

GREEN'S FUNCTIONS, APPLIED TO NON HOMOGENEOUS ELLIPTIC PROBLEMS

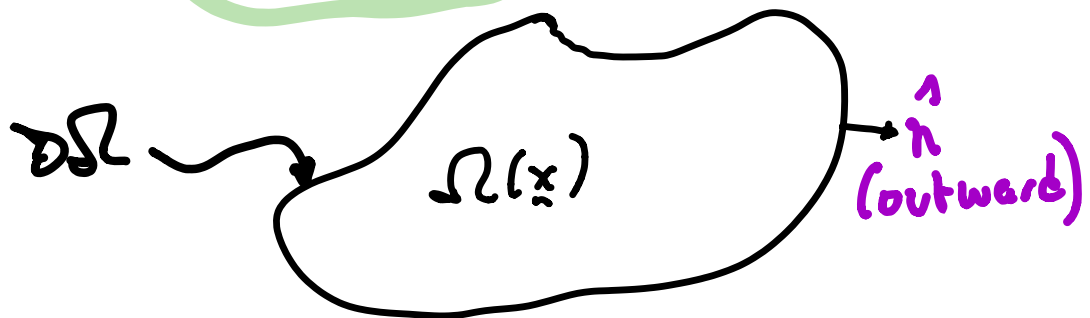
$$\textcircled{A} \quad \mathcal{L}u \equiv -\nabla \cdot (p \nabla u) + qu = \rho F \text{ in } \Omega$$

$$\left(\alpha u + \beta \frac{\partial u}{\partial n} \right) \Big|_{\partial \Omega} = B$$

makes this
problem
non
homogeneous

where

$$\frac{\partial u}{\partial n} = \hat{n} \cdot \nabla u$$



$\rho F(x)$ is known (inhomogeneous term in the PDE), so is B .

$$\text{Form } \mathcal{L}v = -\nabla \cdot (p \nabla v) + qv = \rho F \quad \textcircled{B}$$

Multiply (A) by v and (B) by u
subtract & integrate:

$$\int_{\Omega} (v \Delta u - u \Delta v) d\vec{x} \\ = - \int_{\Omega} \nabla \cdot \underbrace{(p v \nabla u - p u \nabla v)}_{\vec{Q}} d\vec{x} = 0$$

The RHS of (A) & (B) are the same, so zero

Use Divergence Theorem, i.e.

$$\oint_{\Omega} \nabla \cdot \vec{Q} d\vec{x} = \int_{\partial\Omega} \vec{Q} \cdot \hat{n} dS$$

$$\text{here } \vec{Q} = p v \nabla u - p u \nabla v$$

So

$$= - \int_{\Omega} \nabla \cdot (p v \nabla u - p u \nabla v) d\vec{x}$$

$$= - \int_{\partial\Omega} p (v \nabla u - u \nabla v) \cdot \hat{n} dS$$

since $\hat{n} \cdot \nabla w \equiv \frac{\partial w}{\partial n}$, then

$$= \int_{\partial\Omega} p \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) dS$$

we have data for $u, \frac{\partial u}{\partial n}, v, \frac{\partial v}{\partial n}$ at

$\partial\Omega$.

$$\therefore \int_{\Omega} (v \Delta u - u \Delta v) d\vec{x} = \int_{\partial\Omega} p \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) dS \quad (*)$$

Now, suppose v is the solution to

$$\textcircled{D} \quad \Delta v = \delta(\underline{x} - \underline{\xi})$$

$\delta(\underline{x} - \underline{\xi})$ is the Dirac Delta function,
centered at $\underline{\xi}$ inside Ω .

Recall that

$$\left\{ \begin{array}{l} \int_{\Omega} \delta(\underline{x} - \underline{\xi}) d\underline{x} = 1 \\ \int_{\Omega} h(\underline{x}) \delta(\underline{x} - \underline{\xi}) d\underline{x} = h(\underline{\xi}) \end{array} \right.$$

"sifting property"

Use the sifting property:

$$\textcircled{F} \quad \int_{\Omega} u \Delta v \, d\underline{x} = \int_{\Omega} u \delta(\underline{x} - \underline{\xi}) \, d\underline{x} = u(\underline{\xi})$$

using \textcircled{D} .

