Two functions f(x) & g(x), complex (or real) are sevel to be orthogonal with respect to a weight fination with an interval I=[e,b] if $\int f(x) \overline{g(x)} w(x) dx = 0$ where (.) is complex conjugate. 6 (or $\int \overline{f}(x)g(x)w(x)dx = 0$ The weight function $w(x) \ge 0$ for $x \in (a, b)$ and $\int_{a}^{b} w(x) dx = A = 0$, enumber.

$$e_{X} = \frac{1}{2} \left(\frac{1}{2} e^{i(n-n)} \right) = e^{i(n-n)} = e^{i(n-n)} = \frac{1}{2} \left[\frac{1}{2} e^{i(n-n)} \right] = \frac{1}{2} \left[\frac{1}$$

Ruh: What velies of m bin wake this
equal to 0?
(i) Case might let
$$p=m-n$$

 $h(q) \equiv \frac{2}{p} \sin p \pi$
if p is not an integer ($\neq 0$)
 $h(p)$ is generally not equal to 0.
(ii) $p=0$: for this case, take limit
as $p=0$, use L'Hopital's Rule:
 $\lim_{p \to 0} \frac{2}{p} \sin p \pi = \lim_{p \to 0} \frac{2\pi \cos p \pi}{1}$
 $= 2\pi$

(iii) IF
$$p = \pm 1, \pm 2, \pm 3...$$

 $p = h - n$
 $\frac{2}{(m - n)}$ Sin $[(m - n) J_3]$
 $= \frac{2}{p} s_{1} m p J_3 = 0$
 \therefore For $Q_n(x) = e^{inx}$, $Q_m(x) = e^{inx}$
 $\int_{-\pi}^{\pi} Q_n(x) Q_m(x) \int_{-\pi}^{0} (x) = e^{inx}$
 $\int_{-\pi}^{\pi} Q_n(x) Q_m(x) \int_{-\pi}^{0} (x) = e^{inx}$
For $m, n \in \mathbb{Z}$

STURM- LIOUVILLE PROBLEM & special family of boundary value problems, they are all linear and are ordinary differental equations of even order. Kuch: For the special cose of order? we can always transform the differential equation into storn-Lionville form. Focus un 2nd order Stum-Lionille Problems (SL) BVP: Consider some nice function y(x) Such that

 $(\ddagger) \frac{d}{dx} \left(p(x) \frac{dy}{dx} \right) + \left[q(x) + \lambda w(x) \right] y(x) = 0$ on a bounded interval x=a, x=b bra, plus 2 boundary anditus (to be specified shorthy) See MTH481/581 für more debails P, g, w are real functions $bf = \frac{1}{2}(pcx) = \frac{1}{2}(pcx)$ Her (7) $(*) \quad \forall y + \forall wy = 0$

We can take any 2nd order linear
BVP ode and transform it into (*):
(***)
$$a_0(x)y''(x) + a_1(x)y'(x)$$

 $+[a_2(x) + \lambda a_3(x)]y(x) = 0$
 $(a_0(x) \neq 0 \text{ for}$
 $x \in [a,b]$
if $p(x) \equiv e$
 $g(x) \equiv \frac{a_1}{a_0}P$
 $w(x) \equiv \frac{a_2(x)}{a_0(x)}P$ flue (***) com
be written as (**) and can due be
 $0f$ SL form over some intervel $[a_1b]$
provided $\frac{a_3(x)}{a_0(x)} = 0$

SL
$$\begin{cases} dy + \lambda w G y = 0 \\ q_1 y(a) + q_2 y'(a) = 0 \\ \beta_1 y(b) + \beta_2 y'(b) = 0 \end{cases}$$

if $p(x) > 0$, $w(c) > 0$, λ is a number
 $p, q_1 w$ are antinvous on $\chi \in [a_1b]$
 $q_1, q_2, \beta_1, \beta_2$ are numbers.
We will now show that solutions fo
SL BVP have the property
 $\int_{a}^{b} w(x) q_1(x) q_2(x) dx = 0$

Where Pi(x) solves JP; + Xiw (2) P; = O 9;(x) solves 24; + 2; w(x) q- =0 $\lambda_i & \lambda_j$ are 2 numbers plus some Bonday conditions q. (x) and q. (x) are (PEAL) orthogonal finations. We will show that the solutions to the SL problem, with Litteent X5 have the property Jw(x) (f; (x) (f; (+) dx = 0

i. l. G; (x) & G; (x), solutions to SL one orthogoach w.r. t. w(x) on the interval [2,5]: to see this:

