

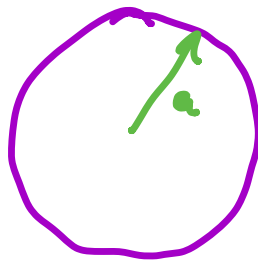
Solve $\Delta u = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2}$

PDE

$$\begin{aligned} 0 < r < a \\ 0 < \theta < 2\pi \\ t > 0 \end{aligned}$$

B.C. $\begin{cases} u(a, \theta, t) = 0 \\ u(0, \theta, t) \text{ bounded} \end{cases}$

$$\begin{cases} u(r, \theta, t) = u(r, \theta + 2\pi, t) & n \in \mathbb{Z} \\ \text{(periodic)} \end{cases}$$



I.C. $u(r, \theta, 0) = f(r), u_t(r, \theta, 0) = g(r)$

$$\Delta u = \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial r^2}, \text{ since } u = u(r, t)$$

$$\text{let } u(r, t) = R(r)W(t)$$

$$\frac{d^2 R}{dr^2} + \frac{1}{r} \frac{\partial R}{\partial r} + k^2 R = 0$$

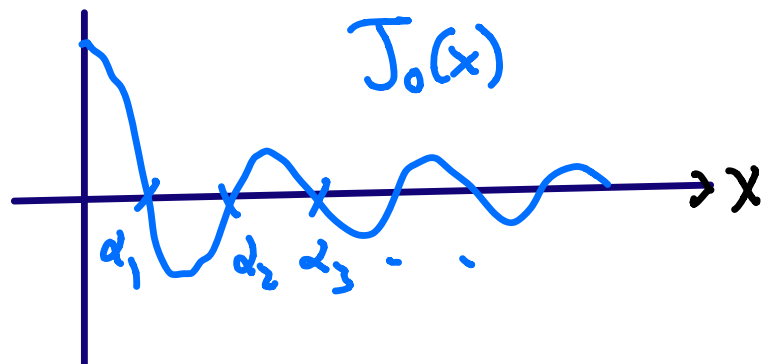
$$R(r) = A J_0(kr) + B Y_0(kr)$$

Since $u(r, \theta, t)$ is bounded everywhere, including at $r=0$, then $B=0$

$$R(r) = A J_0(kr)$$

at $r=a$

$$R(a) = A J_0(ka) = 0$$



$$k_n = \alpha_n / a \quad n=1, 2, \dots$$

The PDE; $\Delta RW = \frac{1}{c^2} RW''$

$$\text{So } W'' + \omega_n^2 W = 0 \quad (*)$$

$$\text{where } \omega_n = c^2 k_n^2$$

The solution to (*) is in terms of $\{\cos \omega_n t, \sin \omega_n t\}$. Let's use the complex representation:

$$W_n = D_n e^{i \omega_n t},$$

D_n is complex.

$$u(r, \theta, t) = \sum_{n=1}^{\infty} D_n J_0(\lambda_n r) e^{i\omega_n t} \quad (\dagger)$$

$$D_n \in \mathbb{C}$$

$$u(r, \theta, 0) = \sum_{n=1}^{\infty} D_n J_0(\lambda_n r)$$

Here we use

$$\int_0^a r J_0(\alpha_n r) J_0(\alpha_m r) dr = \frac{J_1^2(\alpha_n)}{2} \delta_{nm}$$

$$\text{Re}(D_n) = \frac{2}{a^2 J_1^2(\alpha_n)} \int_0^a r f(r) J_0(\alpha_n r/a) dr$$