Also,

$$\textcircled{G} \quad \int_{\Omega} v \Delta u \, d\underline{x} = \int_{\Omega} v \rho F \, d\underline{x},$$

since $\Delta u = \rho F$, see \textcircled{B} .

Substitute \textcircled{F} & \textcircled{G} into \textcircled{I} , into

$$\int_{\Omega} (v \Delta u - u \Delta v) \, d\underline{x} = \int_{\partial\Omega} \underbrace{\left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right)}_{\textcircled{I}} \, dS$$

to get

$$\int_{\Omega} v \rho F d\underline{x} - u(\underline{\xi}) = \int_{\partial\Omega} p \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) dS.$$

Solving for $u(\underline{\xi})$, $\underline{\xi} \in \Omega$,

$$u(\underline{\xi}) = \int_{\Omega} v \rho F d\underline{x} - \int_{\partial\Omega} p \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) dS. \quad (\star)$$

here v is the solution to $\Delta v = \delta(\underline{x} - \underline{\xi})$

so, in principle we know it.

We will call v the Green's function.

We know that $\Delta v = \delta(\underline{x} - \underline{\xi})$

But we need B.C.:

use similar boundary conditions as posed by the BVP, i.e. we'll impose

$$\alpha v + \beta \frac{\partial v}{\partial n} \Big|_{\partial \Omega} = 0$$

i.e. the homogeneous version of the original B.C.

\therefore The boundary terms, i.e.
the RHS of ~~*~~

$$\begin{aligned} & - \int_{\partial \Omega} p \left(u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) dS \\ & = \int -p \frac{1}{2} \left[\alpha u \frac{\partial v}{\partial n} - \alpha v \frac{\partial u}{\partial n} \right] dS \end{aligned}$$

$\left(-\beta \frac{\partial v}{\partial n} \right)$

factor $\partial v / \partial n$

$$\begin{aligned} &= - \int \Gamma \frac{1}{\alpha} \frac{\partial v}{\partial n} \left(\alpha u + \beta \frac{\partial u}{\partial n} \right) dS \\ &= - \int \Gamma \frac{1}{\alpha} \beta \frac{\partial v}{\partial n} dS = \int \Gamma \frac{1}{\beta} \beta v dS \end{aligned}$$

In Summary, using $(*)$:

$$u(\xi) = \int_{\Omega} v \rho F dx + \int_{\partial \Omega} \Gamma \frac{1}{\beta} \beta v dS$$

So $v = G(x, \xi)$ is the Green's Function

it satisfies

$$\Delta G = \delta(x - \xi)$$

$$\left(\alpha G + \beta \frac{\partial G}{\partial n} \right) \Big|_{\partial \Omega} = 0$$

it solves the problem

$$\begin{cases} \Delta u = \rho F & \text{in } \Omega \\ (\alpha u + \beta \frac{\partial u}{\partial n})|_{\partial\Omega} = B \end{cases}$$

as

$$(2) \quad u(x) = \int \rho G F dx - \int \frac{\alpha B}{\alpha} \frac{\partial G}{\partial n} ds + \int \frac{\rho B}{\beta} G ds //$$

Application: in the solution of

the forced Helmholtz Eq

$$(\nabla^2 + k^2)u = \rho F$$

with B.C. as above.

In that case (2) is known as the

Helmholtz Integral Solution

This is a very popular formulation in time harmonic problems in acoustics, electromagnetics, etc.

So let's summarize:
GENERIC RECIPE FOR

NON-HOMOGENEOUS LINEAR ODEs:

① $Lu(x) = f(x) \quad x \in \Omega = (a, b) \subseteq \mathbb{R}$

② $B_k u(x) = h_k(x) \quad x \in \partial\Omega = \{a, b\}$
 $k = 1, 2, \dots, n-1$

L is a linear n^{th} -order differential operator.

③ are the $n-1$ required boundary conditions involving linear combinations of u and its derivatives, up to order $n-1$, evaluated on the boundary.