(1) $\frac{d}{dx}(p(x)\partial_x q_i) + (q+\lambda_iw)q_i = 0$ (2) $\frac{d}{dx}(p(x)\partial_x q_j) + (q+\lambda_jw)q_j = 0$ Hultiph (1) by q_j ad (2) by q_i ad Subtract:

 $\begin{aligned} \mathcal{Q}_{j}(p \partial_{x} \mathcal{Q}_{j})_{x} &= \mathcal{Q}_{i}(p \partial_{x} \mathcal{Q}_{j})_{x} \\ &+ (\lambda_{i} - \lambda_{j}) w \mathcal{Q}_{i} \mathcal{Q}_{j} = 0 \end{aligned}$

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 $(\lambda_{i}-\lambda_{j})$ w $q_{i}q_{j} = q_{i}(p \partial_{x}q_{j})_{x} - q_{j}(p \partial_{x}q_{j})_{x}$

integrate by parts (JBP):
(
$$\lambda_i$$
- λ_j) $\int_{a}^{b} w q_i q_j dx = p \partial_x q_j q_i |_{a}^{b} - \int_{p} q_{ix} q_{jx} dx$
 $-p \partial_x q_i q_j |_{a}^{b} + \int_{p} q_{ix} q_{jx} dx$
The IBP calculation. Take
cancel: let $u = q_i dv = (q^2 \times q_j) x$
 $k_{u} du = q_i dv = (q^2 \times q_j) x$
 $k_{u} du = q_i dv = (q^2 \times q_j) x$
 $k_{u} du = q_i dv = (q^2 \times q_j) dx$
 $(\lambda_i - \lambda_j) \int_{a}^{b} w(x) q_j (x) q_j (x) dx$ $so \int_{a}^{b} q_i (b = q_j) dx$
 $= (p [\partial_x q_j : q_i - \partial_x q_i q_j]_{a}^{b} = p q_i \partial_x q_j]_{a}^{b}$
 $= (p [\partial_x q_j : q_i - \partial_x q_i q_j]_{a}^{b} = -\int_{p} q_{ix} q_{jx} dx$
Rul: Since $\lambda_i - \lambda_j \neq 0$, thu
 $\int_{w(d_j)}^{b} (q_j (x) q_j dx) dx = 0$
provided terms inderlated are zero.
That is, provided the boundary inditions
halve the rhs equal to zero.

CASE (b) if $[q]_{;=0}$ at $x_{=q} & & x_{=b}$ $[q]_{;=0}$ DIRICULET B.C. $\Rightarrow \int w(x) Q; Q; dx = 0$ CASE (B) if $\int \Phi_{ix} = 0$ of x=a, x=b $(\Phi_{ix} = 0)$ NEUMANN B.C. ⇒ (we)Q;Q;dx=0 $C_{NSE} \bigcirc if [P_i + \delta P_{ix} = 0 \quad ROBIN or \\ P_i + \delta P_{ix} = 0 \quad MVXED B.C.$ ⇒ ∫w(x)Q;Q;dx = O Ruh: could also have and institus of Dirichlet B.C. et one end & Neumann et the other. Ruch: #if p:0 at x=e, x=3 the Sw q; q; dx=0 * if Pishite at x=a,b ad P'rp P'tend

to zero at
$$x=a,b$$
 the P_i, P_j are orthing
wrt w as well.
* if $p(a)=p(b) \Rightarrow \int w P_i P_j dx = 0$
 $\int P_i(a) = P_i(b) a d P_i(c) = P_i'(b)$
 $i P_j(c) = P_j(b) a d P_i(c) = P_i'(b)$
 $i.e. periodiz B.C.j.e.$
if $P(c) = P(x+L)$ where $L = b \cdot a$.
 $P(x) \int y'' + a^2 y = 0$ here $y = y(x)$
 $(4) \int y(0) = y(2) = 0$ λ is a number
or $[P Y_x]_x + \lambda y = 0$
 $q = 0, P = 1, \lambda = a^2, W = 1$
we see that (4) is a S.L. problem.
The solution of the ODE $y'' + a^2 y = 0$

$$y = 1 \cos 2x + B \sin 2x$$

apply B.C. $y = B \sin n \frac{\pi x}{Q}$
i.e. $\alpha = \frac{\pi \pi}{Q}$ $n = 1/2...$
B is on orbitrary creatent.
i.e. $y(x) = B P_n(x)$
 $R_n(x) = \sinh n \frac{\pi \pi}{Q}$ $\ln = \frac{\pi \pi}{Q}$
and they are orthogonal for different n
over $x = [0, R]$ with $w = 1$
i.e. $\int_{0}^{Q} \sinh \pi \pi x \, dx = \begin{cases} 0 & \sin \pi \pi x \\ Q & Q \end{cases}$

