

16.1.1 The potential for the harmonic oscillator is

$$V(x) = \frac{1}{2} m \omega^2 x^2$$

Let's try a Gaussian trial wave function in the variational method to estimate the ground state energy and ground state wave function. The Hamiltonian is

$$H = T + V(x) = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2$$

The trial function is

$$\psi(x, \alpha) = e^{-\alpha x^2}$$

The energy functional

$$E[\psi] = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} \geq E_0$$

is then

$$E[\psi] = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} = \frac{\int e^{-\alpha x^2} \left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m \omega^2 x^2 \right) e^{-\alpha x^2} dx}{\int e^{-2\alpha x^2} dx}$$

The denominator is

$$\int_{-\infty}^{\infty} e^{-2\alpha x^2} dx = \sqrt{\frac{\pi}{2\alpha}}$$

The numerator is

$$\begin{aligned} \int_{-\infty}^{\infty} e^{-\alpha x^2} \left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m \omega^2 x^2 \right) e^{-\alpha x^2} dx &= \int_{-\infty}^{\infty} e^{-\alpha x^2} \left(-\frac{\hbar^2}{2m} (4\alpha^2 x^2 - 2\alpha) + \frac{1}{2} m \omega^2 x^2 \right) e^{-\alpha x^2} dx \\ &= \int_{-\infty}^{\infty} \left(\frac{\hbar^2 \alpha}{m} + \left(-\frac{2\hbar^2 \alpha^2}{m} + \frac{1}{2} m \omega^2 \right) x^2 \right) e^{-2\alpha x^2} dx \\ &= \frac{\hbar^2 \alpha}{m} \sqrt{\frac{\pi}{2\alpha}} + \left(-\frac{2\hbar^2 \alpha^2}{m} + \frac{1}{2} m \omega^2 \right) \frac{1}{4} \sqrt{\frac{\pi}{2\alpha^3}} \\ &= \sqrt{\frac{\pi}{2\alpha}} \left(\frac{\hbar^2 \alpha}{2m} + \frac{m \omega^2}{8\alpha} \right) \end{aligned}$$

Hence, the energy function is

$$E(\alpha) = \frac{\hbar^2 \alpha}{2m} + \frac{m \omega^2}{8\alpha}$$

Minimize:

$$\frac{d}{d\alpha} E(\alpha) = \frac{\hbar^2}{2m} - \frac{m\omega^2}{8\alpha^2} = 0 \Rightarrow \alpha_0 = \frac{m\omega}{2\hbar}$$

The minimum energy is thus

$$E(\alpha_0) = \frac{\hbar\omega}{4} + \frac{\hbar\omega}{4} = \frac{1}{2}\hbar\omega$$

which turns out to be the exact (known) answer.

16.1.3 The potential for the delta function system is

$$V(x) = -aV_0\delta(x)$$

Let's try a Gaussian trial wave function in the variational method to estimate the ground state energy. The Hamiltonian is

$$H = T + V(x) = \frac{p^2}{2m} - aV_0\delta(x)$$

The trial function is

$$\psi(x, \alpha) = e^{-\alpha x^2}$$

The energy functional

$$E[\psi] = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} \geq E_0$$

is then

$$E[\psi] = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} = \frac{\int e^{-\alpha x^2} \left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} - aV_0\delta(x) \right) e^{-\alpha x^2} dx}{\int e^{-2\alpha x^2} dx}$$

The denominator is

$$\int_{-\infty}^{\infty} e^{-2\alpha x^2} dx = \sqrt{\frac{\pi}{2\alpha}}$$

The numerator is

$$\begin{aligned} \int_{-\infty}^{\infty} e^{-\alpha x^2} \left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} - aV_0 \delta(x) \right) e^{-\alpha x^2} dx &= \int_{-\infty}^{\infty} e^{-\alpha x^2} \left(-\frac{\hbar^2}{2m} (4\alpha^2 x^2 - 2\alpha) - aV_0 \delta(x) \right) e^{-\alpha x^2} dx \\ &= \int_{-\infty}^{\infty} \left(\frac{\hbar^2 \alpha}{m} - \frac{2\hbar^2 \alpha^2}{m} x^2 - aV_0 \delta(x) \right) e^{-2\alpha x^2} dx \\ &= \frac{\hbar^2 \alpha}{m} \sqrt{\frac{\pi}{2\alpha}} + \left(-\frac{2\hbar^2 \alpha^2}{m} \right) \frac{1}{4} \sqrt{\frac{\pi}{2\alpha^3}} - aV_0 \\ &= \sqrt{\frac{\pi}{2\alpha}} \left(\frac{\hbar^2 \alpha}{2m} \right) - aV_0 \end{aligned}$$

