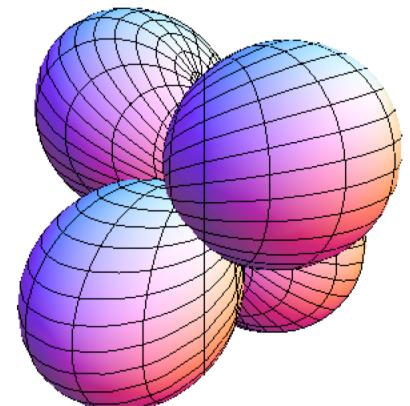
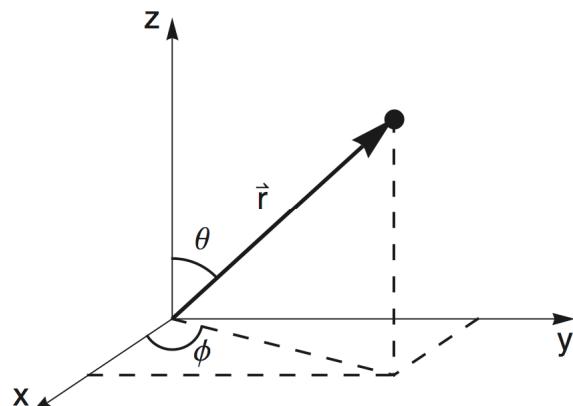


$$\psi_{nlm}(r, \theta, \phi) = R_{nl}(r) \Theta_\ell^m(\theta) \Phi_m(\phi)$$

$$\psi(r, \theta, \phi, t) = \sum_{nlm} c_{nlm} R_{nl}(r) \Theta_\ell^m(\theta) \Phi_m(\phi) e^{-iE_n t/\hbar}$$

The quantum central force problem:  
three dimensions & separation of  
variables

Reading: McIntyre 7.1-7.4, Appendix E



# The 2-body problem

$$H = \left[ \frac{\vec{p}_1^2}{2m_1} + \frac{\vec{p}_2^2}{2m_2} \right] + V(\vec{r}_1, \vec{r}_2)$$

$$H = \left[ \frac{\vec{P}_{CM}^2}{2M} + \frac{\vec{p}_{rel}^2}{2\mu} \right] + V(r)$$

