

# Research Projects in Math Modeling and Numerical Analysis

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Department of Mathematics



Graduate Student Seminar  
May 14, 2014

- 1 **Electromagnetics**
  - Maxwell's Equations
  - Numerical Analysis
  - Dispersive Media
  - Inverse Problems

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- River system and modeling equations
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- Model
- Solution Method
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# Acknowledgements

## Collaborators

- H. T. Banks (NCSU)
- V. A. Bokil (OSU)
- W. P. Winfree (NASA)

## Students

- Karen Barrese and Neel Chugh (REU 2008)
- Anne Marie Milne and Danielle Wedde (REU 2009)
- Erin Bela and Erik Hortsch (REU 2010)
- Megan Armentrout (MS 2011)
- Brian McKenzie (MS 2011)

# Maxwell's Equations

$$\frac{\partial \mathbf{D}}{\partial t} + \mathbf{J} = \nabla \times \mathbf{H} \quad (\text{Ampere})$$

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E} \quad (\text{Faraday})$$

$$\nabla \cdot \mathbf{D} = \rho \quad (\text{Poisson})$$

$$\nabla \cdot \mathbf{B} = 0 \quad (\text{Gauss})$$

$\mathbf{E}$  = Electric field vector

$\mathbf{D}$  = Electric flux density

$\mathbf{H}$  = Magnetic field vector

$\mathbf{B}$  = Magnetic flux density

$\rho$  = Electric charge density

$\mathbf{J}$  = Current density

With appropriate initial conditions and boundary conditions.

## Constitutive Laws

Maxwell's equations are completed by constitutive laws that describe the response of the medium to the electromagnetic field.

$$\mathbf{D} = \epsilon \mathbf{E} + \mathbf{P}$$

$$\mathbf{B} = \mu \mathbf{H} + \mathbf{M}$$

$$\mathbf{J} = \sigma \mathbf{E} + \mathbf{J}_s$$

$\mathbf{P}$  = Polarization                       $\epsilon$  = Electric permittivity

$\mathbf{M}$  = Magnetization                     $\mu$  = Magnetic permeability

$\mathbf{J}_s$  = Source Current                     $\sigma$  = Electric Conductivity

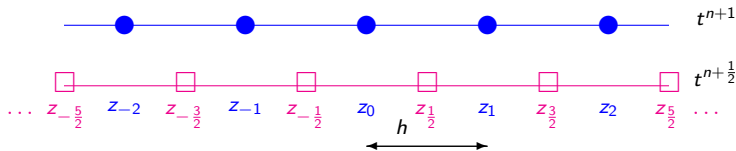


# Yee Scheme in One Space Dimension

- **Staggered Grids:** The electric field/flux is evaluated on the primary grid in both space and time and the magnetic field/flux is evaluated on the dual grid in space and time.
- The Yee scheme is

$$\frac{H|_{\ell+\frac{1}{2}}^{n+\frac{1}{2}} - H|_{\ell+\frac{1}{2}}^{n-\frac{1}{2}}}{\Delta t} = -\frac{1}{\mu} \frac{E|_{\ell+1}^n - E|_{\ell}^n}{\Delta z}$$

$$\frac{E|_{\ell}^{n+1} - E|_{\ell}^n}{\Delta t} = -\frac{1}{\epsilon} \frac{H|_{\ell+\frac{1}{2}}^{n+\frac{1}{2}} - H|_{\ell-\frac{1}{2}}^{n+\frac{1}{2}}}{\Delta z}$$



## Yee Scheme in One Space Dimension

- This gives an explicit second order accurate scheme in both time and space.

## Yee Scheme in One Space Dimension

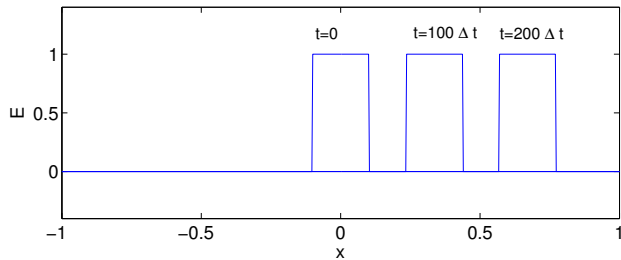
- This gives an explicit second order accurate scheme in both time and space.
- It is conditionally stable with the CFL condition

$$\nu = \frac{c\Delta t}{\Delta z} \leq 1$$

where  $\nu$  is called the Courant number and  $c = 1/\sqrt{\epsilon\mu}$ .

