

Momentum operator in the position basis

Consider momentum as a generator of translation
(let's work in 1D for simplicity)

$$\hat{U}_{\Delta x} = \hat{I} - \frac{i}{\hbar} \hat{P}_x \Delta x, \quad \hat{U}_{\Delta x} |x\rangle = |x + \Delta x\rangle$$

Consider an arbitrary state $|\alpha\rangle$:

$$|\alpha\rangle = \int dx |x\rangle \langle x| \alpha\rangle \leftarrow \begin{array}{l} \text{expansion in} \\ \text{terms of eigenbasis} \end{array}$$

$$\begin{aligned} \hat{U}_{\Delta x} |\alpha\rangle &= \int dx \underbrace{\hat{U}_{\Delta x} |x\rangle \langle x| \alpha\rangle}_{\substack{\uparrow \\ \text{Change of} \\ \text{variables} \\ |x + \Delta x\rangle}} = \\ &= \int dx' |x'\rangle \langle x' - \Delta x | \alpha\rangle = \quad \nwarrow \text{Taylor expansion} \\ &= \int dx' |x'\rangle \left(\langle x' | \alpha\rangle - \Delta x \frac{\partial}{\partial x'} \langle x' | \alpha\rangle \right) = \\ &= |\alpha\rangle - \Delta x \int dx' |x'\rangle \frac{\partial}{\partial x'} \langle x' | \alpha\rangle \quad (12,1) \end{aligned}$$

Now go back to left-hand-side of (13.1), (2)

$$\hat{U}_{\Delta x} |\alpha\rangle = (\hat{I} - \frac{i}{\hbar} \hat{P}_x \Delta x) |\alpha\rangle = \\ = |\alpha\rangle - \underbrace{\frac{i}{\hbar} \Delta x \hat{P}_x |\alpha\rangle}_{(13.2)}$$

Compare (13.1) and (13.2) \Rightarrow

$$\hat{P}_x |\alpha\rangle = -i\hbar \int dx' |x'\rangle \frac{\partial}{\partial x'} \langle x'| \alpha\rangle \\ \langle x'' | \hat{P}_x | \alpha\rangle = -i\hbar \int dx' \underbrace{\langle x'' | x'\rangle \frac{\partial}{\partial x'}}_{\delta(x''-x')} \langle x'| \alpha\rangle = \\ = -i\hbar \frac{\partial}{\partial x''} \langle x'' | \alpha\rangle \quad (13.3)$$

If $|\alpha\rangle = |x'\rangle \Rightarrow \langle x'' | \hat{P}_x | x'\rangle = -i\hbar \frac{\partial}{\partial x''} \langle x'' | x'$

$$= -i\hbar \frac{\partial}{\partial x'} \delta(x'-x'')$$

matrix element of

P_x in the x -representation

If $|\alpha\rangle, |\beta\rangle$ are
arbitrary states

$$= -i\hbar \frac{\partial}{\partial x'} \delta(x'-x'')$$

$$\langle \beta | \hat{P}_x | \alpha\rangle = \int dx' dx'' \langle \beta | x'\rangle \langle x' | \hat{P}_x | x''\rangle \langle x'' | \alpha\rangle$$

$$= \int dx' \langle \beta | x' \rangle \left(-i\hbar \frac{\partial}{\partial x'} \right) \langle x' | \alpha \rangle =$$

$$= \underbrace{\int dx' \Psi_\beta^*(x') \left(-i\hbar \frac{\partial}{\partial x'} \right) \Psi_\alpha(x')}$$

Similarly, $\langle x' | P_x^n | \alpha \rangle = (-i\hbar)^n \underbrace{\frac{\partial^n}{\partial x'^n} \langle x' | \alpha \rangle}$

$$\langle \beta | P_x^n | \alpha \rangle = \underbrace{\int dx' \Psi_\beta^*(x') (-i\hbar)^n \frac{\partial^n}{\partial x'^n} \Psi_\alpha(x')}$$

$$\frac{\partial^n}{\partial x'^n} \Psi_\alpha(x')$$

Example

The system is in a state described by a real wave function $\Psi(x)$. Find the expectation value of the momentum (in 1D case), i.e. $\langle \hat{P}_x \rangle$

$$\langle \hat{P}_x \rangle = \langle \Psi | \hat{P}_x | \Psi \rangle = \int_{-\infty}^{+\infty} dx' dx'' \langle \Psi | x' \rangle \langle x' | \hat{P}_x | x'' \rangle.$$

$$\cdot \langle x'' | \Psi \rangle = \int_{-\infty}^{+\infty} \Psi^*(x') \left(-i\hbar \frac{d}{dx'} \delta(x' - x'') \right) \Psi(x'') dx' dx''$$

$$= -i\hbar \int_{-\infty}^{+\infty} dx' \Psi^*(x') \frac{d}{dx'} \Psi(x') = -i\hbar \left(\Psi^* \Psi \Big|_{-\infty}^{+\infty} - \right)$$

integrate
by parts

for a well-
behaved function

$$-\int_{-\infty}^{+\infty} dx' \psi(x') \frac{d\psi^*(x')}{dx'} = i\hbar \underbrace{\int_{-\infty}^{+\infty} dx' \psi(x') \frac{d\psi^*(x')}{dx'}}_{\text{---}} \quad (9)$$

$$\Downarrow$$

$$-i\hbar \int_{-\infty}^{+\infty} dx' \psi^*(x') \frac{d\psi(x')}{dx'} = i\hbar \int_{-\infty}^{+\infty} dx' \psi(x') \frac{d\psi^*(x')}{dx'} \quad \Downarrow$$