We form:

$$(G, Lu) \equiv \int_a^b G(x, y) L u(y) dy$$

Integrate by parts,

$$(\star) (G, Lu) = (L^* G, u) + \text{integration boundary terms}$$

L^* is the adjoint to L

Remark: if $L = L^*$ then we say that L is self-adjoint.

$$\text{If } L^* G(x, y) = \delta(x - y) \\ + \text{terms involving b.c.}$$

then (\star) yields

$$u = -(\text{integration terms}) + \int_a^b G(x, y) f(y) dy //$$

GENERAL RECIPE FOR NONHOMOGENEOUS LINEAR PDE

$$\text{PDE} \quad Lu(P) = f(P) \quad P \in \Omega \subseteq \mathbb{R}^n$$

$$\text{B.C.} \quad Bu_k(p) = h(p) \quad p \in \partial\Omega$$

$$\text{here } P = (x_1, x_2, \dots, x_n) \in \Omega \subseteq \mathbb{R}^n$$

and p is a point on the boundary $\partial\Omega$

let $Q = (x_1, x_2, \dots, x_n)$ another point in Ω .

We have

$$(G, Lu) = \int_{\Omega} G(P, Q) Lu(P) dP$$

we compute

$$(G, Lu) = (L^*G, u) + \text{integrated terms,}$$

by integrating by parts. Then, solve

$$\begin{cases} L^* G(P, Q) = \delta(P, Q), \text{ the Dirac delta function} \\ \text{with the homogeneous B.C.} \end{cases}$$

then

$$u(Q) = -\text{integrated terms} + \int G(P, Q) f(P) dP$$

ADJOINT OPERATORS

They came up in Differential Equations, i.e.

$$\text{in } (v, Lu) = (L^* v, u)$$

+ boundary terms.

but they came up in a variety of other contexts.

ex) Find L^* , if

$$L = \frac{d}{dx} \quad \text{and boundary conditions}$$

$$\text{or } 0 \leq x \leq 1 \quad v(0) - 2v(1) = 0.$$

Integrate by parts...

$$\begin{aligned}
 (u, Lv) &= \int_0^1 u \frac{\partial v}{\partial x} dx \\
 &= [uv] \Big|_0^1 - \int_0^1 v \frac{du}{dx} dx \\
 &= u(1)v(1) - u(0)v(0) - \int_0^1 v \frac{du}{dx} dx
 \end{aligned}$$

Apply $v(0) = 2v(1)$:

$$\begin{aligned}
 &= u(1)v(1) - u(0)2v(1) - \int_0^1 v \frac{du}{dx} dx \\
 &= v(1)[u(1) - 2u(0)] - \int_0^1 v \frac{du}{dx} dx \\
 &= (L^* u, v)
 \end{aligned}$$

$\therefore L^*$ is $-\frac{d}{dx}$ with $u(1) = 2u(0)$
for boundary conditions.

Remark: $L^* \neq L$, i.e. L is not self adjoint. //

ex) Find L^* for

$$L = e^x \frac{d^2}{dx^2} + e^x \frac{d}{dx} \quad 0 \leq x \leq 1$$

and boundary conditions $u'(0) = 0$ and $u(1) = 0$.

$$(v, Lu) = \int_0^1 v \left[e^x \frac{d^2}{dx^2} + e^x \frac{d}{dx} \right] u dx$$

Integrate by parts: to do this expediently

we rewrite

$$(v, Lu) = \int_0^1 v (e^x u')' dx$$

$$= [v e^x u]_0^1 - [v' e^x u]_0^1 - \int_0^1 u (e^x v')' dx$$

$$= u'(1)v(1)e^1 + v'(0)u(0) - \int_0^1 u [e^x v'' + e^x v'] dx$$

$$= (L^* v, u)$$

with $L^* = L$ (self adjoint)