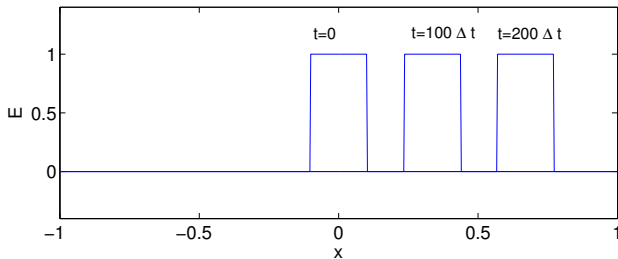
# Numerical Stability: A Square Wave

- Case  $c\Delta t = \Delta z$

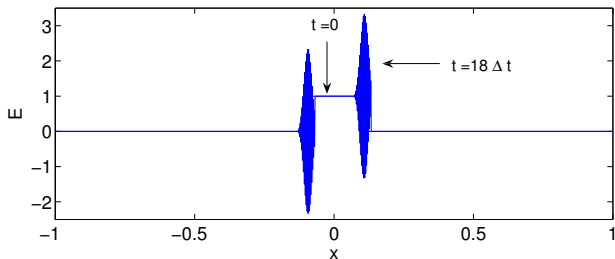


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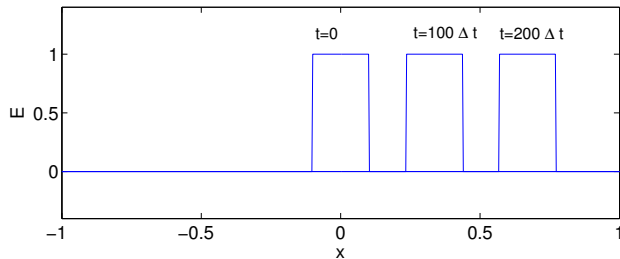


- Case  $c\Delta t > \Delta z$

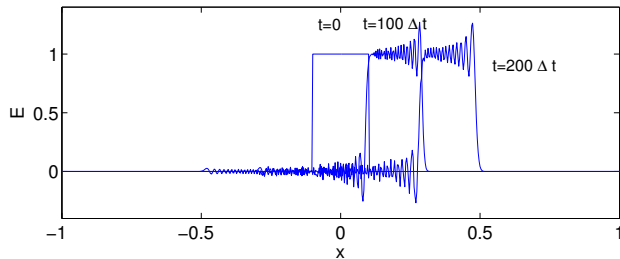


# Numerical Dispersion: A Square Wave

- Case  $c\Delta t = \Delta z$



- Case  $c\Delta t < \Delta z$



## Dispersion Error

- The Yee scheme can exhibit **numerical dispersion**

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### Theorem (CFL for $(2, 2M)$ [Bokil-G, 2011] )

*The method is conditionally stable with CFL condition*

$$\nu \left( \sum_{p=1}^M \gamma_{2p-1} \right) < 1,$$

$$\begin{aligned} \gamma_{2p-1} &:= \sum_{j=p}^M \frac{(-1)^{3j-p-2} (j+p-2)! [(2M-1)!!]^2 2^{2p-1}}{(2p-1)! (j-p)! (2j-1) (2M-2j)! (2M+2j-2)!!} \\ &= \frac{[(2p-3)!!]^2}{(2p-1)!}. \end{aligned}$$

## Dispersion Error (cont.)

- Mimetic methods adapt free parameters in the scheme to reduce certain errors, e.g., dispersion error [Bokil-G-Gyrya-McGregor, submitted 2014].

