

Integral Equations in Quantum Mechanics I

I Bound States, II Scattering*

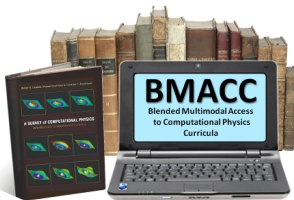
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Based on *A Survey of Computational Physics* by Landau, Páez, & Bordeianu

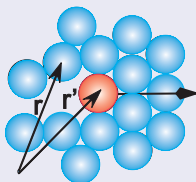
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Course: **Computational Physics II**



Problem: Bound States in Momentum Space

Integro-Differential Equation



- N-body interaction reduces to **nonlocal** $V_{\text{eff}}(r)$:

$$-\frac{1}{2m} \frac{d^2\psi(r)}{dr^2} + \int dr' V(r, r')\psi(r') = E\psi(r) \quad (1)$$

- Integro-differential equation
- **Problem:** Solve for $l = 0$ bound-state E_i & ψ_i

Theory: Momentum-Space Schrödinger Equation

Integral Schrödinger Equation Equally Valid

- Transform Schrödinger Equation to **momentum space**
- Replace integro-differential by integral equation:

$$\frac{\hbar^2 k^2}{2\mu} \psi(k) + \frac{2}{\pi} \int_0^\infty dp p^2 V(k, p) \psi(p) = E \psi(k) \quad (1)$$

- $V(k, p) =$ p-space representation (TF) of V :

$$V(k, p) = \frac{1}{kp} \int_0^\infty dr dr' \sin(kr) V(r, r') \sin(pr') \quad (2)$$

- $\psi(k) =$ p-space representation (TF) of ψ :

$$\psi(k) = \int_0^\infty dr kr \psi(r) \sin(kr) \quad (3)$$

- Will transform into matrix equation (see matrix Chapter)

Algorithm: Integral Equations \rightarrow Linear Equations

Solve on p-Space Grid

k_1 k_2 k_3 \dots k_N

- Integral \simeq weighted sum (see Integration chapter)

$$\int_0^\infty dp p^2 V(k, p) \psi(p) \simeq \sum_{j=1}^N w_j k_j^2 V(k, k_j) \psi(k_j) \quad (1)$$

- Integral equation \rightarrow algebraic equation

$$\frac{k^2}{2\mu} \psi(k) + \frac{2}{\pi} \sum_{j=1}^N w_j k_j^2 V(k, k_j) \psi(k_j) = E \quad (2)$$

- N unknowns $\psi(k_j)$
- 1 unknown E
- Unknown function $\psi(k)$
- Solve on grid, $k = k_j$
- $\rightarrow N$ couple equations
- $(N + 1)$ unknowns:

Algorithm: Integral Equations \rightarrow Linear Equations

Solve on p-Space Grid

k_1 k_2 k_3 \dots k_N

$$\frac{k_i^2}{2\mu} \psi(k_i) + \frac{2}{\pi} \sum_{j=1}^N w_j k_j^2 V(k_i, k_j) \psi(k_j) = E \psi(k_i), \quad i = 1, N \quad (1)$$

- e.g. $N = 2 \Rightarrow 2$ coupled linear equations

$$\frac{k_1^2}{2\mu} \psi(k_1) + \frac{2}{\pi} w_1 k_1^2 V(k_1, k_1) \psi(k_1) + w_2 k_2^2 V(k_1, k_2) \psi(k_2) = E \psi(k_1) \quad (2)$$

$$\frac{k_2^2}{2\mu} \psi(k_2) + \frac{2}{\pi} w_1 k_1^2 V(k_2, k_1) \psi(k_1) + w_2 k_2^2 V(k_2, k_2) \psi(k_2) = E \psi(k_2) \quad (3)$$

Algorithm: Integral Equations \rightarrow Linear Equations

Solve on p-Space Grid

k_1 k_2 k_3 \dots k_N

$$\frac{k_i^2}{2\mu}\psi(k_i) + \frac{2}{\pi} \sum_{j=1}^N w_j k_j^2 V(k_i, k_j)\psi(k_j) = E\psi(k_i), \quad i = 1, N \quad (1)$$