$$\text{and } v'(0) = 0, v(1) = 0$$

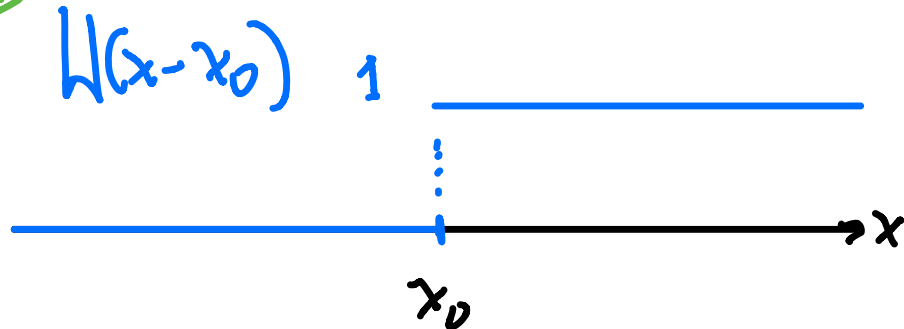
Let's now examine in more detail
the following ODE problem:

Solve

$$(*) \quad \frac{d^2 K}{dx^2} = \delta(x - \xi)$$

We require 2 boundary conditions, more later on this.

Introduce the Heaviside Function:



$$H(x-x_0) = \begin{cases} 0 & \text{for } x < x_0 \\ 1 & \text{for } x \geq x_0 \end{cases}$$

The $H'(x) = \delta(x)$

Integrate (†)

$$\frac{dK}{dx}(x, \xi) = H(x-\xi) + \alpha(\xi)$$

arbitrary

Integrate again

$$K(x, \xi) = \int_{-\infty}^{\infty} H(x-\xi) dx + x\alpha(\xi) + \beta(\xi)$$

arbitrary

$$\therefore (\ddagger) \quad K(x, \xi) = (x-\xi)H(x-\xi) + x\alpha(\xi) + \beta(\xi)$$

K is a continuous, differentiable function.

Let's solve Poisson's Equation in 1D:

$$(y) \left\{ \begin{array}{l} \frac{d^2 u}{dx^2} = f(x) \quad 0 \leq x \leq 1 \\ u(0) = a \quad u(1) = b \end{array} \right.$$

We will use (f) . We need 2
boundary conditions to find α, β .

A side Calculation:

Use sifting property of Dirac Delta:

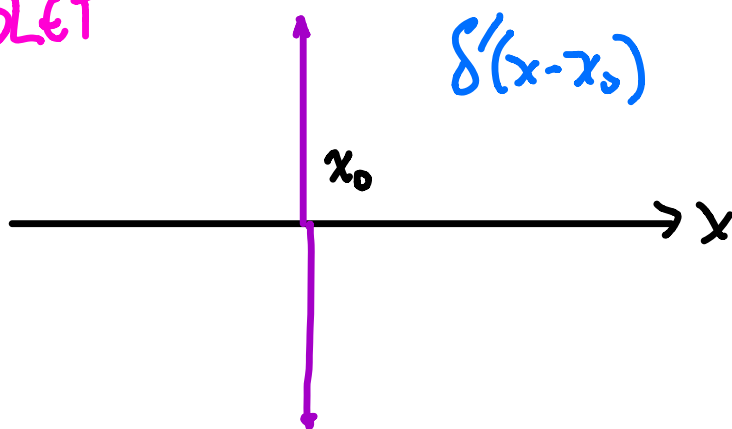
$$g'(0) = \int_{-\infty}^{\infty} g'(x) \delta(x) dx = - \int_{-\infty}^{\infty} g(x) \delta'(x) dx$$

by integrating by parts.

Also,

$$g''(0) = - \int_{-\infty}^{\infty} \delta'(x) g'(x) dx = \int \delta''(x) g(x) dx$$

DOUBLET



// Another Side Calculation

Suppose $Lu = f, (\neq)$

a linear differential equation. It could be a PDE or ODE, but let's assume an ODE and $u(x)$.

Formally, define L^{-1} , is an operator, s.t.

$L^{-1}L = I$, the identity. Hence
 $L^{-1}Lu = u(x)$.