## Complex permittivity

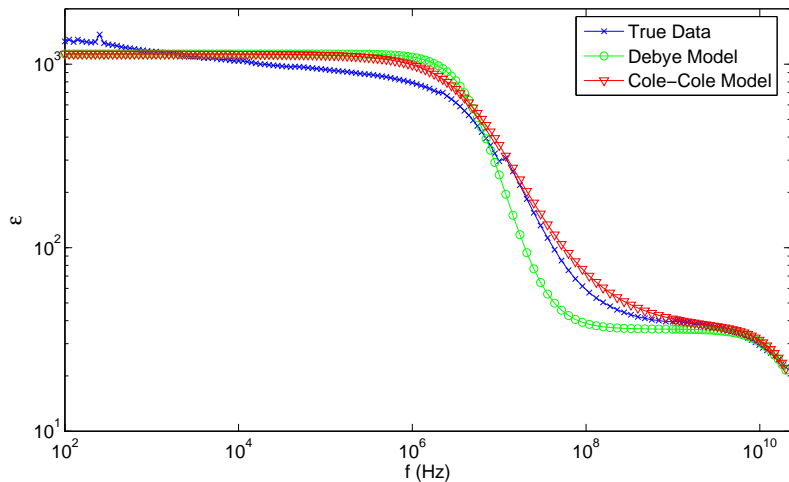
- We can usually define  $\mathbf{P}$  in terms of a convolution

$$\mathbf{P}(t, \mathbf{x}) = g * \mathbf{E}(t, \mathbf{x}) = \int_0^t g(t-s, \mathbf{x}; \mathbf{q}) \mathbf{E}(s, \mathbf{x}) ds,$$

where  $g$  is the dielectric response function (DRF).

- In the frequency domain  $\hat{\mathbf{D}} = \epsilon \hat{\mathbf{E}} + \hat{\mathbf{g}} \hat{\mathbf{E}} = \epsilon_0 \epsilon(\omega) \hat{\mathbf{E}}$ , where  $\epsilon(\omega)$  is called the **complex permittivity**.
- $\epsilon(\omega)$  described by the polarization model
- We are interested in ultra-wide bandwidth electromagnetic pulse interrogation of dispersive dielectrics, therefore we want an accurate representation of  $\epsilon(\omega)$  over a broad range of frequencies.

## Dry skin data



**Figure:** Real part of  $\epsilon(\omega)$ ,  $\epsilon$ , or the permittivity [GLG96].

$$\mathbf{P}(t, \mathbf{x}) = g * \mathbf{E}(t, \mathbf{x}) = \int_0^t g(t-s, \mathbf{x}; \mathbf{q}) \mathbf{E}(s, \mathbf{x}) ds,$$

- Debye model [1929]  $\mathbf{q} = [\epsilon_d, \tau]$

$$g(t, \mathbf{x}) = \epsilon_0 \epsilon_d / \tau e^{-t/\tau}$$

$$\text{or } \tau \dot{\mathbf{P}} + \mathbf{P} = \epsilon_0 \epsilon_d \mathbf{E}$$

$$\text{or } \epsilon(\omega) = \epsilon_\infty + \frac{\epsilon_d}{1 + i\omega\tau}$$

with  $\epsilon_d := \epsilon_s - \epsilon_\infty$  and  $\tau$  a relaxation time.

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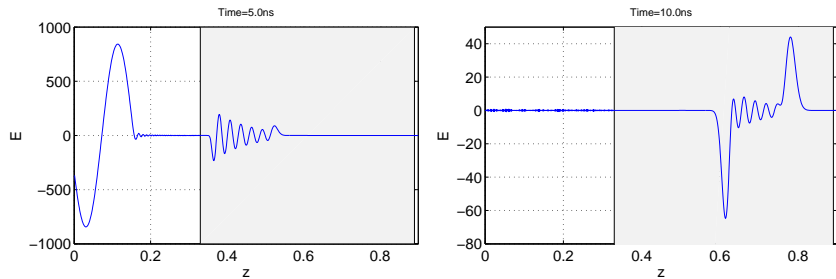
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- Cole-Cole model [1936] (heuristic generalization)