We say that Pn(x) are eigenfination

of the SL problem with eigenvalues

$$\alpha_n = \lambda_n = \frac{n}{R} \quad n = 1, 2...,$$

Because
 $\int -\Omega P_n = \lambda_n CP_n$
with $P_n(0) = P_n(\Omega) = 0$
 $\Omega = \int_{X^2}^{Z} , \quad \lambda_n = \frac{n^2 \pi^2}{R^2}$
We sometrives hid it convenient to normalize
the eigenformations, then we call then conthemorphis:
For the excerpte we just did
(1) $\int P_n(\Omega) P_n(\Omega) dx = \begin{cases} 1 & \text{if } n = n \\ 0 & \text{otherwise}. \end{cases}$
Defore $Snm = \begin{cases} 1 & n = n \\ 0 & n \neq n \end{cases}$ then $\int_{\Omega}^{R} P_n(\Omega) P_n(\Omega) dx = Snm.$

Let $\hat{\varphi}_{j}(x) = \lambda \hat{\varphi}_{j}(x)$, λ is the "normalization factor." Find λ s.t. ($\hat{\varphi}$) is true: $\int_{0}^{Q} \hat{\varphi}_{i}(x) \hat{\varphi}_{j}(x) dx = \frac{1}{2} \delta_{ij}$ or $\frac{2}{2} \int_{0}^{Q} \hat{\varphi}_{i}(x) \hat{\varphi}_{j}(x) dx = \delta_{ij}$ $\therefore (\frac{2}{2} \hat{\varphi}_{i}(x) = \hat{\varphi}_{i}(x))$

EXPANSION OF & FUNCTION IN A SERIES OF ORTHOGONAL FUNCTIONS, on & FIXED INTERVAL I=[R,6]

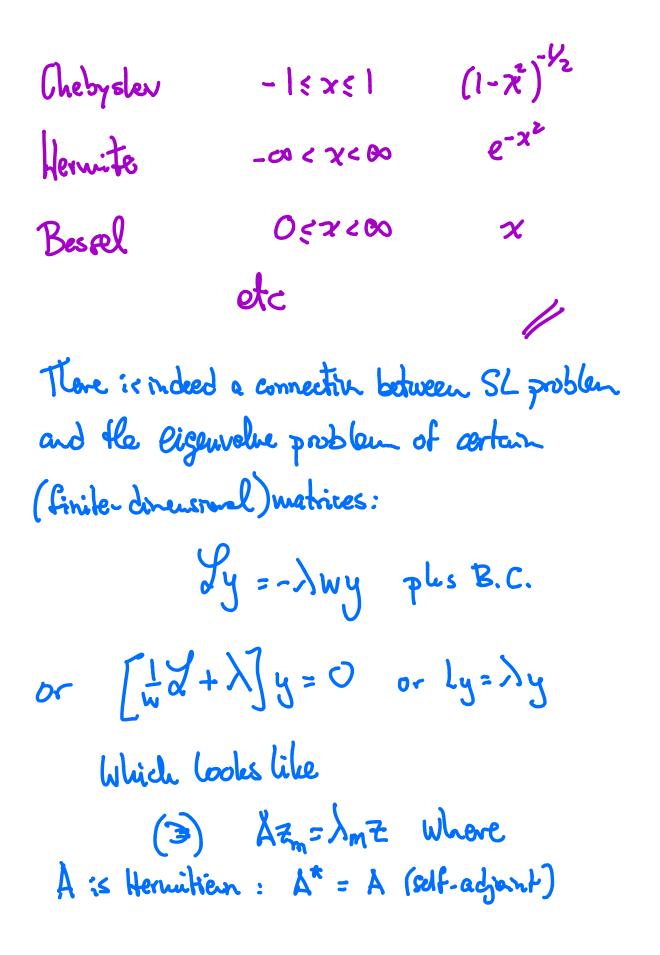
Suppose we know of a family of osthyand bucks {Pn(x)} on a fixed interval [2,5]. Also,