Using (*):

$$u(x) = L^{-1}Lu = L^{-1}f$$

and, from previous discussion

$$L^{-1}u(x) = \int K(x, \xi) u(\xi) d\xi$$

$$\therefore u(x) = \int K(x, \xi) u(\xi) d\xi$$

Since L operates on x , we will put it into the integral:

$$u(x) = \int L K(x, \xi) u(\xi) d\xi$$

write as

$$LK(x, \xi) \equiv g(x, \xi)$$

then

$$u(x) = \int g(x, \xi) u(\xi) d\xi.$$

If true for all $u(x)$ continuous, then

$$g(x, \xi) = 0 \text{ when } x \neq \xi \text{ and}$$

$$\text{when } x = \xi \quad \int g(x, \xi) u(\xi) d\xi = u(x)$$

Back to
$$\begin{cases} u''(x) = f(x) & 0 < x < 1 \\ u(0) = a \quad u(1) = b \end{cases}$$

formally, write (see previous calculation (~~X~~))

$$Ku = -f(x) + b \delta'(x-1) - a \delta'(x) \quad (\star)$$

$$\text{i.e. } K = -\frac{d^2}{dx^2} \quad \text{s.t.}$$

$$Kg = -\delta(x-\xi)$$

it was found that:

$$\therefore (\dagger) \quad K(x, \xi) = (x - \xi) H(x - \xi) + x \alpha(\xi) + \beta(\xi) //$$

Propose a solution to (f) of the form

$$u(x) = u_1(x) + u_2(x).$$

Focus on u_1 : compare (A) and (†) to suggest

$$\text{let } g(x, \xi) \equiv \underline{(x - \xi) H(x - \xi)} - \alpha(1 - \xi) \quad (\S)$$

for $0 < x, \xi < 1$,

where α is to be determined.

Multiply (‡) by $f(\xi)$ and integrate

$$u_1(x) = \int_0^1 f(\xi) g(x, \xi) d\xi$$

$$(\S) \quad u_1(x) = \int_0^1 \overset{x}{\leftarrow \text{due to } H(x-\xi) \text{ in } g} (x - t) f(t) dt - x \alpha \int_0^1 (1 - t) f(t) dt.$$

We can confirm that u_1 obeys $Ku_1 = -f$:

Multiply u_1 by K : $Ku_1 = K \int_0^x (x-t)f(t)dt - Kx \int_0^1 (1-t)f(t)dt$

$$\begin{aligned}\therefore Ku_1 &= K \int_0^1 f(t)g(x,t)dt \\ &= \int_0^1 f(t)Kg(x,t)dt = - \int_0^1 f(t)\delta(x-t)dt \\ &= -f(x).\end{aligned}$$

To complete the solution: we require
 u_2 s.t.

$$Ku_2 = b\delta'(x-1) - a\delta'(x)$$

The idea is to write the δ' above in terms of K : We know that $-\delta(x-t) = Kg$

$$\textcircled{i} \quad -\delta(x-t) = Kg$$

and that

$$\textcircled{3} \quad g(x, t) = (x-t) H(x-t) - x(1-t)$$

Take

$$\frac{d}{dt} \textcircled{A}: \quad \delta'(x-t) = K \frac{dg}{dt}$$

$$\text{where } \frac{dg}{dt} = -H(x-t) + x,$$

obtained by $\frac{d}{dt} \textcircled{B}$.

We can thus infer that

$$u_2 = -b \left. \frac{dg}{dt} \right|_{t=1} - a \left. \frac{dg}{dt} \right|_{t=0}$$

$$\text{where } Ku_2 = b \delta'(x-1) - a \delta'(x-1).$$

$$\text{Since } \left. \frac{dg}{dt} \right|_{t=1} = -H(x-1) + x = x \quad \forall x \in [0, 1]$$

" 0 for $0 \leq x < 1$

$$\left. \frac{dy}{dt} \right|_{t=0} = -h(x) + x = x - 1 \quad \forall x \in [0,1]$$

-1" for $0 < x \leq 1$

$$\therefore u_2(x) = bx - a(x-1).$$

finally $u = u_1 + u_2$

$$= \int_0^x (x-t) f(t) dt - x \int_0^1 (1-t) f(t) dt$$

$$+ bx - a(x-1)$$

Rule: How did we know to construct a solution as a linear superposition of two subproblems? experience... it's a trick, but in linear problems you should always assume you can write

the solution as a linear superposition
of subproblems.