Hence, the energy function is

$$E(\alpha) = \frac{\hbar^2 \alpha}{2m} - aV_0 \sqrt{\frac{2\alpha}{\pi}}$$

Minimize:

$$\frac{d}{d\alpha} E(\alpha) = \frac{\hbar^2}{2m} - aV_0 \frac{1}{2} \sqrt{\frac{2}{\pi\alpha}} = 0 \Rightarrow \alpha_0 = \frac{2m^2 a^2 V_0^2}{\pi \hbar^4}$$

The minimum energy is thus

$$\begin{aligned} E(\alpha_0) &= \frac{\hbar^2 \alpha}{2m} - aV_0 \sqrt{\frac{2\alpha}{\pi}} = \frac{\hbar^2}{2m} \frac{2m^2 a^2 V_0^2}{\pi \hbar^4} - aV_0 \sqrt{\frac{2}{\pi}} \frac{\sqrt{2} m a V_0}{\sqrt{\pi} \hbar^2} \\ &= \frac{m a^2 V_0^2}{\pi \hbar^2} - \frac{2 m a^2 V_0^2}{\pi \hbar^2} = -\frac{m a^2 V_0^2}{\pi \hbar^2} \approx -(0.32) \frac{m a^2 V_0^2}{\hbar^2} \end{aligned}$$

The exact (known) answer is

$$E_{gnd} = -\frac{m a^2 V_0^2}{2 \hbar^2} = -(0.5) \frac{m a^2 V_0^2}{\hbar^2}$$

so our estimate is somewhat higher, as expected.

16.2.4 The wave function on the outside of the barrier (at x_e) in terms of the value inside (at x_0) is (Eq.16.2.24)

$$\psi(x_e) = \psi(x_0) \exp \left[\frac{i}{\hbar} \int_{x_0}^{x_e} i \sqrt{2m(V(x) - E)} dx \right] \equiv \psi(x_0) e^{-\gamma/2}$$

The lifetime is (Eq. 16.2.26, with $V_0 = 0$ for our case)

$$\tau = \frac{1}{R} = \frac{2m x_0}{\sqrt{2mE}} e^{\gamma}$$

The potential is just the Coulomb potential of the alpha particle in the presence of the nucleus with charge Z after emission:

$$V(x) = \frac{(Z_\alpha e)(Ze)}{x} = \frac{2Ze^2}{x}$$

For an alpha particle of energy E , the outer turning point is

$$\frac{2Ze^2}{x_e} = E \Rightarrow x_e = \frac{2Ze^2}{E}$$

Now find γ :

$$\begin{aligned} \gamma &= -2 \frac{i}{\hbar} \int_{x_0}^{x_e} i \sqrt{2m(V(x) - E)} dx = \frac{2}{\hbar} \int_{x_0}^{x_e} \sqrt{2m \left(\frac{2Ze^2}{x} - E \right)} dx \\ &= \frac{2}{\hbar} \int_{x_0}^{x_e} \sqrt{2m \left(\frac{x_e E}{x} - E \right)} dx = \frac{2}{\hbar} \sqrt{2mE} \int_{x_0}^{x_e} \sqrt{\left(\frac{x_e}{x} - 1 \right)} dx \end{aligned}$$

Let $u = x/x_e$ and $y = x_0/x_e$, and note that $v = \sqrt{2E/m}$:

$$\begin{aligned} \gamma &= \frac{2x_e}{\hbar} \sqrt{2mE} \int_y^1 \sqrt{\left(\frac{1}{u} - 1 \right)} du = \frac{2}{\hbar} \frac{2Ze^2}{E} \sqrt{2mE} \int_y^1 \sqrt{\left(\frac{1}{u} - 1 \right)} du \\ &= \frac{8Ze^2}{\hbar v} \int_y^1 \sqrt{\left(\frac{1}{u} - 1 \right)} du \end{aligned}$$

To do the integral, make the substitution $u = \sin^2 z$

$$\begin{aligned} \gamma &= \frac{8Ze^2}{\hbar v} \int_{\sin^{-1}\sqrt{y}}^{\pi/2} \sqrt{\left(\frac{1}{\sin^2 z} - 1 \right)} 2 \sin z \cos z dz = \frac{8Ze^2}{\hbar v} 2 \int_{\sin^{-1}\sqrt{y}}^{\pi/2} \cos^2 z dz \\ &= \frac{8Ze^2}{\hbar v} 2 \left[\frac{z}{2} + \frac{1}{4} \sin 2z \right]_{\sin^{-1}\sqrt{y}}^{\pi/2} = \frac{8Ze^2}{\hbar v} \left[z + \sin z \cos z \right]_{\sin^{-1}\sqrt{y}}^{\pi/2} \\ &= \frac{8Ze^2}{\hbar v} \left[\frac{\pi}{2} - \sin^{-1} \sqrt{y} - \sin(\sin^{-1} \sqrt{y}) \cos(\sin^{-1} \sqrt{y}) \right] \\ &= \frac{8Ze^2}{\hbar v} \left[\frac{\pi}{2} - \sin^{-1} \sqrt{y} - \sqrt{y} \sqrt{1-y} \right] \\ &= \frac{8Ze^2}{\hbar v} \left[\cos^{-1} \sqrt{y} - \sqrt{y} \sqrt{1-y} \right] \end{aligned}$$