suppose f(x) has boundary values crisisterst with gra. Then $f(x) = \sum_{n=1}^{\infty} \lambda_n \varphi_n(x)$ where by one coefficients. This is called an eigenfunction exponsion Of f(x) in terms of eigenfunctions {Pa(x)}, orthogoal over [a,b] w.r.t WGe), the weight (WGO > 0) with Jwkedx = C, a number. We will also require that $\int_{a}^{b} |f(x)|^{2} dx < \infty$

i.e. f(x) is an Lz function. To find the inknown we ficients An: $w(x) f(x) = \sum_{n=1}^{\infty} \Delta_n w(x) q_n(x)$ Hultiply both sides by 9m (2) and Integrate from x=a, b: Sw(x) f(x) cfm (x) dx $= \int_{n=0}^{b} \sum_{n=0}^{\infty} A_n w(x) P_n(x) P_n(x) dx$ $= \sum_{n=1}^{\infty} A_n \int w(x) c_n(x) c_n(x) dx$ = $A_m \int w(x) \varphi_m(x) \varphi_m(x) dx$

= Am Q
where
$$Q = \int w(\omega) Q_m^2(\omega) dx$$

Solving for $\Delta m = \frac{1}{Q} \int w(\omega) f(\omega) Q_m(\omega) dx$
 $M_m = \frac{1}{Q} \int w(\omega) f(\omega) Q_m(\omega) dx$
 $M = 0,1,2...$
There are many familized of orthogonal functions
there are associated with SL problems
Here are associated with SL problems
Family interval weight with
 $Sihmmer$, commer $x = \left[\frac{Q}{2}, \frac{Q}{2}\right] = 1$
legendre $-1 \le x \le 1 = 1$



A is nxn metrix. ZER, m=0,1... and h_GR', En iske eigenvector and Im the essocieted evelve. Like the SL, (I) has real eigen finettes and real eigenvalues. Like the SL, (F) has an minite, unique eigenvelves that can be arrayed as 700 x1 x x2 x ... xn x ... where lim ton = 00 So each eigenvalue In has only 1 eigenfinction Zn.

Ex) 6izenfinction expansion: let
$$f(x) = x$$

let $I = (0, 2]$
expand $f(x)$ subs a series of eizenfunctions
of the SL problem
SL $\begin{cases} y'' + k^2y = 0 \\ y(0) = y(2) = 0 \\ y(0) = y(2) = 0 \\ y(0) = 0$

$$\int (x) = \sum_{n=0}^{\infty} a_n q_n(x)$$

$$\lim_{h \to 0} \int f(x) q_{m}(x) = \sum_{k=0}^{\infty} a_n q_n(x) q_n(x) u(x)$$

$$\lim_{h \to 0} \int g(x) q_{m}(x) = \sum_{k=0}^{\infty} a_n q_n(x) q_n(x) u(x)$$

$$\lim_{k \to 0} \int g(x) q_m(x) dx = \sum_{k=0}^{\infty} a_n \int g(x) dx$$

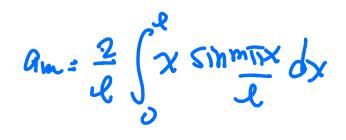
$$= a_{nn} \int g(x) dx$$

$$= a_{nn} \int g(x) dx dx$$

$$= a_{nn} \int g(x) dx$$

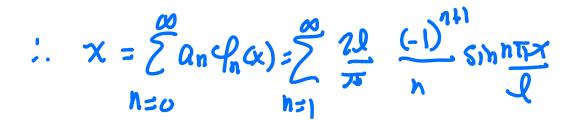
$$= a_{nn} \int g(x) dx$$

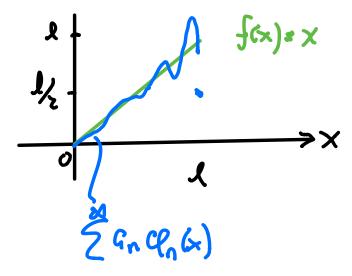
$$= a_{nn} \int g(x) dx$$



(prercise)

$$G_{m} = -\frac{2l}{m\tau_{i}} Cosm\tau_{i} = \frac{2l}{m\tau_{i}} (-1)^{m+1}$$







Rink: Since for isnot periodic, the eigenforther
expansion will not an very in the unitary
Noran:
lim
$$|f(x) - \sum_{n=0}^{N} a_n c_n(x)|$$
 we will not
 $N \rightarrow co$ set 0.
for $x \in [9, 0]$
instead
 $\lim_{N \rightarrow co} |f(x) - \sum_{n=0}^{N} a_n c_n(x)|^2 dx = 0$ be average

See notes for tricles muching periodic extensions