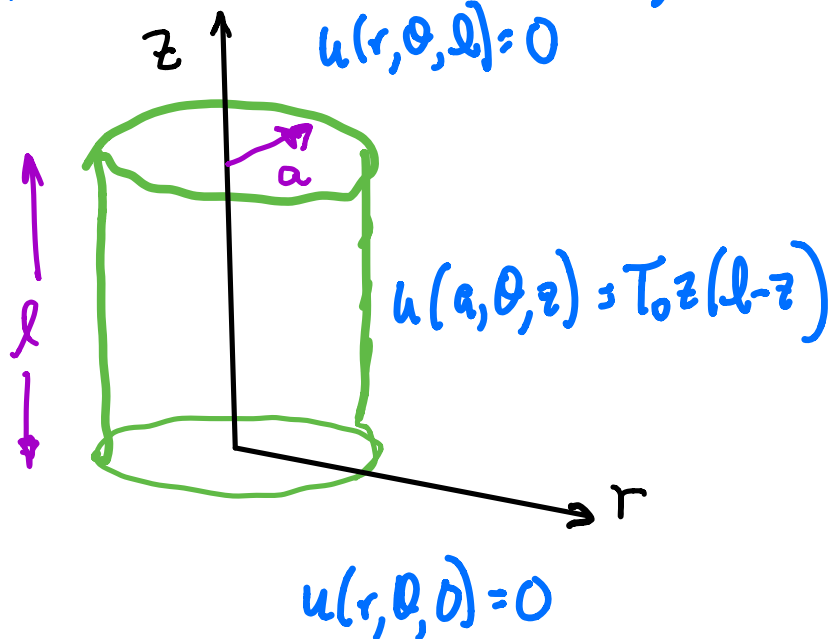
Differentiate $(\dagger)$ wrt t :

$$u_t(r, \theta, t) = \sum_{n=1}^{\infty} i\omega_n D_n J_0\left(\frac{\alpha_n r}{a}\right) e^{i\omega_n t}$$

$$u_f(r, \theta, 0) = g(r)$$

$$\therefore i \omega_n D_n = \frac{2}{a \alpha_n J_1'(\alpha_n)} \int_0^a r g(r) J_0(\alpha_n r/a) dr //$$

ex) Find the solution $u(r, \theta, z)$, $\Delta u = 0$
 $u(r, \theta, l) = 0$



Since no θ dependence $u = u(r, z)$

$$\Delta u = \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{\partial^2 u}{\partial z^2} = 0 \quad \text{PDE}$$

$$u(r,0) = u(r,l) = 0 \quad 0 \leq r \leq a$$

$$u(a,z) = T_0 z(l-z) \quad 0 \leq z \leq l$$

u is bounded inside the cylinder

$u = R(r) Z(z)$, sub into PDE:

$$Z R_{rr} + \frac{1}{r} Z R_r = - Z_{zz} R$$

divide by $Z R$:

$$\frac{R_{rr}}{R} + \frac{1}{r} \frac{R_r}{R} = - \frac{Z_{zz}}{Z} = + \lambda^2$$

$$- Z_{zz} = + \lambda^2 Z$$

$$\text{or } Z_{zz} + \lambda^2 Z = 0$$

$$r^2 R_{rr} + r R_r - \lambda^2 r^2 R = 0$$

The "R" equation has a solution
in terms

$$R(r) = C_1 I_0(\lambda r) + C_2 K_0(\lambda r)$$

I_0 and K_0 are a type of Bessel functions
known as Modified Bessel functions.

$$\begin{cases} z'' + \lambda^2 z = 0 \\ z(0) = 0 \quad z(L) = 0 \end{cases}$$

$$z_n = \sin \frac{n\pi z}{L} \quad n=1, 2, \dots$$

$$\lambda_n = \left(\frac{n\pi}{L} \right) \quad n=1, 2, \dots$$

$$\therefore R(r) = C_1 I_0\left(\frac{n\pi r}{l}\right) + C_2 K_0\left(\frac{n\pi r}{l}\right)$$

K_0 is unbounded at $r=0 \therefore C_2=0$

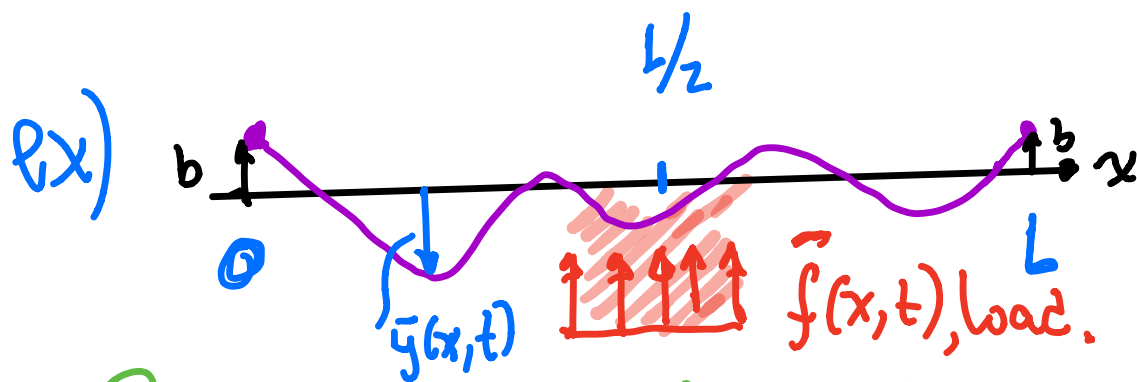
$$u(r, z) = \sum_{n=1}^{\infty} C_n I_0\left(\frac{n\pi r}{l}\right) \sin \frac{n\pi z}{l}$$

