# Discrete and continuous quantum walks

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(based on joint work with R.Burton, Z.Dimcovic and T.Nguyen)

Stochastic processes:  $\ell^1(\mathbb{R})$  norm preserving linear evolution

$$\frac{d}{dt}\mu_t = \mu_t Q$$

For  $\mu_t = (\mu_1(t), \mu_2(t), \dots)$  being  $\ell^1(\mathbb{R})$  norm preserving means

$$\mu_1(t) + \mu_2(t) + \cdots = 1$$

at all times.

Quantum evolution:  $\ell^2(\mathbb{C})$  norm preserving linear evolution

$$\frac{d}{dt}\psi_t = -iH\psi_t$$
,  $H$  is self-adjoint

For  $\psi_t = (\psi_1(t), \psi_2(t), \dots)$  being  $\ell^2(\mathbb{C})$  norm preserving means

$$|\psi_1(t)|^2 + |\psi_2(t)|^2 + \dots = 1$$

at all times.

Dirac notations:  $\frac{d}{dt}|\psi_t>=-iH|\psi_t>$ 

Shrödinger Eq.

$$\frac{d}{dt}\psi_t = -iH\psi_t, \quad H \text{ is self-adjoint}$$

Dirac notations:  $\frac{d}{dt}|\psi_t>=-iH|\psi_t>$ 

Hamiltonian operator H: Eigenvalues must be real  $\lambda_j \in \mathbb{R}$ , and the eigenvectors  $v_j$  are orthonormal.

Operator  $U_t=e^{-itH}$  will have eigenvectors  $e^{-it\lambda_j}$  of unit magnitude, and the same orthonormal eigenvectors  $v_j$ 

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Take 
$$\psi=\sum_j a_j v_j$$
 s.t.  $\sum_j |a_j|^2=1$ , then 
$$U_t\psi=\sum_j a_j e^{-it\lambda_j}v_j,$$

where  $\sum_{i} |a_{i}e^{-it\lambda_{j}}|^{2} = 1$ 

Dirac notation:  $|\psi>=\sum_{j}a_{j}|v_{j}>$ , then  $U_{t}|\psi>=\sum_{j}a_{j}e^{-it\lambda_{j}}|v_{j}>$ 

## Classical randomized algorithms

Randomized algorithms is an effective tool for speeding up computations and is an important field for applications of stochastic processes, e.g. Markov chain Monte Carlo (MCMC).

In short: classical computation makes use of the  $\ell^1(\mathbb{R})$  norm preserving linear Markov evolution  $\frac{d}{dt}\mu_t=\mu_tQ$ 

## Randomized algorithms

In short: classical computation makes use of the  $\ell^1(\mathbb{R})$  norm preserving linear Markov evolution  $\frac{d}{dt}\mu_t=\mu_tQ$ 

Quantum computation: analogous tool is being developed, called the **quantum walk**. Idea: make use of the  $\ell^2(\mathbb{C})$  norm preserving linear Shrödinger evolution  $\frac{d}{dt}|\psi_t>=-iH|\psi_t>$ 

Both the classical and quantum computers provide the framework for implementation.

## Quantum computation: qubits

One qubit system: two basis vectors |0> and |1>

Two qubit system: four basis vectors |00>, |01>, |10> and |11>

Another notation: |0>, |1>, |2> and |3>

Tensor notation:  $|0>\otimes|0>$ ,  $|0>\otimes|1>$ ,  $|1>\otimes|0>$  and  $|1>\otimes|1>$ 

### **Quantum walk**

Hilbert space  $\mathcal{H}_C = \{|\downarrow\rangle, |\uparrow\rangle\}$  represents the outcome of a "coin toss"

