit we'll study.

http://www.ngsir.netfirms.com/englishhtm/RLC.htm

Taylor 5.29, Main 3.1 – 3.8

## B. REQUIRED

- 1. Main, problem 3.5
- 2. The current in a circuit is represented by the complex number  $I(t) = I_0 e^{i(\omega t + \phi)}$ . Current cannot be complex, but we represent it this way with the understanding that the real part will represent the actual current in the circuit.

(a) Represent the current on an Argand diagram at t = 0 and with your choice of  $\phi$ .

(b) Represent  $\dot{I}$  on the same diagram at t = 0 and with the same choice of  $\phi$ . What is the phase of *I* relative to  $\dot{I}(t)$  at any time?

(c) Represent the charge q flowing in that circuit on the same diagram at t = 0 and the same choice of  $\phi$ . What is the phase of I relative to q(t) at any time?

3. A series LRC circuit is driven by a sinusoidal voltage that by convention we write as  $V_0 cos(\omega t)$ . Draw phasor diagrams (*i.e.* on an Argand plot) representing the driving voltage and each of the voltages across the capacitor  $V_c$ , resistor  $V_R$ , inductor  $V_L$  in a driven LRC circuit for three different cases: (1)  $\omega \ll \omega_0$ , (2)  $\omega = \omega_0$  (resonance frequency), (3)  $\omega \gg \omega_0$ . An example is done for you. For each of the 3 frequency regimes, explain whether and why the circuit as a whole behaves predominantly as a resistor, capacitor, or inductor.

	$V_R$		V <sub>L</sub>		V <sub>C</sub>	
$\omega << \omega_0$		Im		Im		Im
		Re		Re		R
$\omega \approx \omega_0$		Im		Im		Im
		Re		Re		R
		I		I		I



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$\omega >> \omega_0$	Im	Im	Im
	Re	Re	R

## 4. A forced oscillator characterized by a damping $\beta$ and frequency $\omega_0$ has a generalized displacement response $x(t) = \frac{F_0 / m}{\sqrt{(\omega_0^2 - \omega^2)^2 + 4\omega^2 \beta^2}} \cos(\omega t + \delta)$ .

This function is of the form  $x(t) = A(\omega)\cos(\omega t + \delta)$ , where  $A(\omega)$  measures the amplitude of the response and it is a (strong) function of frequency.

- (a) Sketch the form of  $A(\omega)$ .
- (b) Find the exact frequency at which the amplitude  $A(\omega)$  is maximal.

(c) If the damping is light ( $\omega_0\beta << 1$ ), show that the frequency of maximum response (resonance frequency) is approximately  $\omega_0$ , and that it differs from  $\omega_0$  by a term of order

 $\beta^2$ . You will need to use the approximation  $(1+z)^p \approx 1 + pz + \frac{p(p-1)}{2!}z^2 + ...$ 

In what follows, you will be looking for values of  $\omega$  that are very close to  $\omega_0$ , *i.e.* 

 $\omega \approx \omega_0 (1 + something small)$ . Think hard about what level of approximation is appropriate (you will learn much more about this in PH320).

(d) In the light damping approximation, what is the value of the maximum value of  $A(\omega)$ ? (e) In the light damping approximation, find the "width" of the  $A(\omega)$  function, by finding the two frequencies at which  $A(\omega)$  drops to  $\frac{1}{\sqrt{2}}$  of its maximum.

(f) Define the difference between these two frequencies as  $\Delta \omega$ , and find an expression for

 $Q \equiv \frac{\omega_0}{\Delta \omega}$  in terms of  $\beta$ .

[Note: Generally, it is the *power* that is measured (proportional to the square of the displacement), not the displacement itself. Thus  $\Delta \omega$  as defined above corresponds to the full-width at half-maximum (FWHM) of the *power* response curve.]

