

$$dE_x = \frac{1}{4\pi\epsilon_0} \frac{dQ}{(d+L-x)^2} = \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \frac{dx}{(d+L-x)^2}$$

For the portion of the rod from $x=0$ to $x=L/2$, E points in the negative x direction.

$$\vec{E}_x = \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \left[-\int_0^{L/2} \frac{dx}{(d+L-x)^2} + \int_{L/2}^L \frac{dx}{(d+L-x)^2} \right]$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \left[-\frac{1}{d+L-x} \Big|_0^{L/2} + \frac{1}{d+L-x} \Big|_{L/2}^L \right]$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \left[-\frac{2}{d+L/2} + \frac{1}{d} + \frac{1}{d+L} \right]$$

(b) At large d the rod should look like a dipole with $E_x \propto 1/d^3$.

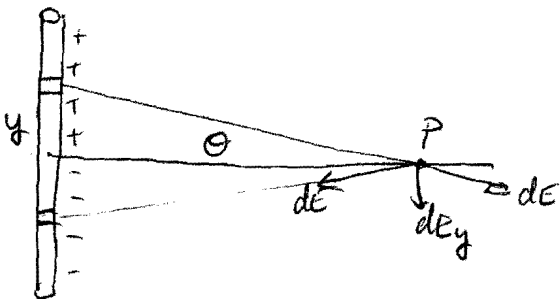
$$E_x \approx \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \frac{1}{d} \left[1 + \frac{1}{1+L/d} - \frac{2}{1+L/2d} \right]$$

$$\approx \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \frac{1}{d} \left[1 + \left(1 - \frac{L}{d} + \left(\frac{L}{d}\right)^2\right) - 2\left(1 - \frac{L}{2d} + \frac{L^2}{4d^2}\right) \right]$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \frac{1}{d} \left[\frac{L^2}{d^2} - \frac{L^2}{2d^2} \right] = \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \frac{L^2}{2d^3}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q \cdot L/2}{d^3}$$

②



(a) For each element dq in the top half with + charge, there is a corresponding element in the bottom half with - charge, such that the x components of the fields cancel and only the y component survives.

$$dE_y = dE \sin \theta = \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \frac{dy}{d^2+y^2} \frac{y}{\sqrt{d^2+y^2}}$$

Integrate from 0 to $L/2$ and include factor of 2 to account for lower half:

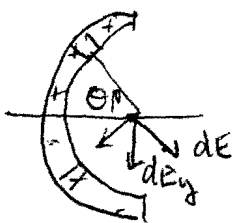
$$\begin{aligned} E_y &= \frac{2}{4\pi\epsilon_0} \frac{Q}{L} \int_0^{L/2} \frac{y dy}{(d^2+y^2)^{3/2}} = \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \int_0^{L/2} \frac{2y dy}{(d^2+y^2)^{3/2}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{Q}{L} \left[-2(d^2+y^2)^{-1/2} \right]_0^{L/2} \\ &= \frac{1}{2\pi\epsilon_0} \frac{Q}{L} \left(-\frac{1}{d} + \frac{1}{\sqrt{d^2+L^2/4}} \right) \end{aligned}$$

(b) For $d \gg L$

$$E_y \approx \frac{1}{2\pi\epsilon_0} \frac{Q}{L} \left(-\frac{1}{d} + \frac{1}{d} \left(1 - \frac{1}{2} \frac{L^2}{4d^2} \right) \right) = \frac{-1}{4\pi\epsilon_0} \frac{Q/4L}{d^3}$$

The field depends on $1/d^3$, like a dipole.

(3)



$$dQ = \frac{Q}{\pi R} R d\theta = \frac{Q}{\pi} d\theta$$

For every $+dQ$ on the upper half, there will be a $-dQ$ on the lower half such that the x components of the field will cancel.

$$dE_y = dE \sin \theta = \frac{1}{4\pi\epsilon_0} \frac{dQ}{R^2} \sin \theta = \frac{1}{4\pi\epsilon_0} \frac{Q}{\pi R^2} \sin \theta d\theta$$

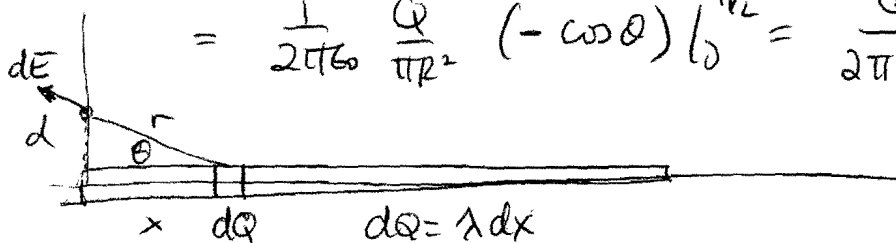
Integrate from $\theta = 0$ to $\pi/2$ and double the result to account for the lower half:

$$E_y = 2 \int_0^{\pi/2} \frac{1}{4\pi\epsilon_0} \frac{Q}{\pi R^2} \sin \theta d\theta$$

$$= \frac{1}{2\pi\epsilon_0} \frac{Q}{\pi R^2} \int_0^{\pi/2} \sin \theta d\theta$$

$$= \frac{1}{2\pi\epsilon_0} \frac{Q}{\pi R^2} (-\cos \theta) \Big|_0^{\pi/2} = \frac{Q}{2\pi^2\epsilon_0 R^2}$$

(4)



$$dE = \frac{dQ}{4\pi\epsilon_0 r^2} = \frac{\lambda}{4\pi\epsilon_0} \frac{dx}{x^2 + d^2}$$

$$dE_x = dE \cos \theta = \frac{\lambda}{4\pi\epsilon_0} \frac{dx}{x^2 + d^2} \frac{x}{\sqrt{x^2 + d^2}}$$

$$E_x = \frac{\lambda}{4\pi\epsilon_0} \int_0^{\infty} \frac{x dx}{(x^2 + d^2)^{3/2}} = \frac{\lambda}{4\pi\epsilon_0} \frac{1}{2} \int_0^{\infty} \frac{2x dx}{(x^2 + d^2)^{3/2}}$$

$$= \left(\frac{\lambda}{4\pi\epsilon_0} \frac{1}{2} \right) \left(-2(x^2 + d^2)^{-1/2} \Big|_0^{\infty} \right) = \frac{\lambda}{4\pi\epsilon_0 d} \text{ in } -x \text{ direction}$$

$$dE_y = dE \sin \theta = \frac{\lambda}{4\pi\epsilon_0} \frac{dx}{x^2+d^2} \frac{d}{\sqrt{x^2+d^2}}$$

$$E_y = \frac{\lambda d}{4\pi\epsilon_0} \int_0^\infty \frac{dx}{(x^2+d^2)^{3/2}}$$

$$= \frac{\lambda d}{4\pi\epsilon_0} \left. \frac{x}{d^2 \sqrt{x^2+d^2}} \right|_0^\infty$$

$$= \frac{\lambda}{4\pi\epsilon_0 d} \quad \text{in } +y \text{ direction}$$