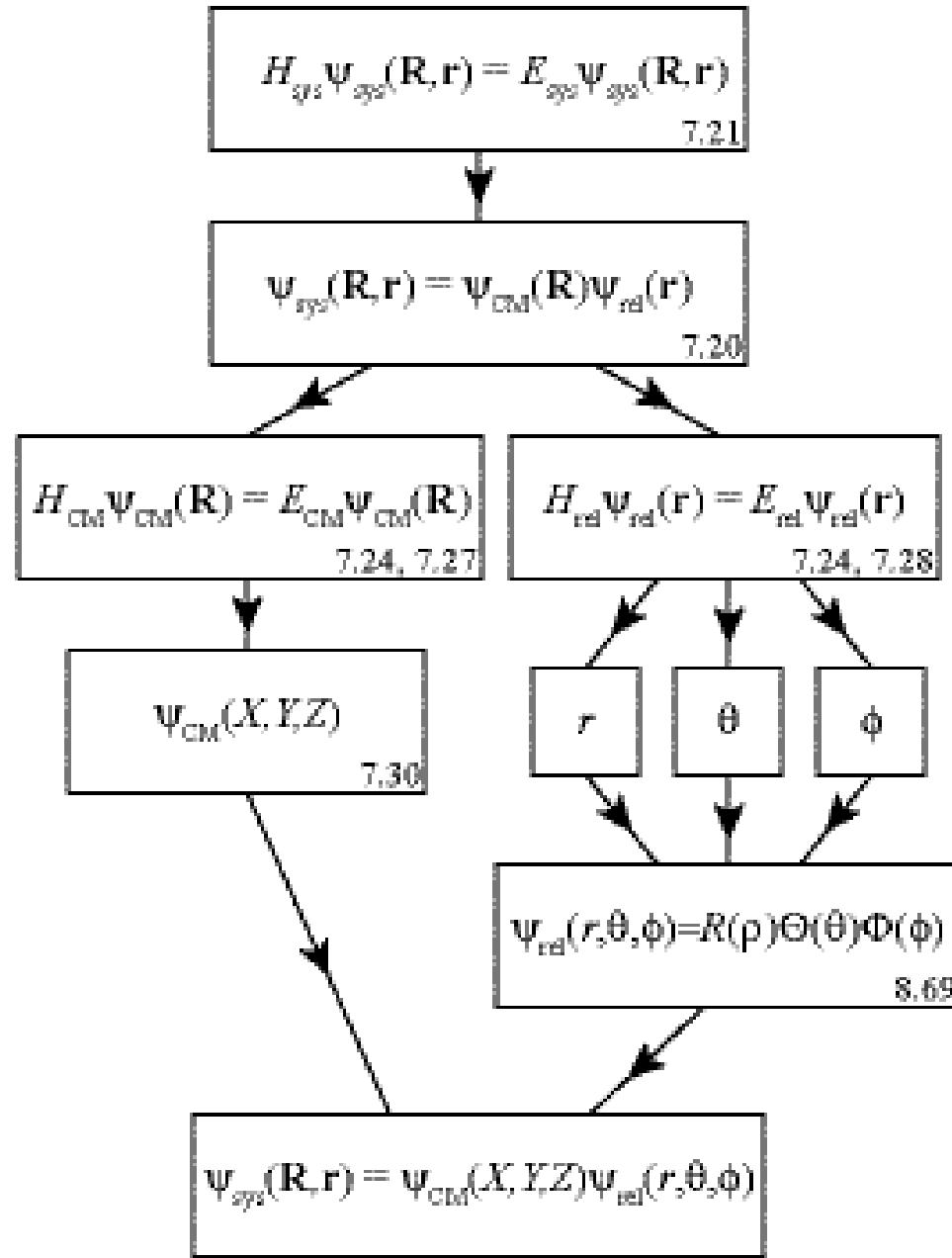
$$p_{(rel)} \doteq -i\hbar \left( \mathbf{i} \frac{\partial}{\partial x} + \mathbf{j} \frac{\partial}{\partial y} + \mathbf{k} \frac{\partial}{\partial z} \right) = -i\hbar \nabla_{(rel)}$$

$$H_{(rel)} \doteq -\frac{\hbar^2}{2\mu} \nabla_{(rel)}^2 + V(r_{(rel)})$$

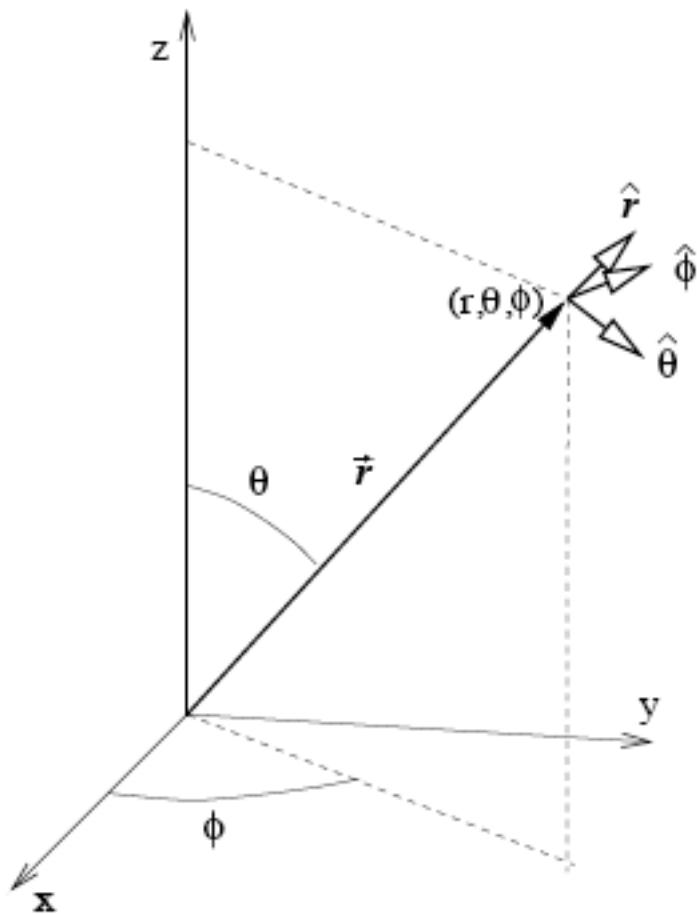
- Separates into a CoM part and a relative part with same definitions as in classical problem for center of mass, relative coordinates. Again, we will treat the CoM part of the problem as “solved”.