$$u(a, z) = \sum_{n=1}^{\infty} C_n I_0\left(\frac{n\pi a}{l}\right) \sin \frac{n\pi z}{l}$$

$$= T_0 z (l-z)$$

We use orthogonality of $\sin \frac{n\pi z}{l}$

$$C_n = \frac{T_0}{I_0 \left(\frac{n\pi a}{L} \right)} \frac{2}{L} \int_0^L z(L-z) \sin \frac{n\pi z}{L} dz$$



Beam, loaded in the middle and
hinged at the 2 end points

The beam is of length L .
 $\bar{y} = \bar{y}(x,t)$ is the displacement

$$\rho A \frac{\partial^2 \bar{y}}{\partial t^2} = \bar{f}(x,t) - \frac{\partial^2}{\partial x^2} \left(EI \frac{\partial^2 \bar{y}}{\partial x^2} \right) \quad (\text{PDE})$$

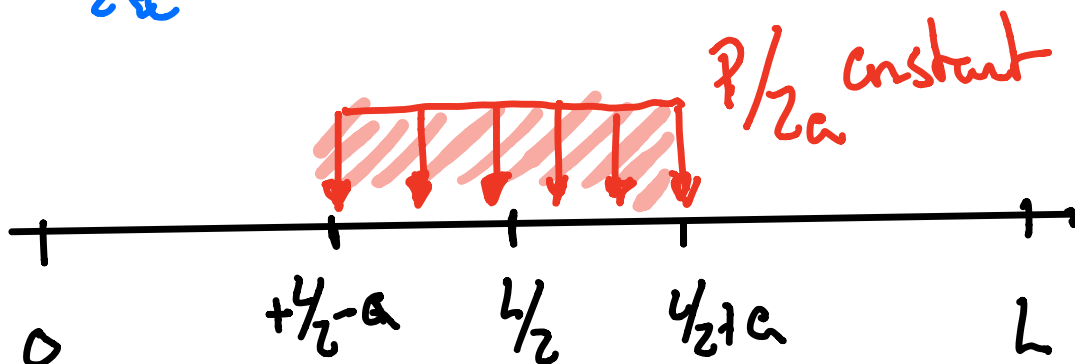
here $\Delta(x)$ is beam cross section
 $\rho(x)$ the density of beam
 $E(x)$ is the Young's modulus
 $I(x)$ is the inertia.

We'll assume Δ, ρ, E, I are constant.

LOAD:

$$\bar{f}(x,t) = \sin \omega t \begin{cases} \frac{P}{2a} & \frac{L}{2} - a < x < \frac{L}{2} + a \\ 0 & \text{otherwise} \end{cases}$$

$\frac{P}{2a}$ is constant



We require 4 B.C. (in space)

$$\bar{y}(0, t) = b \quad \frac{d^2 \bar{y}}{dx^2}(0, t) = 0$$

$$\bar{y}(L, t) = b \quad \frac{d^2 \bar{y}}{dx^2}(L, t) = 0$$

$$\text{let } \bar{y}(x, t) = \sin \omega t \, y(x)$$

substitute into the PDE

$$\underline{-\rho A \omega^2 \sin \omega t \, y(x)}$$

$$\rho A \frac{\partial^2 \bar{y}}{\partial t^2}$$

$$= f(x) \sin \omega t - \frac{\partial^2}{\partial x^2} \left(EI \frac{\partial^2 \bar{y}}{\partial x^2} \right) \sin \omega t$$