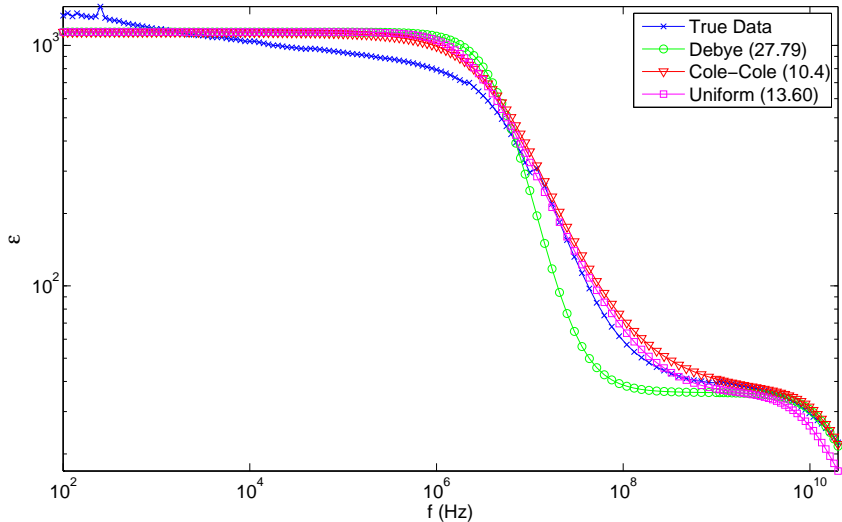
$$\mathbf{q} = [\epsilon_d, \tau, \alpha]$$

$$\epsilon(\omega) = \epsilon_\infty + \frac{\epsilon_d}{1 + (i\omega\tau)^{1-\alpha}}$$

# Dispersive Media



**Figure:** Debye model simulations.



**Figure:** Real part of  $\epsilon(\omega)$ ,  $\epsilon$ , or the permittivity [REU2008].



## Random Polarization

We can define the **random polarization**  $\mathcal{P}(t, \mathbf{x}; \tau)$  to be the solution to

$$\tau \dot{\mathcal{P}} + \mathcal{P} = \epsilon_0 \epsilon_d \mathbf{E}$$

where  $\tau$  is a random variable with PDF  $f(\tau)$ , for example,

$$f(\tau) = \frac{1}{\tau_b - \tau_a}$$

for a uniform distribution.

The electric field depends on the macroscopic polarization, which we take to be the expected value of the random polarization at each point  $(t, \mathbf{x})$

$$\mathbf{P}(t, \mathbf{x}) = \int_{\tau_a}^{\tau_b} \mathcal{P}(t, \mathbf{x}; \tau) f(\tau) d\tau.$$

## Polynomial Chaos

Apply Polynomial Chaos method to approximate the random polarization

$$\tau \dot{\mathcal{P}} + \mathcal{P} = \epsilon_0(\epsilon_s - \epsilon_\infty)E, \quad \tau = \tau(\xi) = r\xi + m$$

resulting in

$$(rM + ml)\dot{\vec{\alpha}} + \vec{\alpha} = \epsilon_0(\epsilon_s - \epsilon_\infty)E\vec{e}_1$$

or

$$A\dot{\vec{\alpha}} + \vec{\alpha} = \vec{g}.$$

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The macroscopic polarization, the expected value of the random polarization at each point  $(t, x)$ , is simply

$$P(t, x; F) = \alpha_0(t, x).$$

# Inverse Problems

## Definition

An **inverse problem** estimates quantities *indirectly* by using measurements of other quantities.

# Inverse Problems

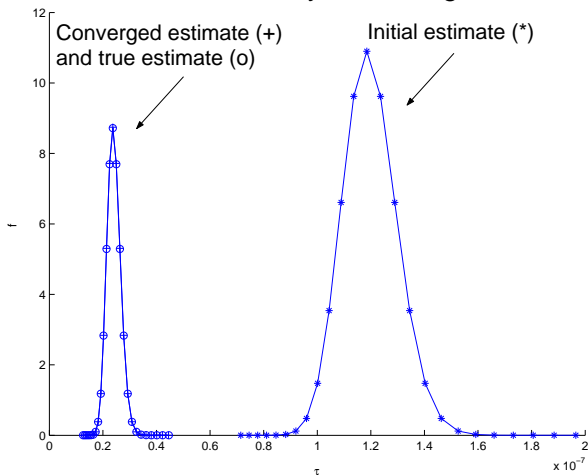
## Definition

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For example, a **parameter estimation** inverse problem attempts to determine values of a parameter set  $q$  given (discrete) observations of (some) state variables.