Warning: always ask yourself -> is this an operator or a number? H, P and p are operators here.

# CM Separation Flowchart



# Spherical polar coordinates



Spherical Coordinates

$$\mathbf{r} = r \sin \theta \cos \phi \mathbf{i} + r \sin \theta \sin \phi \mathbf{j} + r \cos \theta \mathbf{k}$$

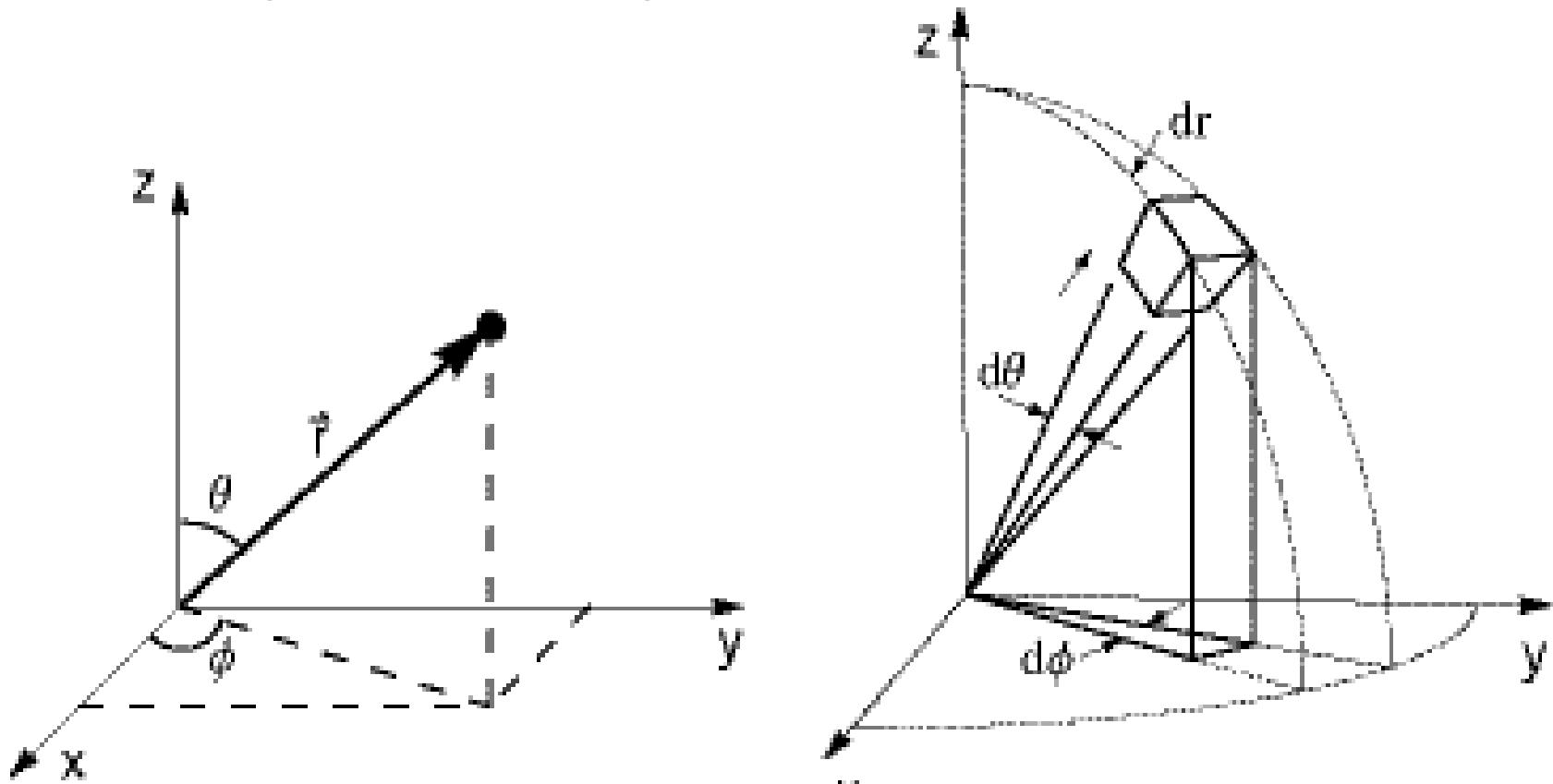
$$\hat{\mathbf{r}} = \sin \theta \cos \phi \mathbf{i} + \sin \theta \sin \phi \mathbf{j} + \cos \theta \mathbf{k}$$