Hilbert space  $\mathcal{H}_P$  represents the position of the walker

Distribution

$$|\psi\rangle = \sum_{j} a_{j} |\uparrow\rangle \otimes |j\rangle + |b_{j}| \downarrow\rangle \otimes |j\rangle$$

means the walker is at site j with probability

$$|a_i|^2 + |b_i|^2$$

#### Quantum walk

Hilbert space:  $\mathcal{H}_C \otimes \mathcal{H}_P$ 

Shrödinger Eq. 
$$\frac{d}{dt}|\psi_t>=-iH|\psi_t>$$

Discrete time: 
$$U = e^{-iH}$$
,  $|\psi_{t+1}\rangle = U|\psi_t\rangle$ 

Quantum evolution:

$$|\psi_t> = \sum_j a_j(t)|\uparrow>\otimes|j> +b_j(t)|\downarrow>\otimes|j>$$

means the walker is at site j with probability

$$|a_j(t)|^2 + |b_j(t)|^2$$

## Hadamard quantum walk

Conside the following (Hdamard) coin on the two qubit space  $\mathcal{H}_C = \{|\downarrow\rangle, |\uparrow\rangle\}$ 

$$C = \frac{1}{\sqrt{2}} \left( \begin{array}{cc} 1 & 1 \\ 1 & -1 \end{array} \right)$$

Let the transition matrix for the Hadamard quantum walk be the following operator on  $\mathcal{H}_C \otimes \mathcal{H}_P$ 

$$U = S(C \otimes I),$$

where

$$S = |\uparrow\rangle < \uparrow|\otimes \sum_{s} |s+1\rangle < s| + |\downarrow\rangle < \downarrow|\otimes \sum_{s} |s-1\rangle < s|$$

## **Example: Hadamard quantum walk**

Take  $|\psi_0>=|\downarrow>\otimes|0>$ . First iteration:

$$(C \otimes I)|\psi_0> = C|\downarrow>\otimes I|0>$$

$$=\frac{1}{\sqrt{2}}|\uparrow>\otimes|0>-\frac{1}{\sqrt{2}}|\downarrow>\otimes|0>$$

Now,

$$S = |\uparrow\rangle < \uparrow|\otimes \sum_{s} |s+1\rangle < s| + |\downarrow\rangle < \downarrow|\otimes \sum_{s} |s-1\rangle < s|$$

and  $|\psi_1>=U|\psi_0>=S(C\otimes I)|\psi_0>$ 

$$=rac{1}{\sqrt{2}}|\uparrow>\otimes|1>-rac{1}{\sqrt{2}}|\downarrow>\otimes|-1>$$

$$|\psi_0>=|\downarrow>\otimes|0>$$

$$|\psi_1>=\frac{1}{\sqrt{2}}|\uparrow>\otimes|1>-\frac{1}{\sqrt{2}}|\downarrow>\otimes|-1>$$

Next iteration:

$$(C\otimes I)|\psi_1>=\frac{1}{\sqrt{2}}C|\uparrow>\otimes|1>-\frac{1}{\sqrt{2}}C|\downarrow>\otimes|-1>$$

$$= \frac{1}{2} |\uparrow\rangle \otimes |1\rangle + \frac{1}{2} |\downarrow\rangle \otimes |1\rangle - \frac{1}{2} |\uparrow\rangle \otimes |-1\rangle$$

$$+ \frac{1}{2} |\downarrow\rangle \otimes |-1\rangle \quad \text{Now}$$

$$+\frac{1}{2}|\downarrow>\otimes|-1>$$
 Now,

$$S = |\uparrow\rangle < \uparrow| \otimes \sum_{s} |s+1\rangle < s| + |\downarrow\rangle < \downarrow| \otimes \sum_{s} |s-1\rangle < s|$$
 and  $|\psi_2\rangle = \frac{1}{2}|\uparrow\rangle \otimes |2\rangle + \frac{1}{2}|\downarrow\rangle \otimes |0\rangle + \frac{1}{2}|\uparrow\rangle \otimes |0\rangle + \frac{1}{2}|\downarrow\rangle \otimes |-2\rangle$ 