Examples:

- distance of an object using echo-location (easily invertible)
- amount of oil/water/cave in the ground using RADAR backscatter
- geometry or composition of a defect using measurements of EM fields (CT, MRI, etc.)

Estimated density of  $\tau$  as log normal

Shown are the initial density function, the minimizing density function and the true density function (the latter two being practically identical).

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## Group Members

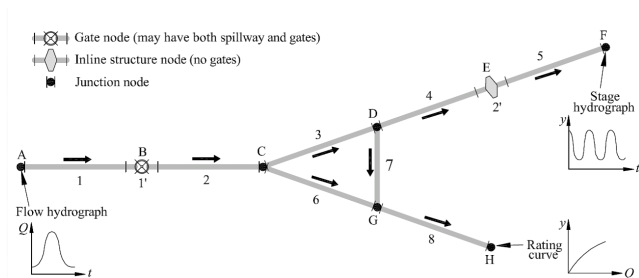
- Department of Civil and Construction Engineering
  - Dr. Arturo Leon (PI)
  - Dr. Duan Chen (Post-doc)
  - Mahabub Alam (Ph.D. student)
  - Luis Gomez-Cunya (Ph.D. student)
  - Parnian Hosseini (Ph.D. student)
  - Christopher Gifford-Miears (MS student)
- Department of Mathematics
  - Dr. Nathan Gibson (Co-PI)
  - Dr. Veronika Vasylykivska (Post-doc)
- Department of Mechanical Engineering
  - Dr. Christopher Hoyle (Co-PI)
  - Matthew McIntire (Ph.D. student)

Funded by BPA Technology Innovation Program: "Towards reduction of uncertainty in the operation of reservoir systems"



# Simple River System

Consider this simple network system



**Unknowns:** flow discharge upstream  $Q_u$  and downstream  $Q_d$ , water surface elevation downstream  $WSE_d$  for each reach  $i = \overline{1, 8}$ .

## Simulation of Unsteady Flows

- Most free surface flows are unsteady and nonuniform.
- Unsteady flows in river systems are typically simulated using one-dimensional models.

**Saint-Venant equations**, PDEs representing conservation of mass and momentum for a control volume:

$$B \frac{\partial y}{\partial t} + \frac{\partial Q}{\partial x} = 0, \quad (1)$$

$$\frac{\partial Q}{\partial t} + \frac{\partial}{\partial x} \left( \frac{Q^2}{A} \right) + gA \left( \frac{\partial y}{\partial x} + S_f - S_0 \right) = 0, \quad (2)$$

where  $x$  is a distance along the channel in the longitudinal direction,  $t$  is time,  $y$  is a water depth,  $Q$  is a flow discharge,  $B$  is a width of the channel,  $g$  is an acceleration due to gravity,  $A$  is a cross-sectional area of the flow,  $S_f$  is a friction slope,  $S_0$  is a river bed slope.

Initial and boundary conditions are required to close the system.

## Objective and Constraints

The broad context of the problem of interest is a PDEs-constrained optimal control problem with uncertainty. In particular, an objective is expected to be optimized by choices of a control functions.

Let  $P$  be a price, and  $E$  be a produced hydro-power energy, then  $R = P \cdot E$  is a revenue.

**Objective:**  $\max R$ ,

Let  $Q_t$  denote a turbine flow,  $Q_s$  a spill flow, and  $S$  a storage.