$$\hat{\theta} = \cos \theta \cos \phi \mathbf{i} + \cos \theta \sin \phi \mathbf{j} - \sin \theta \mathbf{k}$$

$$\hat{\phi} = -\sin \phi \mathbf{i} + \cos \phi \mathbf{j}$$

$$\nabla = \hat{\mathbf{r}} \frac{\partial}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi}$$

# Spherical polar coordinates

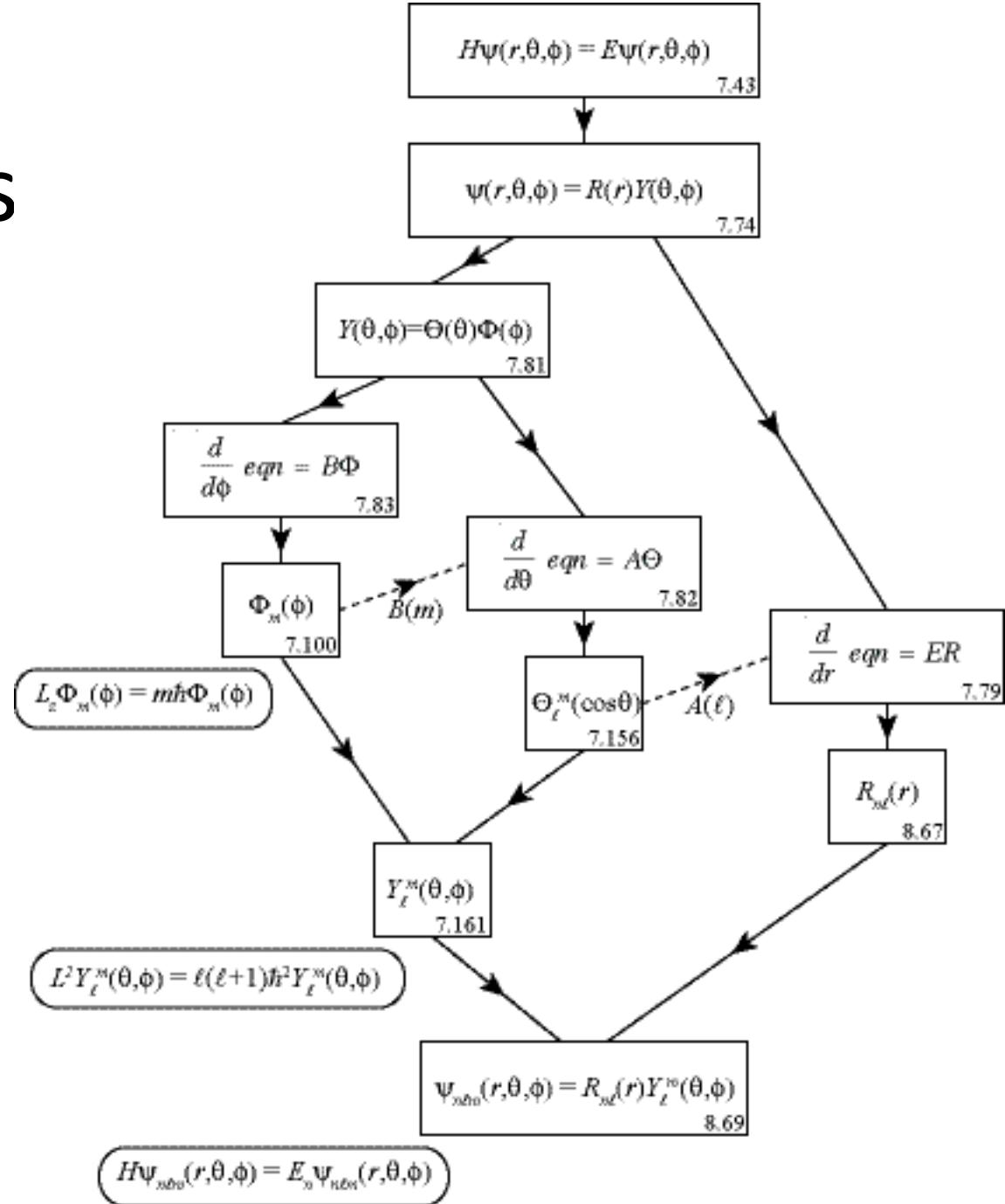


(a)

$$\begin{aligned}dV &= r^2 \sin\theta \, d\theta \, d\phi \, dr \\&= (r \, d\theta) (r \sin\theta \, d\phi) (dr) \\&= (\sin\theta \, d\theta) (d\phi) (r^2 \, dr) \\&= r^2 \, dr \, d\Omega\end{aligned}$$

(b)

# Separation of Variables Flowchart



# Spherical coordinates

- Energy eigenvalue equation for reduced mass part of the 2-body problem with kinetic energy operator explicitly in spherical coordinates

$$\left[ \underbrace{-\frac{\hbar^2}{2\mu} \nabla^2}_{\text{kinetic energy operator}} + \underbrace{V(r)}_{\text{potential energy operator}} \right] \underbrace{\psi(r, \theta, \phi)}_{\text{eigenfunction}} = \underbrace{E}_{\text{eigenvalue}} \underbrace{\psi(r, \theta, \phi)}_{\text{eigenfunction}}$$

$$-\frac{\hbar^2}{2\mu} \left[ \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \left( \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right] \psi$$

# Angular momentum

- This is angular momentum in rectangular coordinates