$$\text{where } f(x) = \begin{cases} \frac{P}{2a} & \frac{L}{2} - a < x < \frac{L}{2} + a \\ 0 & \end{cases}$$

factor out sin wt :

$$\frac{\partial^2}{\partial x^2} (EI y'') - \rho A y = f(x) \quad (*)$$

$$y = y(x)$$

$$\text{let } \beta^4 = \frac{\rho A}{EI} \quad \gamma = \frac{f(x)}{EI}$$

$$y^{iv} - \beta^4 y = \gamma \quad \text{is } (*)$$

$$y(0) = y(L) = b$$

$$y''(0) = y''(L) = 0$$

$$\text{let } y = y_I + y_{II}$$

$$\left\{ \begin{array}{l} y_I^{IV} - \beta^4 y_I = 0 \\ y_I(0) = y_I(L) = b \\ y_I''(0) = y_I''(L) = 0 \end{array} \right.$$

$$\left\{ \begin{array}{l} y_{II}^{IV} - \beta^4 y_{II} = f(x) \\ y_{II}(0) = y_{II}(L) = 0 \\ y_{II}'(0) = y_{II}'(L) = 0 \end{array} \right.$$

① Find eigenfunctions:

$$\begin{cases} u^{IV} - \beta^4 u = 0 & \text{D.E.} \\ u(0) = u(L) = u''(0) = u''(L) = 0 \end{cases}$$

The solution is found by assuming $u = e^{\alpha x}$ in D.E.

We get $\alpha^4 - \beta^4 = 0$. There are 2 real roots & 2 complex conjugate roots:

$$u = C_1 \sin \beta x + C_2 \cos \beta x + C_3 \sinh \beta x + C_4 \cosh \beta x$$

Apply B.C.

$$u(0) = 0 = C_2 + C_4 \quad u''(0) = -C_2 + C_4 = 0 \Rightarrow C_2 = C_4 = 0$$

$$u(L) = 0 = C_1 \sin \beta L + C_3 \sinh \beta L$$

$$u''(L) = 0 = -C_1 \beta^2 \sin \beta L + C_3 \beta^2 \sinh \beta L$$

$$\text{or } \sin \beta L = 0 \quad \therefore \beta_n = \frac{n\pi}{L} \quad n = 1, 2, \dots$$

$$\therefore u_n = \sin \frac{n\pi x}{L} \quad n = 1, 2, \dots$$

② Expand $f(x)$ in a series in u_n wrt $w(x) = 1$:

$$\frac{f}{2AEI} = \sum_{n=1}^{\infty} b_n \sin \frac{n\pi x}{L}$$

$$b_n = \frac{2}{L} \frac{P}{2aEI} \int_{y/2-a}^{y/2+a} \sin \frac{n\pi x}{L} dx$$

$$\therefore b_n = -\frac{2}{L} \frac{P}{2aEI} \left(\frac{L}{n\pi} \right) \cos \left(\frac{n\pi x}{L} \right) \Big|_{y/2-a}^{y/2+a}$$

$$\text{But } \cos \left(\frac{n\pi}{2} + \frac{n\pi a}{L} \right) - \cos \left(\frac{n\pi}{2} - \frac{n\pi a}{L} \right) = -2 \sin \frac{n\pi}{2} \sin \frac{n\pi a}{L}$$

$$\therefore b_n = \frac{2}{aEI} \frac{1}{n\pi} \sin \frac{n\pi}{2} \sin \frac{n\pi a}{L} \quad n = \text{even.}$$

③ Substitute series of $F(x)$ into CDE:

$$y_{xx}^{IV} - \beta^4 y_{xx} = \sum_{n=1}^{\infty} b_n \sin \frac{n\pi x}{L}$$

$$\text{Let } y_{xx} = \sum_{n=1}^{\infty} a_n \sin \frac{n\pi x}{L} = \sum_{n=1}^{\infty} a_n \sin \beta_n x$$

$$y_{xx}^{IV} = \sum_{n=1}^{\infty} a_n \beta_n^4 \sin \beta_n x$$

$$\therefore \sum_{n=1}^{\infty} [a_n \beta_n^4 \sin \beta_n x - \beta^4 a_n \sin \beta_n x] = \sum_{n=1}^{\infty} b_n \sin \beta_n x$$