Let's first estimate y . It is convenient to recall that $13.6eV = e^2/2a_0$:

$$y = \frac{x_0}{x_e} = \frac{x_0}{\frac{2Ze^2}{E}} = \frac{Ex_0}{2Ze^2} = \frac{4.2MeV(10^{-12}cm)}{2(90)(2 \times 0.0529nm)13.6eV} = 0.162$$

Hence we can use the small angle approximation on the trig function to get

$$\begin{aligned}\gamma &\approx \frac{8Ze^2}{\hbar v} \left[\frac{\pi}{2} - \sqrt{y} - \sqrt{y(1-y)} \right] \approx \frac{8Ze^2}{\hbar \sqrt{2E/m}} \left[\frac{\pi}{2} - \sqrt{y} - \sqrt{y(1-y)} \right] \\ &\approx \frac{8Z}{\sqrt{2}} \frac{e^2}{\hbar c} \sqrt{\frac{mc^2}{E}} \left[\frac{\pi}{2} - \sqrt{y} - \sqrt{y(1-y)} \right] \approx \frac{8Z}{\sqrt{2}} \alpha \sqrt{\frac{mc^2}{E}} \left[\frac{\pi}{2} - \sqrt{y} - \sqrt{y(1-y)} \right]\end{aligned}$$

Putting in numbers gives (m is the mass of the alpha: $3.73 \text{ GeV}/c^2$)

$$\begin{aligned}\gamma &\approx \frac{8(90)}{\sqrt{2}} \frac{1}{137} \sqrt{\frac{3.73 \text{ GeV}}{4.2 \text{ MeV}}} \left[\frac{\pi}{2} - \sqrt{0.162} - \sqrt{0.162(1-0.162)} \right] \approx 110.7 [0.7998] \\ &\approx 88.5\end{aligned}$$

and a lifetime

$$\begin{aligned}\tau &= \frac{1}{R} = \frac{2mx_0}{\sqrt{2mE}} e^\gamma = \sqrt{2} \frac{\sqrt{mc^2}}{\sqrt{E}} \frac{x_0}{c} e^\gamma \\ &\approx \sqrt{2} \sqrt{\frac{3.73 \text{ GeV}}{4.2 \text{ MeV}}} \frac{10^{-12} \text{ cm}}{3 \times 10^{10} \text{ cms}^{-1}} e^{88.5} \approx 3.8 \times 10^{17} \text{ s} \approx 1.2 \times 10^{10} \text{ years}\end{aligned}$$

16.2.7 The bound states in the WKB approximation are found by integrating the momentum between the classical turning points:

$$\int_{x_1}^{x_2} p(x) dx = (n + \frac{1}{2}) \pi \hbar$$

For the harmonic oscillator the momentum is

$$E = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2 \Rightarrow p(x) = \sqrt{2m(E - \frac{1}{2} m \omega^2 x^2)} = m\omega \sqrt{\frac{2E}{m\omega^2} - x^2}$$

and the turning points are

$$p(x_{1,2}) = 0 = m\omega \sqrt{\frac{2E}{m\omega^2} - x^2} \Rightarrow x_2 = -x_1 = \sqrt{\frac{2E}{m\omega^2}}$$

Applying the WKB condition gives

$$\begin{aligned}(n + \frac{1}{2}) \pi \hbar &= \int_{x_1}^{x_2} p(x) dx = m\omega \int_{-x_2}^{x_2} \sqrt{x_2^2 - x^2} dx \\ &= m\omega \frac{1}{2} \left[x \sqrt{x_2^2 - x^2} + x_2^2 \sin^{-1} \left(\frac{x}{x_2} \right) \right]_{-x_2}^{x_2} = m\omega \frac{1}{2} \left[x_2^2 \frac{\pi}{2} - x_2^2 \left(-\frac{\pi}{2} \right) \right] \\ &= \frac{1}{2} m\omega \pi x_2^2 = \frac{1}{2} m\omega \pi \frac{2E}{m\omega^2} = \pi \frac{E}{\omega}\end{aligned}$$

Solving for the energy gives

$$E_n = (n + \frac{1}{2}) \hbar \omega$$

which is the exact answer.