Now 
$$|\psi_2\rangle = \frac{1}{2}|\uparrow\rangle \otimes |2\rangle + \frac{1}{2}|\downarrow\rangle \otimes |0\rangle - \frac{1}{2}|\uparrow\rangle \otimes |0\rangle + \frac{1}{2}|\downarrow\rangle \otimes |-2\rangle$$

$$(C \otimes I)|\psi_2\rangle = \frac{1}{2}C|\uparrow\rangle \otimes |2\rangle + \frac{1}{2}C|\downarrow\rangle \otimes |0\rangle$$

$$-\frac{1}{2}|\uparrow\rangle\otimes|0\rangle+\frac{1}{2}|\downarrow\rangle\otimes|-2\rangle$$

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$$(C\otimes I)|\psi_2\rangle = \frac{1}{2}C|\uparrow\rangle\otimes|2\rangle + \frac{1}{2}C|\downarrow\rangle\otimes|0\rangle$$

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$$(C \otimes I)|\psi_{2}\rangle = \frac{1}{2}C|\uparrow\rangle\otimes|2\rangle + \frac{1}{2}C|\downarrow\rangle\otimes|0\rangle - \frac{1}{2}C|\uparrow\rangle\otimes|0\rangle + \frac{1}{2}C|\downarrow\rangle\otimes|-2\rangle = \frac{1}{2\sqrt{2}}|\uparrow\rangle\otimes|2\rangle + \frac{1}{2\sqrt{2}}|\downarrow\rangle\otimes|2\rangle - \frac{2}{2}|\downarrow\rangle\otimes|0\rangle + \frac{1}{2}|\downarrow\rangle\otimes|-2\rangle$$

$$= \frac{1}{2\sqrt{2}} |\uparrow\rangle \otimes |2\rangle + \frac{1}{2\sqrt{2}} |\downarrow\rangle \otimes |2\rangle$$

$$- \frac{2}{2\sqrt{2}} |\downarrow\rangle \otimes |0\rangle + \frac{1}{2\sqrt{2}} |\uparrow\rangle \otimes |-2\rangle$$

$$- \frac{1}{2\sqrt{2}} |\downarrow\rangle \otimes |-2\rangle$$

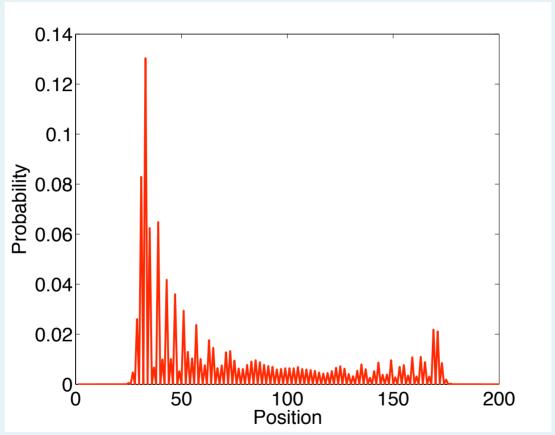
$$-\frac{1}{2\sqrt{2}}|\downarrow>\otimes|-2>$$

$$|\psi_{2}\rangle = \frac{1}{2\sqrt{2}}|\uparrow>\otimes|3> + \frac{1}{2\sqrt{2}}|\downarrow>\otimes|1>$$

$$-\frac{2}{2\sqrt{2}}|\downarrow>\otimes|-1> + \frac{1}{2\sqrt{2}}|\uparrow>\otimes|-1>$$

$$-\frac{1}{2\sqrt{2}}|\downarrow>\otimes|-3>$$

# Quantum Walk with Hadamard coin



### Markov chain with internal states

$$U(\pm |\uparrow>\otimes|s>)$$

$$=\frac{1}{\sqrt{2}}(\pm|\uparrow>\otimes|s+1>)+\frac{1}{\sqrt{2}}(\pm|\downarrow>\otimes|s-1>)$$