- Matrix Schrödinger equation $[H][\psi] = E[\psi]$
- $\psi(k) = N \times 1$ vector

$$\begin{pmatrix} \frac{k_1^2}{2\mu} + \frac{2}{\pi} V(k_1, k_1)k_1^2 w_1 & \frac{2}{\pi} V(k_1, k_2)k_2^2 w_2 & \dots & \frac{2}{\pi} V(k_1, k_N)k_N^2 w_N \\ \dots & \dots & \dots & \frac{k_N^2}{2\mu} + \frac{2}{\pi} V(k_N, k_N)k_N^2 w_N \end{pmatrix} \times \begin{pmatrix} \psi(k_1) \\ \psi(k_2) \\ \vdots \\ \psi(k_N) \end{pmatrix} = E \begin{pmatrix} \psi(k_1) \\ \psi(k_2) \\ \vdots \\ \psi(k_N) \end{pmatrix} \quad (2)$$

Eigenvalue Problem

Search for Solution; N equations for $(N + 1)$ unknowns?

- Solution only sometimes, certain E (eigenvalues)
- Try to solve, multiply both sides by $[H - E\mathbb{I}]$ inverse:

$$[H][\psi] = E[\psi] \quad (1)$$

$$[H - E\mathbb{I}][\psi] = [0] \quad (2)$$

$$\Rightarrow [\psi] = [H - E\mathbb{I}]^{-1}[0] \quad (3)$$

- \Rightarrow if inverse \exists , then only **trivial** solution $\psi \equiv 0$
- For nontrivial solution inverse can't \exists

$$\det[H - E\mathbb{I}] = 0 \quad (\text{bound-state condition}) \quad (4)$$

- Requisite additional equation for $N + 1$ unknowns
- Solve for just eigenvalues, or full e.v. problem

Model: Delta-Shell Potential (Sort of Analytic Solution)

2 Particles Interact When b Apart

$$V(r) = \frac{\lambda}{2\mu} \delta(r - b) \quad (1)$$

$$V(k', k) = \frac{1}{k'k} \int_0^\infty \sin(k'r') \frac{\lambda}{2\mu} \delta(r - b) \sin(kr) dr \quad (2)$$

$$= \frac{\lambda}{2\mu} \frac{\sin(k'b) \sin(kb)}{k'k} \quad (\text{too slow decay}) \quad (3)$$

- 1 Bound state $E = -\kappa^2/2\mu$, if

$$e^{-2\kappa b} - 1 = \frac{2\kappa}{\lambda} \quad (4)$$

- Only if strong & attractive ($\lambda < 0$)
- **Exercise:** Solve transcendental equation ($b = 10$, $\lambda = ?$)

Bound-State Integral-Equation Code



CODE

Sample Code Surveys all Parameters

- Gauss quadrature for pts & wts
- Two possible libe calls
- 1. Search on E , $\det[H - EI] = 0$
- 2. Use eigenproblem solver*
- Both iterative solutions

Your Implementation

Modify or Write Eigenvalues, Eigenproblem

- 1 $2\mu = 1, b = 10, N > 16$
- 2 Set up $[V(i, j)]$ and $[H(i, j)]$ for $N \geq 16$
- 3 Observe monotonic relation $E(\lambda)$
- 4 True bound state stable with N , others = artifacts
- 5 Extract best value for E & estimate precision
- 6 Comparing RHS, LHS $[H][\psi] = E[\psi]$

Exploration: Momentum Space Wave Function*

Bound in p Space?

- 1 Determine $\psi(k)$ (analytic $\psi(p) \propto [p^2 - 2mE]^{-1}$)
- 2 Is this reasonable, normalizable?
- 3 Determine $\psi(r)$ via transform

$$\psi(r) = \int_0^{\infty} dk \psi(k) \frac{\sin(kr)}{kr} k^2 \quad (1)$$

- 4 Is this reasonable $\psi(r)$?
- 5 Compare to analytic $\psi(r)$,

$$\psi_0(r) \propto \begin{cases} e^{-\kappa r} - e^{\kappa r}, & \text{for } r < b, \\ e^{-\kappa r}, & \text{for } r > b \end{cases} \quad (2)$$