

Translation in space

Consider an initial state  $|\vec{x}'\rangle$ , which we want to transform into  $|\vec{x}' + d\vec{x}'\rangle \Rightarrow$

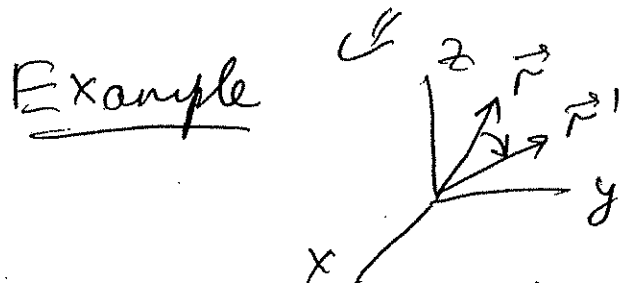
$$\hat{U}_{d\vec{x}'} = \hat{I} - i d\vec{x}' \cdot \hat{G} \quad \hat{U}_{d\vec{x}'} |\vec{x}'\rangle = |\vec{x}' + d\vec{x}'\rangle$$

$\hat{G} = \frac{\hat{\vec{p}}}{\hbar}$  ← momentum operator

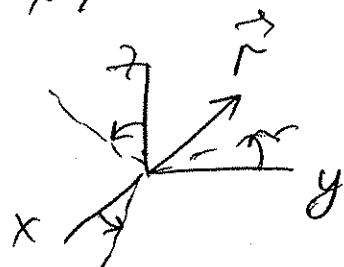
$\uparrow$  infinitesimal

Obviously,  $|\vec{x}'\rangle$  is not an eigenket of  $\hat{U}_{d\vec{x}'}$

Note: Transformation of a state vector itself or operators is an active transformation, whereas transformation of a coordinate system is a passive transformation.



$\vec{r} \rightarrow \vec{r}' \Rightarrow$  active  
(x y z fixed)



$\vec{r}$  fixed  
 $x y z \rightarrow x' y' z' \Rightarrow$  passive

Here we will stick with active transformations



So, the fact that successive transformations <sup>(3)</sup> commute is directly related to the fact that  $[P_i, P_j] = 0$  !  $\Rightarrow$  comes from the fact that translation group is Abelian.

It's not the case for rotations, since  $\hat{G} \sim \vec{J}$ , and  $[J_i, J_j] = i\hbar J_k \neq 0$    
↑  
angular momentum

$\Downarrow$   
rotations about different axes do not commute!

$\Downarrow$   
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Since  $[P_i, P_j] = 0 \Rightarrow P_x, P_y, P_z$  are mutually compatible observables

$\Downarrow$   
form  $|\vec{p}'\rangle = |p'_x, p'_y, p'_z\rangle$   
(common basis)

$\Downarrow$   

$$P_x |\vec{p}'\rangle = p'_x |\vec{p}'\rangle$$

Is  $|\vec{p}'\rangle$  an eigenket of  $\hat{U}_{d\vec{x}'}$ ?  $\Rightarrow$

$$\hat{U}_{d\vec{x}'} |\vec{p}'\rangle = \left( \hat{I} - \frac{i}{\hbar} \vec{P} \cdot d\vec{x}' \right) |\vec{p}'\rangle \stackrel{?}{=} \vec{p}' |\vec{p}'\rangle$$

$$\textcircled{=} \left(1 - \frac{i}{\hbar} \vec{p} \cdot d\vec{x}'\right) |\vec{p}'\rangle \Rightarrow$$

④

$|\vec{p}'\rangle$  is an eigenket of  $\hat{U}_{d\vec{x}'}$

## Wave functions in position and momentum space

1) Wave function in position space

$$\Psi_{\alpha}(x') = \langle x' | \alpha \rangle$$

2) Express  $\langle \beta | \alpha \rangle$  via wave functions:

$$\begin{aligned} \langle \beta | \alpha \rangle &= \int dx' \langle \beta | x' \rangle \langle x' | \alpha \rangle = \\ &= \int dx' \Psi_{\beta}^*(x') \Psi_{\alpha}(x') \end{aligned}$$

3) Expansion in terms of eigenfunctions of  $A \Rightarrow$   
 $\Psi_{\alpha}(x') = \sum_n C_n \Psi_n(x)$ , where  $A \Psi_n(x) = a_n \Psi_n(x)$

$$4) \langle \beta | A | \alpha \rangle = \int dx' \int dx'' \langle \beta | x' \rangle \langle x' | A | x'' \rangle \cdot$$

$$\cdot \langle x'' | \alpha \rangle = \int dx' \int dx'' \Psi_{\beta}^*(x') \langle x' | A | x'' \rangle \Psi_{\alpha}(x'')$$

$$\cdot \Psi_{\alpha}(x'')$$

matrix  
element

Example :  $A = \hat{X}^2$ ,  $\langle x' | \hat{X}^2 | x'' \rangle$ ? (5)

$$\begin{aligned}\langle x' | \hat{X}^2 | x'' \rangle &= \langle x' | x''^2 | x'' \rangle = \\ &= x''^2 \delta(x' - x'') = x'^2 \delta(x' - x'')\end{aligned}$$

Then,  $\langle \beta | \hat{X}^2 | \alpha \rangle = \int dx' \int dx'' \psi_{\beta}^*(x') \cdot x'^2 \delta(x' - x'') \psi_{\alpha}(x'')$

$$= \int dx' \psi_{\beta}^*(x') x'^2 \psi_{\alpha}(x')$$

In general,  $\langle \beta | \underbrace{f(\hat{X})}_{\substack{\uparrow \\ \text{function of the operator } \hat{X}}} | \alpha \rangle =$

$$= \int dx' \psi_{\beta}^*(x') \cdot \underbrace{f(x')}_{\substack{\uparrow \\ \text{function of the eigenvalue}}} \psi_{\alpha}(x')$$