and

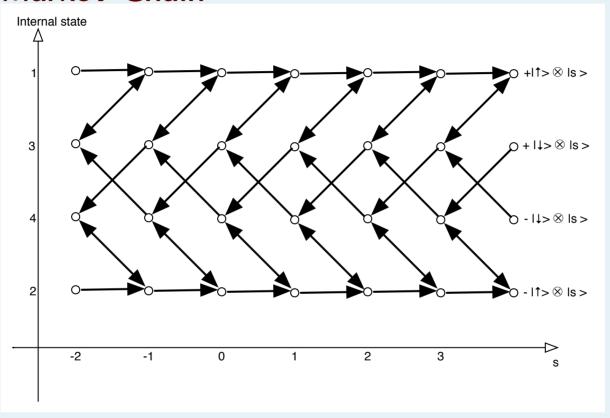
$$U(\pm |\downarrow> \otimes |s>)$$

$$=\frac{1}{\sqrt{2}}(\pm|\uparrow>\otimes|s+1>)+\frac{1}{\sqrt{2}}(\mp|\downarrow>\otimes|s-1>)$$

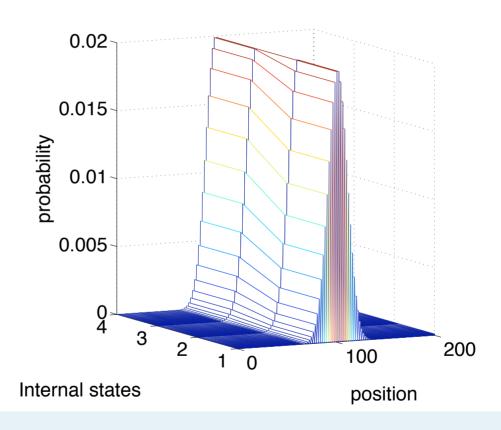
Four internal states:

$$(|\uparrow>\otimes|s>), \quad -(|\uparrow>\otimes|s>), \quad (|\downarrow>\otimes|s>)$$
 and  $-(|\downarrow>\otimes|s>)$ 

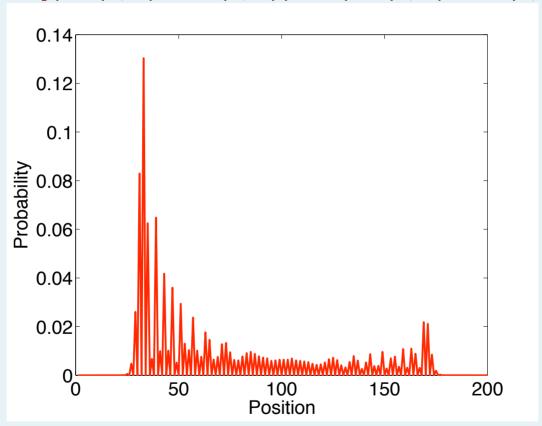
# **Markov Chain**



# Markov chain



$$2^{n}[(P_{n}(s,1)-P_{n}(s,2))^{2}+(P_{n}(s,3)-P_{n}(s,4))^{2}]$$



Four eigenvalues in the Fourier space:

$$\lambda_1=0, \quad \lambda_2=2\cos(k)$$
 and 
$$\lambda_{3,4}=\pm\sqrt{1+\cos^2(k)}+i\sin(k)$$

The distribution of the Hadamard quantum walk is expressed in the closed form as

$$\mu_n(s) = \frac{D_{\uparrow,n}^2(s) + D_{\downarrow,n}^2(s)}{2^n},$$

where

$$D_{\uparrow,n}(s) = \frac{1}{2\pi} \int_{-\pi}^{\pi} \frac{e^{-i(s+1)k}}{2\sqrt{1+\cos^2(k)}} \left[\lambda_3^n - \lambda_4^n\right] dk$$

and  $D_{\downarrow,n}(s)$ 

$$= \frac{1}{2\pi} \int_{-\pi}^{\pi} \frac{e^{-isk} \left[ (\lambda_4^n - \lambda_3^n) \cos(k) + (\lambda_3^n + \lambda_4^n) \sqrt{1 + \cos^2(k)} \right]}{2\sqrt{1 + \cos^2(k)}} dk$$