Since our $\psi(x)$ is real $\Rightarrow \psi = \psi^*$

$$\Downarrow$$

$$\underbrace{\langle P \rangle}_{} = 0$$

Momentum-space wave function

Consider $|P\rangle$ -basis (for simplicity \Rightarrow 1D case)

Similarly to the position space \Rightarrow

$$\hat{P}|p'\rangle = p'|p'\rangle ; \quad \langle p'|p'' \rangle = \delta(p' - p'')$$

Expansion of an arbitrary state $|\alpha\rangle \Rightarrow$
orthogonality

$$|\alpha\rangle = \int dp' |p'\rangle \langle p' | \alpha \rangle$$

Probability that a measurement of P yields an eigenvalue p' within an interval $\Delta p = dp'$

$$\text{is } \left| \langle p' | \alpha \rangle \right|^2 dp' \quad (\text{similar to the measurement of } X \text{ - see lecture #12})$$

$$\langle p' | \alpha \rangle = \phi_\alpha(p') \leftarrow \text{momentum-space wave function} \quad (5)$$

If $|\alpha\rangle$ is normalized $\Rightarrow \int dp' |\langle p' | \alpha \rangle|^2 =$

$$= \int dp' |\phi_\alpha(p')|^2 = 1$$

What is the connection between x - and p -representations?

Recall: $\langle x'' | \hat{P} | \alpha \rangle = -i\hbar \frac{\partial}{\partial x''} \langle x'' | \alpha \rangle$ (13.3)

Let $|\alpha\rangle = |p'\rangle$ \leftarrow momentum eigenket

$$\underbrace{\langle x'' | \hat{P} | p' \rangle}_{p' | p' \rangle} = -i\hbar \frac{\partial}{\partial x''} \langle x'' | p' \rangle \Rightarrow$$

$$\underbrace{p' \langle x'' | p' \rangle}_{\uparrow} = -i\hbar \frac{\partial}{\partial x''} \langle x'' | p' \rangle$$

differential equation for $\langle x'' | p' \rangle$

$$p' f = -i\hbar \frac{\partial}{\partial x''} f \Rightarrow f = N e^{\frac{i}{\hbar} p' x''}$$

say, $f(x'', p')$

Drop "primes" for simplicity $\Rightarrow \langle x | p \rangle = N e^{\frac{i}{\hbar} px}$

prob. amplitude \xrightarrow{x}
for the momentum eigenstate $|p\rangle$
to be found at position x normaliz. const

Normalization: $\langle x' | x'' \rangle = \delta(x' - x'') =$ (6)

$$= \int dp' \langle x' | p' \rangle \langle p' | x'' \rangle =$$

$$= |N|^2 \int dp' e^{\frac{i}{\hbar} p'(x' - x'')} = |N|^2 \cdot 2\pi\hbar \delta(x' - x'') \Rightarrow$$

\uparrow

$\Rightarrow |N| = \frac{1}{\sqrt{2\pi\hbar}}$ Recall

$$\delta(x - x_0) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ik(x-x_0)} dk$$

say, real

\downarrow

$$\langle x | p \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{\frac{ipx}{\hbar}} \leftarrow \begin{matrix} \text{plane} \\ \text{wave} \end{matrix}$$

How is a position-space wave function $\psi_\alpha(x)$ related to a momentum-space one $\phi_\alpha(p)$? \Rightarrow

$$\underbrace{\psi_\alpha(x) = \langle x | \alpha \rangle = \int dp \underbrace{\langle x | p \rangle}_{\text{!}} \underbrace{\langle p | \alpha \rangle}_{\text{!}} =}_{\text{!}}$$

$$= \int dp \cdot \frac{1}{\sqrt{2\pi\hbar}} e^{\frac{ipx}{\hbar}} \phi_\alpha(p)$$

$$\underbrace{\phi_\alpha(p) = \langle p | \alpha \rangle = \int dx \underbrace{\langle p | x \rangle}_{\text{!}} \underbrace{\langle x | \alpha \rangle}_{\text{!}} = \int dx \cdot \frac{1}{\sqrt{2\pi\hbar}} e^{-\frac{ipx}{\hbar}} \psi_\alpha(x)}$$

\downarrow

$\psi_\alpha(x)$ and $\phi_\alpha(p)$ are Fourier transforms of each other!