$$\therefore \sum_{n=1}^{\infty} [a_n(\beta_n^4 - \beta^4) - b_n] \sin \beta_n x = 0$$

$$\therefore a_n = \frac{b_n}{\beta_n^4 - \beta^4}$$

$$y_{II} = \sum_{n=1}^{\infty} \frac{b_n}{\beta_n^4 - \beta^4} \frac{\sin n\pi x}{L}$$

$$\text{with } b_n = \frac{2P}{AEI} \frac{1}{n\pi} \frac{\sin \frac{n\pi}{2}}{2} \frac{\sin \frac{n\pi a}{L}}$$

$$\text{or } b_n = \frac{2PL}{EI} \frac{\sin \frac{n\pi}{2}}{2} \frac{\sin \theta}{\theta}, \quad \theta = \frac{n\pi a}{L}$$

$$\lim_{\theta \rightarrow 0} b_n = \frac{2PL}{EI} \frac{\sin \frac{n\pi}{2}}{2}$$

④ Solve y_I

$$(\#) \left\{ \begin{array}{l} y_{II}^{IV} - \beta^4 y_{II} = 0 \\ y_{II}(0) = y_{II}(L) = b \\ y_{II}''(0) = y_{II}''(L) = 0 \end{array} \right. \quad \text{We show that the solution is } y_{II} = b:$$

$$\text{let } y_{II} = b + v(x)$$

y_{II} requires v & derivatives are 0 at boundaries.

Substitute in $(\#)$:

$$v^{IV} - b\beta^4 - \beta^4 v = 0$$

$$(\dagger) \left\{ \begin{array}{l} v^{IV} - \beta^4 v = b\beta^4 \quad (\text{D.E.}) \\ v(0) = v(L) = v''(L) = v''(0) = 0 \end{array} \right.$$

by inspection,

$$v = \alpha_1 x^4 + \alpha_2 x^3 + \alpha_3 x^2 + \alpha_4 x + \alpha_5$$

$$v' = 4\alpha_1 x^3 + 3\alpha_2 x^2 + 2\alpha_3 x + \alpha_4$$

$$v'' = 12\alpha_1 x^2 + 6\alpha_2 x + 2\alpha_3$$

$$v''' = 24a_1x + 6a_2$$

$$v'''' = 24a_1$$

The D.E.

$$24a_1 - \beta^4(a_1x^4 + a_2x^3 + a_3x^2 + a_4x + a_5) = b\beta^4$$

$$\therefore \text{let } 24a_1 - \beta^4a_5 = b\beta^4$$

$$v''(L) = 0 = 12a_1L^2 + 6La_2 + 2a_3 = 0$$

$$v(L) = 0 = a_1L^4 + a_2L^3 + a_3L^2 + a_4L + a_5 = 0$$

$$\text{let } a_3 = 0, a_4 = 0$$

$$a_2 = -2a_1L$$

$$a_5 = \frac{24a_1 - b\beta^4}{\beta^4}$$

$$a_1L^4 - 2a_1L^4 - \frac{24a_1 - b\beta^4}{\beta^4} = 0$$

$$\therefore a_1 \left(-L^4 - \frac{24}{\beta^4} \right) = b$$

$$\alpha_1 (\beta^4 L^4 - 24) = -\beta^4 b$$

$$\alpha_1 = \frac{\beta^4 b}{24 - \beta^4 L^4}$$

$$\alpha_5 = \frac{\beta^2 b}{24 - \beta^4 L^4} - b$$

$$\alpha_2 = \frac{-2L\beta^4 b}{24 - \beta^4 L^4}$$

$$\therefore v(x) = \frac{\beta^4 b}{24 - \beta^4 L^4} x^4 - \frac{2L\beta^4 b}{24 - \beta^4 L^4} x^3 - \frac{\beta^2 b}{24 - \beta^4 L^4} - b$$

$$y_5 = b + v(x)$$

$$y_1 = \frac{\beta^4 b}{24 - \beta^4 L^4} x^4 - \frac{2L\beta^4 b}{24 - \beta^4 L^4} x^3 - \frac{\beta^2 b}{24 - \beta^4 L^4} //$$

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