$$\mathbf{L} = \mathbf{r} \times \mathbf{p}$$

$$L_x = yp_z - zp_y \doteq -i\hbar \left( y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right)$$

$$L_y = zp_x - xp_z \doteq -i\hbar \left( z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right)$$

$$L_z = ?$$

$$\mathbf{L}^2 = \mathbf{L}_x^2 + \mathbf{L}_y^2 + \mathbf{L}_z^2$$

$$[L_x, L_y] = i\hbar L_z$$

$$[L_y, L_z] = i\hbar L_x$$

$$[L_z, L_x] = i\hbar L_y$$

$$[\mathbf{L}^2, L_{x,y,z}] = 0$$

$$\mathbf{L}^2 |\ell m_\ell\rangle = \ell(\ell+1)\hbar^2 |\ell m_\ell\rangle$$

$$L_z |\ell m_\ell\rangle = m_\ell \hbar |\ell m_\ell\rangle$$

BUT JUST LIKE SPIN!<sup>8</sup>

# Angular momentum

- This is angular momentum in spherical coordinates (homework)

$$L_x \doteq i\hbar \left( \sin\phi \frac{\partial}{\partial\theta} + \cos\phi \cot\theta \frac{\partial}{\partial\phi} \right)$$

$$L_y \doteq i\hbar \left( -\cos\phi \frac{\partial}{\partial\theta} + \sin\phi \cot\theta \frac{\partial}{\partial\phi} \right)$$

$$L_z \doteq -i\hbar \frac{\partial}{\partial\phi}$$

$$\mathbf{L}^2 \doteq -\hbar^2 \left[ \frac{1}{\sin\theta} \frac{\partial}{\partial\theta} \left( \sin\theta \frac{\partial}{\partial\theta} \right) + \frac{1}{\sin^2\theta} \frac{\partial^2}{\partial\phi^2} \right]$$