**Examples of additional constraints:**

$$0 < Q_t^n < Q_{crit},$$

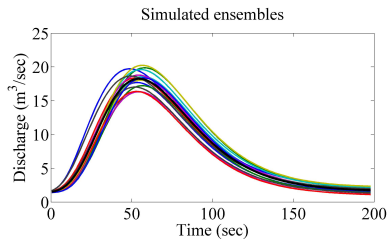
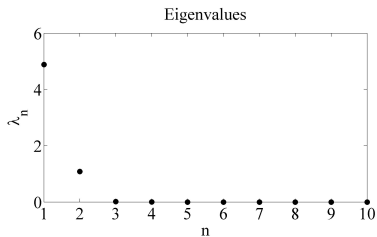
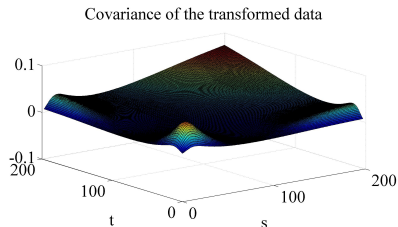
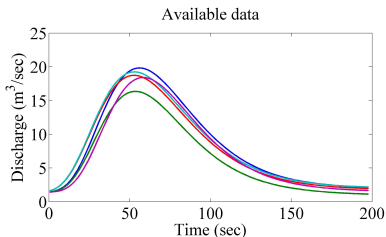
$$0 \leq Q_s^n,$$

$$0 < |Q_t^{n+1} + Q_s^{n+1} - Q_t^n - Q_s^n| < \text{Allowed Value},$$

$$S_{min} < S^n < S_{max}.$$

# Numerical Experiments. Stochastic Parametrizations

## Experiment 1: 5 predictions



# Robust Optimization

The deterministic constrained optimization problem can be formulated as

$$\text{find} \quad \max_q R(q), \quad (3)$$

$$\text{subject to} \quad y(x, t; q) < y_{crit}(x), \quad (4)$$

where  $q$  is a control vector.

We assume that price is deterministic and reformulate our problem as follows

$$\text{find} \quad \max_q \left( E[R(q)] - \alpha \text{Var}[R(q)] \right), \quad (5)$$

$$\text{subject to} \quad P(y(x, t; q) < y_{crit}(x)) > R_0, \quad (6)$$

where  $\alpha > 0$  is a risk tolerance coefficient,  $R_0$  is a reliability level the decision maker wishes to achieve.

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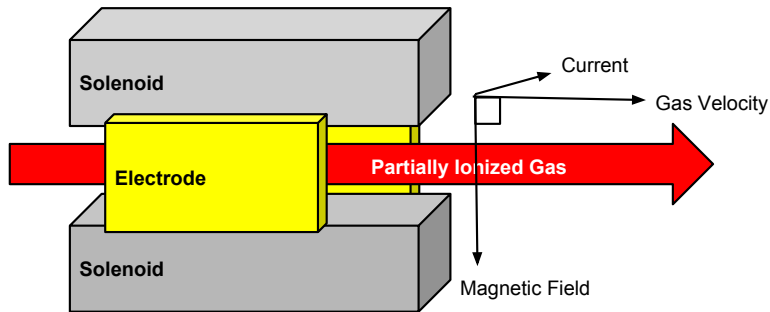
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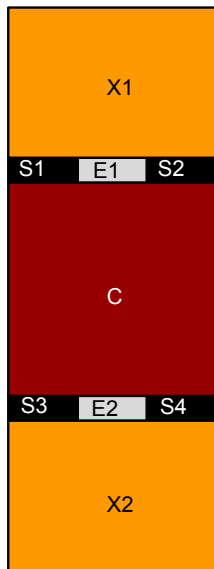
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# MHD generator



Funded by National Energy Technology Laboratory Office of Research & Development Fund: "Applying Computational Methods to Determine the Electric Current Densities in a Magneto-hydrodynamic Generator Channel from External Magnetic Flux Density Measurements", with Vrushali Bokil (Co-PI)

## A slice of simplified generator geometry



-  External Casing (Copper?)
-  Electrode (Ceramics?)
-  Insulating Segment
-  Generator Channel



We assume the fluid-electro-magnetic system model by the following system of ten non-linear, partial differential equations.