# Angular momentum

- Here comes the big simplification:

$$H_{KE} = -\frac{\hbar^2}{2\mu} \left[ \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]$$

$$-\frac{\hbar^2}{2\mu} \left[ \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) - \frac{1}{\hbar^2 r^2} \mathbf{L}^2 \right] \psi(r, \theta, \phi) + V(r) \psi(r, \theta, \phi) = E \psi(r, \theta, \phi)$$

- Assume a separable solution:

$$\psi(r, \theta, \phi) = R(r) Y(\theta, \phi)$$

# Separation of variables:

- Blue has angular dependence, red is radial:

$$-\frac{\hbar^2}{2\mu} \left[ \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) - \frac{1}{\hbar^2 r^2} \mathbf{L}^2 \right] R(r) Y(\theta, \phi) + V(r) R(r) Y(\theta, \phi)$$
$$= E R(r) Y(\theta, \phi)$$

Algebra (follow book 7.4):

$$\underbrace{\frac{1}{R(r)} \frac{d}{dr} \left( r^2 \frac{dR(r)}{dr} \right)}_{function \ of \ r \ only} - \frac{2\mu}{\hbar^2} (E - V(r)) r^2 = \underbrace{\frac{1}{\hbar^2} \frac{1}{Y(\theta, \phi)} \mathbf{L}^2 Y(\theta, \phi)}_{function \ of \ \theta, \phi \ only} \equiv A$$

# Separation of variables:

- Once we solve the (blue) angular problem, it is the solution to the angular part of ALL central force problems!

$$\frac{1}{\hbar^2} \mathbf{L}^2 Y(\theta, \phi) = A Y(\theta, \phi)$$

- Once we find  $A$  (and  $Y$ ), plug back into red equation and solve to find  $E$  (and  $R(r)$ ).

$$\frac{d}{dr} \left( r^2 \frac{dR(r)}{dr} \right) - \frac{2\mu}{\hbar^2} (E - V(r)) r^2 R(r) \equiv A R(r)$$

# Separate angular equation

- But we need to work on blue equation more, first.

$$\mathbf{L}^2 \doteq -\hbar^2 \left[ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]$$

$$Y(\theta, \phi) = \Theta(\theta) \Phi(\phi)$$

$$\left[ \frac{1}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{d}{d\theta} \right) - B \frac{1}{\sin^2 \theta} \right] \Theta(\theta) = -A \Theta(\theta)$$

$$\frac{d^2 \Phi(\phi)}{d\phi^2} = -B \Phi(\phi)$$

# Summary

- Here's the plan:

$$\frac{d^2\Phi(\phi)}{d\phi^2} = -B\Phi(\phi)$$

$$\left[ \frac{1}{\sin\theta} \frac{d}{d\theta} \left( \sin\theta \frac{d}{d\theta} \right) - B \frac{1}{\sin^2\theta} \right] \Theta(\theta) = -A\Theta(\theta)$$

$$\frac{d}{dr} \left( r^2 \frac{dR(r)}{dr} \right) - \frac{2\mu}{\hbar^2} (E - V(r)) r^2 R(r) \equiv AR(r)$$

$$\psi_{n\ell m}(r, \theta, \phi) = R_{n\ell}(r) \Theta_\ell^m(\theta) \Phi_m(\phi)$$

$$\psi(r, \theta, \phi, t) = \sum_{n\ell m} c_{n\ell m} R_{n\ell}(r) \Theta_\ell^m(\theta) \Phi_m(\phi) e^{-iE_n t/\hbar}$$

# Summary

- Here's the plan:
- We'll consider 3 different systems, a ring (to solve the  $\phi$  problem), a sphere (to solve the  $\theta$  and  $\phi$  problem), and the full hydrogen atom (to solve the  $(r, \theta, \phi)$  problem)
- We'll find the quantum numbers and wave functions that solve each problem
- We'll apply all the things you've learned in PH424 and PH425
- **Please read ahead** – the math is much more intense (though not harder) than before