$$\rho(\partial_t - \mathbf{u} \cdot \nabla) \mathbf{u} = \mathbf{j} \times \mathbf{b} - \nabla p \quad \text{Momentum} \quad (7a)$$

$$\rho \partial_t \left( \frac{\mathbf{u} \cdot \mathbf{u}}{2} + U + \frac{\epsilon \mathbf{e} \cdot \mathbf{e}}{2} + \frac{\mathbf{b} \cdot \mathbf{b}}{2\mu} \right) \quad (7b)$$

$$+ \rho \mathbf{u} \cdot \nabla \left( \frac{\mathbf{u} \cdot \mathbf{u}}{2} + U \right) = -\nabla \cdot (\mathbf{u} \rho) - \nabla \cdot (\mu^{-1} \mathbf{e} \times \mathbf{b}) \quad \text{Energy} \quad (7c)$$

$$(\partial_t + \mathbf{v} \cdot \nabla) \rho = -\rho \nabla \cdot \mathbf{v} \quad \text{Mass} \quad (7d)$$

$$\partial_t \epsilon \mathbf{e} + \mathbf{j} = \nabla \times \mu^{-1} \mathbf{b} \quad \text{Ampere's Law} \quad (7e)$$

$$\partial_t \mathbf{b} = -\nabla \times \mathbf{e} \quad \text{Faraday's Law} \quad (7f)$$

$$\mathbf{j} = \sigma(\mathbf{e} + \mathbf{u} \times \mathbf{b}) + \frac{\beta}{\sqrt{\mathbf{b} \cdot \mathbf{b}}} \mathbf{j} \times \mathbf{b} \quad \text{Ohm's Law} \quad (7g)$$

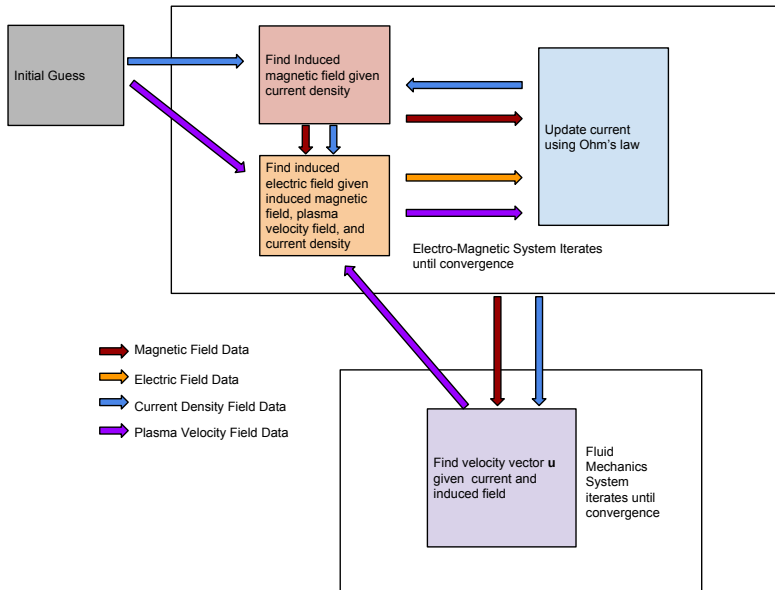
$$\nabla \cdot \epsilon \mathbf{e} = \rho_c \quad \text{Gauss' Law} \quad (7h)$$

$$\nabla \cdot \mathbf{b} = 0 \quad \text{Gauss' Law for Magnetism} \quad (7i)$$

With the following variable definitions:

- $\mathbf{u}$  gas velocity
- $p$  pressure
- $\rho$  gas density
- $U$  thermal energy
- $\mathbf{e}$  electric field
- $\mathbf{b}$  magnetic flux density field
- $\mathbf{j}$  current density field
- $\rho_c$  charge density
- $\epsilon$  electrical permittivity (tensor)
- $\mu$  magnetic permeability (tensor)
- $\sigma$  conductivity (tensor)
- $\beta$  Hall parameter (tensor)

# Fixed Point Approach



# Inverse Problem

